

A P P E N D I X,

O R,

S U P P L E M E N T

T O T H E

Treatise of Artillery:

C O N T A I N I N G

The true Projectile described by Bodies in the Air.	the greatest Effects; with an Introduction of Fluxions.
The greatest Velocity and Resistance they can have.	To which is added, the true
The most advantageous Length of Guns, and their Charges which produce	Figure of the Earth, deduced from actual Mensuration.

By J O H N M U L L E R,

Preceptor to his Royal Highness WILLIAM Duke
of GLOUCESTER.

L O N D O N:

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APPENDIX,

OR

SUPPLEMENT

TO THE

Treatise of Artillery;

CONTAINING

the greatest Effects; with
an Introduction of Experiments
on the Velocity and Force
of the Shot, and the true
Resistance they can have
in the most advantageous
Position of Guns, and the
Charges which produce
the greatest Effects; with
an Introduction of Experiments
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of the Shot, and the true
Resistance they can have
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Position of Guns, and the
Charges which produce

By JOHN MULLER,

Professor to his Royal Highness William Duke
of Gloucester.

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P R E F A C E.

THE improvements in artillery and gunnery are the chief aim of this work ; but as this could not be done without the highest and most intricate parts of the mathematics ; and as several propositions, which are absolutely necessary in what follows, are not to be found elsewhere ; I was induced to begin with a concise treatise of fluxions ; deduced from the principles laid down by Sir Isaac Newton, the inventor, in his Introduction to the Quadrature of Curves. By this the reader will find every thing that is wanting to understand the subject, without being obliged to have recourse to other authors.

In a subject so intricate as this, it is unsafe to depend on what has been said by others without a very strict examination, to avoid inaccuracies that might mislead and occasion wrong conclusions. The want of this precaution has been the cause that writers upon this subject, have not succeeded so well as might have been expected from their knowledge in other respects.

This work is the production of many years labour, as the great number of capital errors committed by others in their theories required time and application to discover. Thus the specific gravity of air, as commonly supposed, we found not to be true : the common estimation of the resistance of bodies moving in the air is quite void of truth, at least as Mr. Robins will have it ; for he makes that of a 24 pound ball above 23 times more than it possibly can be, as we have proved in this work. Again, the greatest velocity a body can possibly acquire by falling in the air, has not been rightly determined, though it is absolutely necessary to be

known, before projectiles in the air can be determined.

Seeing that the greatest velocity that a shot can possibly have depends on its diameter, I was enabled to determine the most advantageous length of guns and their best charges (a question sought for in vain, ever since their invention) and to correct the vulgar error of making small calibers longer, and loading them with more powder, in proportion to the large, in order, as they imagine, to make their shot go farther; whereby they were obliged to make them heavier than they need be, at a much greater expence in metal and carriage.

From what has been said in this work, artillery may be brought to its greatest perfection by a few good experiments: for all those that have hitherto been published, I may say, are trifling, and not the least to the purpose: the only good are those made by General Williamson, Major Hislop, and several other artillery officers at Minorca, in 1745; which they were so kind to communicate to me, and are inserted in the introduction to our artillery.

The use of fluxions, in the improvement of artillery and gunnery, may possibly startle many artillery officers, but as it cannot be done otherwise, I cannot help it: though an officer may understand his business very well in all other respects, without so complex a theory as this, yet some of them will be pleased to find how they may improve their knowledge; the gallant behaviour of all those employed in the two last wars has been so conspicuous to all the world that they need not be under any apprehension, that the utmost malice or spite could invent can be the least detraction of their merit.

Though the figure of the earth has no immediate connection with what we have treated before, yet I hope the curious reader will not be displeased to see it treated in an easy and familiar manner.

C O N-

CONTENTS.

INTRODUCTION.

SECTION I.

Of fluxions in general.

Art.

5.	<i>THE fluxion of a variable square</i>	—	6
6.	<i>The fluxion of a variable rectangle</i>	—	6
8.	<i>The fluxion of a variable quantity raised to any power</i>	— — —	8
10. 11.	<i>The fluxion of a single quantity raised to any power</i>	— — —	8
14.	<i>The fluxion of a binomial raised to any power</i>		9
	<i>General rules for finding the fluxions of various quantities</i>	— — —	9
16.	<i>The fluxion of curvilinear areas</i>	— — —	11
17.	<i>The fluxion of solids described about an axis</i>		11
18.	<i>The fluxion of solids described about a tangent</i>		12
19.	<i>The fluxion of surface described by curves about an axis</i>	— — —	13
20.	<i>The fluxion of a circular arc</i>	— — —	13
25.	<i>The fluxion of a circular sector</i>	— — —	14
28.	<i>The fluxions of various lines drawn in a figure</i>		15
	<i>Of various orders of fluxions</i>	— — —	16

SECTION II.

The method of finding the greatest and least values of variable expressions

26.	<i>General problem, to determine the properties of such expressions</i>	— — —	17
28.	<i>To divide a given line into two parts, such as the product of any powers shall be the greatest</i>		18
29.	<i>To divide a given line into three parts, such as the product of any powers shall be the greatest</i>		18
30.	<i>Of all triangles of the same base and the same angle at the vertex, the isosceles contains the greatest area</i>	— — —	19
33.	<i>The</i>		

CONTENTS.

33. *The most advantageous situations of braces to support timbers* ———— 20
34. *The most advantageous situation to support timber* ———— 21
36. *The most advantageous angle of a roof, in regard to the strength of its principal rafters* — 21
37. 38. *To find a point within a triangle or a trapezium, so that the lines drawn from the angles to that point shall be the least possible* 22

SECTION III.

The radii of curvature.

39. *To find the radius of curvature* ———— 23
40. *To find the radius of curvature of a parabola* 24
41. *To find the radius of curvature, in an ellipsis or hyperbola* ———— 24
42. *To find the same in another manner* ———— 25

Of the cycloid.

44. *To find the length of a cycloid, or that of any arc* 26
45. *To find the radius of curvature of a cycloid* — 26
46. *To find the evoluta of a cycloid* ———— 47

SECTION IV.

Of fluents.

47. Rule 1. *To find the fluent of a single quantity raised to any power* ———— 27
48. Rule 2. *To find the fluent of a binomial raised to any power* ———— 28
50. *To express a circular arc by the radius and parts of the tangent* ———— 29
51. *The tangents of any two arcs being given, to find the tangent of their difference* ———— 29
53. *Compendious computation of a circular arc to 25 places* ———— 30
56. *The area of an equilateral hyperbola, terminated by two ordinates to the asymptote* ———— 31
57. *An application to logarithms* ———— 32

Examp.

CONTENTS.

Examp. The logarithm of two computed	33
60. The area of an hyperbolic sector expressed by parts of the tangent	34
61. The same area expressed by parts of the ordinate to the first axis	35
63. The area of an elliptic sector expressed by parts of the tangent	36
64. The same area expressed by the ordinate to the first axis	37
66. General fluent of a binomial	38
71. General fluent of the same expressed in a different manner	41
73. The fluent of a binomial expression, when the variable quantity vanishes	44
Sir Isaac Newton's method of differences.	
74. General problem for finding the equation of the curve	45
75. From the given equation of the curve to find the area	45
76. Table of areas expressed by ordinates	46
General method for finding the root of equations by approximation	48
74. General rule to find the roots, with several particular examples	50
80. Of exponentials and their fluxions	54

SECTION V.

Of motion in general considered in a non-resisting medium	55
80. The fluxion of the velocity is as the fluxion of the time and the measure of the force that generates it conjunctly	58
81. The fluxion of the space described is as the fluxion of the time and the velocity conjunctly	59
88. General problem of projectiles, when the ordinates pass through a given point	62
92. The same when the ordinates are parallel	64
98. When	

CONTENTS.

98. *When the force is constant and the ordinates parallel, the projectile is a parabola* ——— 66
101. *The time of describing any arc of a cycloid, by the force of gravity, is always the same* — 68
103. *The length of a pendulum is to the force of gravity, as the square of the diameter is to the square of the circumference* ——— 69
107. *The length of a pendulum vibrating seconds in latitude $51^{\circ} : 32'$ of London, is 39.125 inches, and the space fallen through in a second, 16.088 feet, English measure* ——— 70
- The length of such a pendulum, in the latitude $48^{\circ} : 51'$ of Paris, is 440.567 lines, and the space fall through in a second 15.097 feet, French measure* ——— 70
109. *When the curve described is a conic section, the centripetal force, which tends to the focus, is as the square of the distance inversely* — 71
112. *When the centripetal force is as the square of the distance inversely, the curve described will be a conic section, and the directions of the force passes through the focus* ——— 72
113. *The periodic time of a body revolving in an ellipsis* ——— 73
114. *The periodic time of a body revolving in a circle, and the several comparisons between the centrifugal force, time, and velocity* ——— 73
- From thence follows, from the known periodic times of the planets, 1^o. their mean distances from the sun; 2^o. their diameters; 3^o. their solid contents; 4^o. their quantity of matter; 5^o. their attractions of equal bodies placed on their surfaces; 6^o. and lastly, the densities of the Sun, Saturn, Jupiter, and the Earth* — 75
- Problems of the greatest and least quantities.*
119. *Lemma, in which the general properties of such quantities are generated* ——— 78
125. To

CONTENTS.

125. To find the least time in which a line passes thro' two given points, is described — — 80
126. To find the greatest area contained within the same bounds — — 81
127. To find the greatest solid contained within the same bounds — — 82

SECTION VI.

The motion of fluids and the resistance of bodies.

130. The force with which a fluid strikes a plane obliquely is as the square of the sine of the incident angle — — 83
132. If a machine be moved by water, to find the velocity of the wheel so as to produce the greatest effect possible — — 84
135. The angles which the direction of the wind makes with the axis and sails of a windmill being given, to find the force with which the mill turns: the same thing in ships — — 85
138. To find the time and velocity of water moving in a crooked tube — — 89
140. The density of the air being as the force of compression, to find the density at any given height 90
The particle of air being attracted by forces proportional to the squares of the distances inversely from the center of the earth, to find the density in any given place — — 91
144. The density and elastic force of a medium being given, to find the velocity and time of a pulse produced by a sonorous body — — 92
- Remark, in which the motion of sound and the specific gravity of air are determined — 95
148. The dimensions of a cylindric vessel constantly full of water being given, as well as the diameter of the orifice in the bottom, to find the quantity of water run out in a given time 97
150. To find the pressure of the water on a small circle placed in the middle of the orifice — 100

CONTENTS.

155. *If a globe moves uniformly in a compressed, infinite, and non-elastic fluid, its resistance will be to the force by which his whole motion may be generated or destroyed, in the time it moves over a space of six times its length, as the density of the fluid is to the density of the globe* — — — 104
156. *The resistance of globes in an infinite, compressed, non-elastic medium, is in the compound ratio of the square of the velocity, of the square of the diamenter, and the density of medium directly, and the length of the globe inversely* — — — 104
158. *The greater velocity a body can acquire in a resisting medium, is equal to that it may acquire by falling through a space in the same medium, that is to three of its diameter as the density of the body is to the density of the medium* — — —

SECTION VIII.

- Of motion in a resisting medium* — 107
160. *The fluxion of the velocity is as the fluxion of the time and difference between the generating power and resistance conjointly* — — 111
161. *The fluxion of the space described is as the fluxion of the time and velocity conjointly* 112
163. *If a body projected from a given point in a given direction, with a given velocity, by a centripetal force in a medium of a given resistance, to find the nature of the curve* — 112
166. *If a body falls from a given height by the force of gravity, and the resistance is as the square of the velocity, to find the velocity and time* 113
167. *If a body be projected upwards with a given velocity, acted upon by the force of gravity, and the resistance is as the square of the velocity, to find the velocity and time at any given height* — — — 115

CONTENTS.

Remark, <i>A falling body can never acquire the velocity which we call the greatest possible, but may continually approach it</i>	116
168. <i>If a body moves in a canal, or horizontal right line, to find the velocity, time, and space moved over</i>	117
169. <i>The velocity with which the body sets out can never be equal to the greatest possible</i>	118
170. <i>Suppose the powder all fired before the shot is sensibly moved from its place, to find the velocity with which the shot is discharged</i>	118
171. <i>The resistance of a body arising from its motion in the medium can never be equal, much less greater, than the weight of the shot</i>	119
<i>The force of gravity is to the force of powder as the diameter of the shot is to 27810 nearly</i>	120
172. <i>To find the velocity of an 18 pound shot, being charged with 9 pounds of powder</i>	121
<i>Three Remarks relating to the most advantageous length of guns and their greatest charge, from 121 to</i>	130
173. <i>If a body projected with a given velocity in a given direction, to find the nature of the curve described by the body</i>	131
FIRST METHOD <i>of solving the problem by an infinite series</i>	133
<i>Contrivance to make the infinite series always converge</i>	135
179. <i>The angle of elevation being 45 degrees, and the velocity the greatest possible, to find the horizontal range</i>	136
180. <i>To find the abscissa corresponding to the greatest ordinate</i>	136
181. <i>To find the greatest ordinate</i>	137
Remark, <i>Where is shewn the agreement between our theory with the best experiments</i>	137
b 2	183. To

CONTENTS.

183. To find the velocity at the highest point and the time — — — 139
184. To find the velocity and angle in which the body strikes the ground — — — 140
185. The projectile velocity being one half of the greatest possible, and the rest the same as before, to find the horizontal range — — — 141
186. To find the abscissa corresponding to the greatest ordinate — — — 141
187. The projectile velocity being one tenth of the greatest possible, and the rest as before, to find the horizontal range — — — 141
188. To find the abscissa corresponding to the greatest ordinate — — — 142
190. The angle of elevation being 30° , and the rest as before, to find the horizontal range — — — 143
191. The angle of elevation being 60° , and the rest the same as before, to find the horizontal range — — — 143
192. The inclination of a plane being given, to find the point where the shot will hit it — — — 143
193. The angle of elevation and the horizontal range being given, to find the ratio of the greatest and projectile velocities — — — 144
196. The ratio between the greatest and projectile velocities being given, as well as the horizontal range, to find the sine of elevation — — — 145

SECOND METHOD of solving the same Problem by the differences.

197. The ratio of the greatest and projectile velocities being given, together with the angle of elevation, to find the velocity at the highest point — — — 146
200. The angle of elevation being 45 degrees, and the projectile velocity the greatest possible, to find the velocity at the highest point — — — 148
201. To find the abscissa corresponding to the greatest ordinate — — — 149

CONTENTS.

202. To find the greatest ordinate — — 150
 203. To find the remainder of the horizontal range 150

THIRD METHOD.

204. The projectile velocity at the vertex being given,
 to find the abscissa corresponding to the great-
 est ordinate — — — 152
 205. The angle of elevation being 45 degrees in the
 case above — — — 153
 206. To find the remainder of the horizontal range 153

SECTION VIII.

*The true figure of the earth determined from actual
 measures — — — 155*

207. A particle of matter placed on the surface of a
 sphere is attracted towards the center by a
 force proportionally to the radius — — 156
 209. If a sphere, whose particles are freely disposed,
 revolves about a diameter as an axis, with
 such a force as may be compared to that of
 attraction, the sphere will be changed into a
 spheroid — — — 157
 213. The direction of gravity on the surface of a
 planet is every-where perpendicular to the
 surface — — — 159
 214. The attraction towards the center of a particle
 placed on the surface of a planet is as its
 distance from the center — — 160
 218. The ratio between the actual gravities in two
 different latitudes being given, to find the
 ratio between the arcs of the ellipsis — — 162
 221. From the given ratio between the force of gra-
 vity at Paris and at the equator, to find the
 ratio between the axis and the diameter of the
 equator — — — 164
 222. From the known meridian degrees in two dif-
 ferent latitudes, to find the ratio of the
 axis to the diameter of the equator — — 164
 225. From the known meridian degrees at the equa-

tor

CONTENTS.

- tor and near the pole, to find the ratio of the
axis to the diameter of the equator — — 166
226. The meridian degree at the equator being given,
to find that at Pellow in the lat. of $66^{\circ} : 20'$ 167
227. The mer. deg. at the equat. being given, to find
that at Paris in the lat. of $40^{\circ} : 50'$ — 167
- 228 The mer. deg. at the equat. being given, to find
that at London in the lat. of $51^{\circ} : 32'$ — 168
229. The mer. deg. at the equat. being given, to find
that at Edinburgh in the lat. of $55^{\circ} : 47'$ 168
230. The mer. deg. at the equat. being given, to find
that at Dublin in the lat. of $53^{\circ} : 11'$ — 169
231. The length of a pendulum vibrating seconds in
any latitude being given, to find the length of
a like pendulum in any other latitude — 169
233. The length of a pendulum vibrating seconds at
the equator being given, to find the length of
a like pend. at Pellow in the lat. of $66^{\circ} : 48'$ 170
234. The length of a pend. vibrating seconds at the
equat. being given, to find the length of a like
pend. at London in the lat. of $51^{\circ} : 32'$ 170
235. The length of a pend. vibrating seconds at the
equat. being given, to find that of a like pend.
at Edinburgh in the latitude of $55^{\circ} : 47'$ 171
236. The length of a pend. vibrating seconds at the
equat. being given, to find that of a like pend.
at Dublin in the latitude of $53^{\circ} : 11'$ — 171
237. The length of a pend. vibrating seconds at the
equat. being given, find that of a like pend.
at the pole — — — 171
238. To find the length of the radius of the equator,
and that of half the axis — — 172
239. To find the force of gravity at the equator — 173
240. To find the centrifugal force at the equator — 173
- Observation on the authors who have wrote on
the figure of the earth — — — 174

THE FIRST PRINCIPLES OF FLUXIONS.

SECTION I.

Of FLUXIONS in General.

HOUGH Fluxions have been treated by many able authors, and their application in a very copious manner; yet, if I am not mistaken, the first principles have not hitherto been explained in so easy and perspicuous a manner as might be done. Two different methods are used to demonstrate them; the one by the laws of motion, and the other by the prime and ultimate ratios. The first requires more previous knowledge, and is not so applicable in many cases; for instance, in finding the fluxion of a curve line, or that of its surface, described about an axis, as any one may be convinced who reads *Maclaurin's* work upon this subject. The second is that which has been used by the inventor Sir *Isaac Newton*, as may be seen in the introduction to his quadrature of curves. But those authors who use this method, have generally mistaken the limiting ratio for the prime and ultimate ones, such as Mr. *Robins* and many others, which has occasioned so much disputes about the certainty of their principles.

In the method of fluxions, quantities are not considered as made up by the accumulation of parts, but as generated by a continual flux or motion; that is, a line described by a point, a surface by a line, and a solid by a surface.

DEFINITION I.

The velocities or rates with which quantities, varying by a continual motion, increase or decrease at the same time, are called *fluxions*, and the quantities themselves *fluents*.

DEFINITION II.

The ultimate *ratio* of variable quantities, which vanish or cease to exist at the same time, is the last ratio they can have, and in which they vanish *; and the prime ratio, that in which quantities enter together into existence.

Sir *Isaac Newton* observes, that the fluxions of variable quantities, are as the ultimate ratios of their cotemporary increments. But as this inference may not readily appear to beginners, we shall endeavour to prove it in the two following theorems.

THEOREM. Plate I. Fig. 1.

1. The difference mn or pP , between two abscissas AP , Ap , of a curve AM , and the difference nM between the

* This definition of the ultimate ratio is so exact and precise, that one should think it impossible to mistake its meaning; since it requires that the quantities should absolutely vanish at the same time, and therefore all such quantities as do or cannot actually vanish, tho' they approach ever so near, can have no ultimate ratio: yet the limiting ratio, which may approach ever so near, but the quantities can never vanish, has been confounded with the ultimate ratio. To shew the difference between these two ratios, let a be a constant, and x a variable quantity; then it is said, that the ratio of a to $a+x$, becomes that of equality, when x becomes infinitely small or less than any assignable quantity. Is not this supposing $a \doteq a+x$? It is in vain to say that x can make no difference on account of its smallness; it must be either something or nothing, since there can be no state between existence and non-existence; and if x be something, tho' ever so small, a can never be equal to $a+x$, in a mathematical sense. Tho' this supposition were admitted, it must be owned, that, to establish the principles of so extensive a science as fluxions upon this foundation, must appear very improper, as it introduces a quantity of which we can have no possible idea; and, what is still more absurd, and against good method, when there is not the least occasion for it.

corre-

corresponding ordinates PM, pm, vanish in the ratio of the subtangent TP to the ordinate PM.

Let the line drawn thro' the extremities M, m of the ordinates meet TA in V; then the similar triangles VPM, mnM, give $mn : nM :: VP : PM$; and as this proportion always subsists, where-ever the ordinate pm may be drawn, if we suppose pm to move parallel to itself towards PM, it is plain, that as pm approaches PM, or the point p to the point P, the point V approaches also to the point T, and the lines mn, nM diminish gradually till the point p coincides with the point P, then they vanish together, and the point V in that instant, and not before nor after, coincides with the point T: consequently, by the proportion above, the ultimate or vanishing ratio of mn to nM, will be equal to the ratio of the subtangent TP to the ordinate PM.

N. B. As but one line can be drawn so as to touch a curve in the same point, so can there be but one ultimate ratio between the differences of the increments or decrements of the abscissas and ordinates of the same curve.

It may also be observed, that since $mn : mM :: VP : VM$, the ultimate ratio of mn to the arc mM is equal to the ratio of the subtangent TP to the tangent TM; for if a tangent be drawn thro' the point m, it is easily proved that the sum of the tangents at m and M, terminated by their intersection, will always be greater, and the chord mM less, than the arc mM; and since the sum of these tangents and the chord mM vanish in the ratio of equality, the arc mM and that chord vanish also in the same ratio: consequently the ultimate ratio of mn to mM is equal to the ratio of the subtangent TP to the tangent TM.

THEOREM. *Fig. 2.*

2. *If a line pE moves parallel to itself with an uniform motion, while a point m in that line moves with a*
B 2
retarded

retarded one so as to describe a curve AM; when that line arrives into any position PM, the velocity of the point p is to the velocity of the point m, as the subtangent is to the ordinate.

Let TP be taken to PM, as the velocity of the point p with which it arrives at P, is to the velocity the point m arrives with at M; thro' the point M draw EF parallel to AP, and let the ordinates pm, qn drawn on both sides of PM, meet EF, in E, F and TM, in L, N: Then as the velocity of the point m is greater at m and less at n than at M, by supposition, while the line pE moves from p to P, and from P to q, or the point E from E to M, and from M to F, uniformly, the point m which describes the curve, moves over mE with a greater and over Fn with a less velocity than that point has at M; and since $(TP:PM::) EM:EL::MF:FN$; by supposition the spaces EL, FN, are described uniformly with a velocity equal to that which the point m has at M, at the same time that EM, MF are described uniformly, with the velocity of the point p; and as mE is described with a greater velocity, and Fn with a less, than that with which LE and FN are described at the same time, LE will be less than Em, and FN greater than Fn. The points L, N, are therefore on the outside of the curve AM; and as this will always happen where-ever the ordinates pm, qn, are drawn, excepting at the point M, the line TM touches the curve at M. The velocity of the point p at P is consequently to the velocity of the point m, at M, or the fluxion of the abscissa AP is to the fluxion of the ordinate PM, as the subtangent TP is to the ordinate PM.

If the motion of the point m was accelerated instead of being retarded, the curve would be convex towards the line AP; but the demonstration will remain the same as before.

We have here supposed the common laws of motion, as being treated of in books of geometry,
and

and therefore supposed to be known; but if any reader is unacquainted with them, he may pass over this last theorem, as being only of use to shew the agreement between fluxions and ultimate ratios.

C O R.

3. Hence, the ultimate ratio of the cotemporary increments or decrements of the abscissa AP and ordinate PM of any curve is then the same, or equal to the ratio of the fluxions of these lines, agreeable to Sir *Isaac Newton's* assertion: the demonstrations deduced from the ultimate ratios are consequently equally conclusive and geometrical, as those from fluxions, and in many cases more concise and convenient, as will appear hereafter.

Observe, that no actual motion is required in the method of fluxions, but only the proportion, whereby variable quantities increase or decrease at the same time, which is always assignable from the known relation between the quantities themselves. The word *velocity* is applied in a figurative sense, to denote the degree of increase and decrease at the same time,

C O R.

4. Since the direction of the motion in any point of the curve is always in the tangent at that point, and if the motion with which the point *m* arrives at *M*, was to become uniform, the point *m* would proceed in the direction of the tangent *TM*; the directions of the motions in the abscissa AP, ordinate PM, and the curve, being in the subtangent TP, ordinate PM and tangent TM; the velocities, or fluxions of the abscissa, ordinate and curve, may therefore be represented by the three sides of the triangle TPM, or by those of any other triangle similar to this.

Fig. 3. No. 2. Hence if AM be an arc of a circle, O the center, PM the sine: then because of the similar right-angled triangles OPM, MPT, we have, 1°. OP: OM::PM:MT, that is, the cosine is to the radius as the fluxion of the sine is to the fluxion of the arc AM.

B 3

2°. PM;

2°. $PM:OM::PT:MT$, that is, the sine is to the radius as the fluxion of the cosine is to the fluxion of the arc AM .

N.B. Constant or invariable quantities have no fluxions, and are generally expressed by the initial letters a, b, c, d of the alphabet; the variable or flowing ones, by the finals v, x, y, z , and their fluxions by the same letters with a point over them, as $\dot{v}, \dot{x}, \dot{y}, \dot{z}$.

A X I O M.

Equal magnitudes have equal fluxions, since they must necessarily increase or decrease at an equal rate at the same time.

T H E O R E M.

5. *If the sum of several magnitudes is equal to one or more magnitudes, the sum of their fluxions will be equal to the fluxions of their equals.*

This is evident from the axiom; for if $v+y=x+z$, then will $\dot{v}+\dot{y}=\dot{x}+\dot{z}$, otherwise these magnitudes could not increase or decrease at an equal rate.

T H E O R E M.

6. *When a magnitude decreases while others increase, its fluxion must be negative.*

For if $y+z=a$, and a is a constant magnitude; then when y increases z must decrease, or if y decreases z must increase: and since the fluxion of $y+z=a$ is $\dot{y}+\dot{z}=0$, because a is constant, we have $-\dot{y}=\dot{z}$, or $\dot{y}=-\dot{z}$.

T H E O R E M.

7. *If variable magnitudes are in a constant ratio, their fluxions will be in the same ratio.*

For since variable magnitudes, which increase or decrease in a constant ratio, their rates or fluxions must be in the same ratio: thus, if a, b are constant, and x, y variable, and if $a:b::x:y$, then x and y must vanish together, and in the ratio of a to b .

P R O-

PROBLEM. *Fig. 1.*

8. To find the fluxion of a variable square, yy .

Let y express the ordinate PM of a parabola whose parameter is unity, and abscissa $AP = x$; then will $x = yy$, and because the subtangent TP is double the abscissa AP by conicks, and the fluxion of the abscissa AP is to that of the ordinate PM , as TP is to PM , by *Art. 2.* we have $TP (2x) : PM (y) :: \dot{x} : \dot{y}$, or $2x\dot{y} = y\dot{x}$, or because $x = yy$, by writing the value of x into $2x\dot{y}$, we get $2y\dot{y}^2 = y\dot{x}$, which divided by y gives $2\dot{y}y = \dot{x}$; but $\dot{x}$ is the fluxion of x , therefore $2\dot{y}y$ is that of its equal yy by the axiom.

Otherwise by the ultimate ratios. *Fig. 1.*

Let $Ap = u$, $pm = z$; then will $u = zz$, by the property of the parabola; this equality subtracted from $x = yy$ gives $x - u = yy - zz$; whence $mn (x - u) : nM (y - z) :: y + z : 1$; and by the similarity of triangles $mn : nM :: VP : PM$; therefore $y + z : 1 :: VP : PM$, by equality of ratios. Now, when Ap becomes AP and pm becomes PM , or $u = x$ and $y = z$, then VP becomes the subtangent TP , and consequently $2y : 1 :: (TP : PM ::) \dot{x} : \dot{y}$, by *Art. 1.* or $\dot{x} = 2\dot{y}y$.

C O R.

9. Hence, the fluxion of a variable rectangle yz is $\dot{y}z + y\dot{z}$. For if we suppose $y + z = x$, then will $\dot{y} + \dot{z} = \dot{x}$ by *Art. 5.* and $y + z = x$ multiplied by $\dot{y} + \dot{z} = \dot{x}$, gives $y\dot{y} + \dot{y}z + y\dot{z} + z\dot{z} = x\dot{x}$, and the square of $y + z = x$ is $yy + 2yz + zz = xx$. Hence, since $y\dot{y}$, $z\dot{z}$, $x\dot{x}$, are half the fluxions of the squares yy , zz , xx , by *Art. 8.* it is manifest that $\dot{y}z + y\dot{z}$ is half the fluxion of $2zy$, that is of yz by the axiom.

C O R.

10. The fluxion of y^3 is $3y\dot{y}^2$. For suppose $x = yy$; then $xy = y^3$ by equal multiplication, and the fluxion $\dot{x}y + x\dot{y}$ of xy is equal to the fluxion of y^3 its equal, by the axiom; but the fluxion of $x = yy$, is $\dot{x} = 2y\dot{y}$, by *Art. 8.* and the values of x , $\dot{x}$, being wrote into the

fluxion $\dot{x}y + x\dot{y}$, gives $2\dot{y}y^2 + \dot{y}y^2$, or $3\dot{y}y^2$, for the fluxion of y^3 .

The fluxion of y^4 is $4\dot{y}y^3$; for if $x = yy$, which squared gives $xx = y^4$, and the fluxion $2x\dot{x}$ of x^2 will be equal to the fluxion of y^4 , by the axiom: now the fluxion of $x = yy$ is $\dot{x} = 2y\dot{y}$, by *Art. 8.* and these values of $x, \dot{x}$, wrote into the fluxion $2x\dot{x}$, gives $4\dot{y}y^3$ for the fluxion of y^4 .

In the same manner the fluxion of y^5 is found to be $5y^4\dot{y}$, by supposing $x = y^4$, and that of y^6 is $6y^5\dot{y}$, by making $x = y^5$: in general the fluxion of any quantity y^n , whose index n is any whole positive number, is $n\dot{y}y^{n-1}$, that is equal to the product of the index n multiplied by the fluxion $\dot{y}$ of y , and this product by the quantity y^n , whose index is diminished by unity.

C O R.

11. The fluxion of yz^n is $\dot{y}z^n + ny\dot{z}z^{n-1}$. For if $x = z^n$, then $xy = yz^n$, and the fluxion $\dot{x}y + x\dot{y}$ of xy , by *Art. 9.* will be equal to the fluxion $\dot{y}z^n$, by the axiom; but the fluxion of $x = z^n$, is $\dot{x} = n\dot{z}z^{n-1}$, by the last; and the value of $x, \dot{x}$ wrote into $\dot{x}y + x\dot{y}$, gives $\dot{y}z^n + ny\dot{z}z^{n-1}$, for the fluxion of yz^n .

Hence the fluxion of y^{-n} is $-n\dot{y}y^{-n-1}$, for if $x = -y^n$, then $xy^n = 1$, whose fluxion $\dot{x}y^n + ny\dot{x}y^{n-1}$ is $= 0$, or dividing by y^n , $\dot{x} = -n\dot{y}y^{-1}$; or because $x = -y^n$, we got $\dot{x} = -n\dot{y}y^{-n-1}$.

C O R.

12. The fluxion of $y^{\frac{n}{r}}$ is $\frac{n}{r}\dot{y}y^{\frac{n}{r}-1}$. For if $x = y^{\frac{n}{r}}$, then by raising both sides of the equation to the power r , $x^r = y^n$; whose fluxion is $r\dot{x}x^{r-1} = n\dot{y}y^{n-1}$: the first side being divided by $\frac{r}{x}x^r$, and the second by its equal $\frac{r}{x}y^n$, gives $\dot{x} = \frac{n}{r}\dot{y}y^{-1}$, or because $x = y^{\frac{n}{r}}$, we get $\dot{x} = \frac{n}{r}\dot{y}y^{\frac{n}{r}-1}$ for the fluxion required, which
shews

shews that the fluxion of y^n is always $n\dot{y}y^{n-1}$, whether the index n is a whole or broken rational number, positive or negative.

C O R.

13. The fluxion of xyz , is found by making $u=xy$; for then $uz=xyz$, and the fluxion $\dot{u}z+u\dot{z}$ of uz by Art. 9. is equal to the fluxion of xyz , by the axiom; but the fluxion of $u=xy$ is $\dot{u}=\dot{x}y+x\dot{y}$, and these values of u and $\dot{u}$, being wrote into the fluxion $\dot{u}z+u\dot{z}$, gives $\dot{x}yz+x\dot{y}z+xy\dot{z}$ for that required. Or if $xyz=u$, then the fluxion $\dot{x}yz+x\dot{y}z+xy\dot{z}=\dot{u}$, divided, the first side by xyz , and the second by its equal u , gives

$$\frac{\dot{x}}{x} + \frac{\dot{y}}{y} + \frac{\dot{z}}{z} = \frac{\dot{u}}{u}.$$

C O R.

14. The fluxion of $\overline{a+cz^n}^r$ is found by supposing $x=a+cz^n$, whose fluxion is $\dot{x}=nc\dot{z}z^{n-1}$, and $x=a+cz^n$ raised to the power $r-1$, and multiplied by r , gives $rx^{r-1}=r \times \overline{a+cz^n}^{r-1}$. The first side multiplied by $\dot{x}$, and the second by its equal $nc\dot{z}z^{n-1}$, gives $r\dot{x}x^{r-1}=ncr\dot{z}z^{n-1} \times \overline{a+cz^n}^{r-1}$ for the fluxion of x^r or of its equal $\overline{a+cz^n}^r$.

G E N E R A L R U L E S.

I. To find the fluxion of a single quantity y^n , whose index n is any rational number.

Multiply that quantity by the index n , and change one of its dimensions into the fluxion $\dot{y}$.

II. To find the fluxion of a binomial $\overline{a+cz^n}^r$, whose index r is any rational number.

Multiply the fluxion of the binomial by the index r , and this product, by the binomial raised to a power whose index is unity, less than that of the proposed one.

III. To find the fluxion of any product xyz , of several variable quantities.

Multiply the fluxions $\dot{x}$, $\dot{y}$, $\dot{z}$, of these quantities, each by the product of the others, yz , xz , xy , and the sum

sum of all these products will be the fluxion required.

EXAMPLES.

The fluxion of $au + bx^2 + cy^3 + dz^4$, is $a\dot{u} + 2b\dot{x}x + 3c\dot{y}y^2 + 4d\dot{z}z^3$, by Rule I.

The fluxion of $\sqrt{aa + xx}$, or $(aa + xx)^{\frac{1}{2}}$, is $\frac{1}{2} \times 2\dot{x}x \times (aa + xx)^{\frac{1}{2}-1}$, or $\frac{\dot{x}x}{\sqrt{aa + xx}}$ by Rule II.

The fluxion of $\sqrt{ax + xx}$ or $(ax + xx)^{\frac{1}{2}}$, is $\frac{1}{2}a\dot{x} + \dot{x}x \times (ax + xx)^{\frac{1}{2}-1}$, or $\frac{a\dot{x} + 2\dot{x}x}{2\sqrt{ax + xx}}$ by the same Rule.

And the fluxion of $x\sqrt{aa + xx}$, is $\dot{x}\sqrt{aa + xx} + \dot{x}x^2 \times (aa + xx)^{-\frac{1}{2}}$, by Rule III.

In the same manner may be found the fluxion of any multinomial, by following the same steps as in *Art.* 14. or that of a multinomial multiplied by any other quantity, by *Art.* 13.

THEOREM. *Fig.* 4.

15. *Whatever the magnitude may be which has but one variable dimension, its fluxion will always be in the compound ratio of the fluxion of the variable dimension, and of the constant dimensions.*

This is manifest, because the fluxion of ax is $a\dot{x}$, whatever constant dimensions a may express.

Thus if a given line PM, by moving along another line AD, parallel to itself, describes a rectangle AC; the fluxion of any part ABMP, will be equal to the rectangle made by that line and the fluxion of the part AP, moved over by the point P. If the rectangle AC revolves about the base AD as an axis, so as to describe a cylinder, the fluxion of any part described by the rectangle AM, will be equal to a cylinder whose base is the circle described by the given line PM, and altitude the fluxion of AP. In the same manner, the fluxion of the cylindric surface, described

scribed by the variable line BM, will be equal to the rectangle made by the circumference described by the radius PM, and the fluxion of the variable line BM, or that of its equal AP.

P R O B L E M. *Fig. 5.*

16. *To find the fluxion of the area AMP, terminated by the curve AM, and the ordinate PM, perpendicular to the abscissa AP.*

Draw AB, of any given length, perpendicular, and BC parallel to AP, and let the ordinate QN, PM, produced meet BC in q and p: then if the ordinate mn, be such, that the area QM be equal to the rectangle made by the ordinate mn and the base QP; the difference QM between the the areas AMP, ANQ, will be to the difference qP between the rectangle Ap, Aq as $mn \times QP$ is to $Pp \times QP$, or as mn is to Pp. Now if the ordinate QN moves parallel to itself till it coincides with PM, it is evident that mn, which is always greater than QN, and less than PM, will likewise coincide with PM. The ultimate ratio of the area QM, to the rectangle Qp, or the fluxion of the area AMP to the fluxion of the rectangle Ap, is therefore as PM to Pp, by *Art. 3.* But the fluxion of the rectangle Ap is equal to the rectangle made by the given line Pp and the fluxion of AP, by *Art. 15.* the fluxion of the area AMP, is consequently equal to the rectangle made by the ordinate PM and the fluxion of the abscissa AP.

P R O B L E M.

17. *To find the fluxion of the solid described by the area AMP, about the axis AP.*

If the ordinate mn be such, that the cylinder, whose base is the circle, described by that ordinate and altitude the line QP, be equal to the solid described by the space QM: then the solid described by the space QM will be to the cylinder, described by the rectangle Qp, about the same axis AP, as the circle described by

by mn is to the circle described by Pp , and by a similar argument as in the last, the ultimate ratio of the solid described by the space QM , to the cylinder described by the rectangle Qp , or the fluxion of the solid described by the area AMP , is to the fluxion of the cylinder described by the rectangle Ap , as the circle described by the ordinate PM is to the circle described by Pp . But the fluxion of the cylinder is equal to the cylinder whose base is the circle described by Pp and altitude the fluxion of AP , by *Art. 15*; the fluxion of the solid described by the area AMP , about the axis AP , is consequently equal to the cylinder whose base is the circle described by the ordinate PM and altitude the fluxion of AP .

P R O B L E M.

18. *To find the fluxion of the solid described by the area AMP , about the tangent AT as an axis.*

If the ordinate mn be such, that the solid described by the space QM be equal to the cylinder whose surface is that described by mn , and base by Q' ; then the solid described by the space QM is to the solid described by the rectangle Qp , about the same axis, as the surface described by mn is to the surface described by Pp . Therefore, the ultimate ratio of the solid described by the area QM to the solid described by the rectangle Qp , or the fluxion of the solid described by the area AMP , is to the fluxion of the cylinder described by the rectangle Ap ; as the surface described by the ordinate PM to the surface described by Pp ; and since the fluxion of the cylinder is equal to the solid made by the surface described by Pp and the fluxion, of AP , by *Art. 15*, the fluxion of the solid described by the area AMP about the tangent AT is also equal to the solid made by the surface described by the ordinate PM and the fluxion of AP .

P R O

PROBLEM. *Fig. 6.*

19. To find the fluxion of the surface described by the arc AM, of a curve about the axis AP.

Draw DM parallel to AP, meeting the tangent at A in D, and the ordinate pm produced in n: then if c expresses the circumference of a circle described by the point M, and d another such, that the rectangle made by it and the chord mM, may be equal to the surface described by that chord in the rotation of the figure about the axis AP; the cylindric surface described by nM will be to the conic surface described by the chord mM, as $c \times nM$ to $d \times mM$: and when Ap becomes AP, the circumference d becomes equal to c , and the ultimate ratio of the chord mM, or of the arc mM, by *Art. 4.* to nM, is equal to the ratio of the tangent TM to the subtangent TP. Consequently the fluxion of the cylindric surface described by DM is to the fluxion of the surface described by the arc AM. as $c \times TP$ to $c \times TM$, by *Art. 3.* or as TP to TM; that is as the fluxion of the abscissa to the fluxion of the curve AM, by *Cor. 4*; and since $c \times TP$ expresses the fluxion of the cylindric surface, by *Art. 15.* $c \times TM$, will consequently express the fluxion of the surface described by the arc AM

Having thus given the fluxions of such magnitudes as most commonly occur, we shall now investigate some others no less useful than the former.

THEOREM. *Fig. 7, 8.*

20. If in the tangent at M of a curve AM, the part MT, be taken for the fluxion of the arc AM, and a line OM be drawn from a given point O in the axis AP thro' the point M, the perpendicular TR to that line, will determine its fluxion MR.

Draw the ordinate PM, and TV, perpendicular to it, draw likewise, Vr perpendicular to TR, and Vn to OM. If $OP = x$, $PM = y$, $OM = z$, then will $MV = \dot{y}$,

$=\dot{y}$, $TV=\dot{x}$, by *Art. 4.* and by the similarity of the triangles OPM , VnM , we have $OM(z):PM(y)::MV(\dot{y}):nM$, or $z \times nM = y\dot{y}$; by the similarity of the triangles OPM , VrT , $OM(z):OP(x)::TV(\dot{x}):Vr$, or $z \times Vr = x\dot{x}$; and since $Vr \pm nM = MR$, we have $z \times MR = x\dot{x} \pm y\dot{y}$; but the right angled triangle OPM , gives $zz = xx + yy$, whose fluxion is $z\dot{z} = x\dot{x} \pm y\dot{y}$, according as y increases or decreases, while x increases, and therefore $z \times MR = z\dot{z}$, or $MR = \dot{z}$.

T H E O R E M.

21. *The fluxion of the sector OMA, is equal to $\frac{1}{2}OM \times TR$.*

The similar triangles OPM , VrT , give $z \times Tr = y\dot{x}$, and the similar triangles OPM , VnM , $z \times Vn = x\dot{y}$; and since $Vn \mp Tr = TR$, we get $z \times TR = x\dot{y} \mp y\dot{x}$; now because $y\dot{x}$ is the fluxion of the segment AMP , by *Art. 16.* and $\frac{1}{2}x\dot{y} \mp \frac{1}{2}y\dot{x}$, that of the triangle $OPM = \frac{1}{2}x\dot{y}$, by *Art. 9.* the sum or difference gives $\frac{1}{2}x\dot{y} \mp \frac{1}{2}y\dot{x}$ or its equal $\frac{1}{2}OM \times TR$, for the fluxion of the sector OAM.

C O R.

22. It is manifest that TR expresses the fluxion of the circular arc described from the center O with the radius OM , and whose sine is PM . For if z be supposed constant, the fluxion of $zz = yy + xx$, will be $y\dot{y} \pm x\dot{x} = 0$, and writing the value of $\dot{x}$ into $z \times TR = x\dot{y} \pm y\dot{x}$, we get $zx \times TR = y \times yy + xx$, or because $zz = yy + xx$, $x \times TR = z\dot{y}$; and since the cosine $OP(x)$ is to the radius $QM(z)$ as MV to MT , or as the fluxion $\dot{y}$ of the ordinate PM is to the fluxion TR of the arc, by *Art. 4, No. 2.*

C O R.

23. Hence, by drawing the chord AM , and producing the tangent MT , so as to meet the axis in Q ; if $PQ = \dot{x}$, and $PM = \dot{y}$; the fluxion of the triangle OAM, will be $\frac{1}{2}AQ \times PM$, and that of the sector OMA

OMA is $OPM \mp QPM$, or $\frac{1}{2}OQ \times PM$. Consequently the fluxion of the triangle OAM is to the fluxion of the sector OMA, as OA is to OQ; and the fluxion of the triangle OMA is to the fluxion of the segment AMA, as OA to OA—OQ or AQ.

C O R. *Fig. 9.*

24. If from any other point C in the axis OA, a line CM be drawn to the same point M, and from the same point T in the tangent at M, the lines TR, TS, perpendicular to OM and CM, will express the fluxions of the circular arcs described from centers O, C, thro' the point M, and whose sines is PM, by *Art. 22*, and MR, MS, are the fluxions of the lines OM, CM, by *Art. 20*, when MT expresses the fluxion of curve AM.

C O R. *Fig. 10.*

25. *Lastly*, If from any given point O, in the axis OA, a line be drawn meeting the lines AM, BN, in M and N, and if in the tangents at M, N, the parts MT, NV, are taken for the fluxions of AM, BN, the perpendiculars TR, VX, to ON, will be proportional to the radii OM, ON; since they express the fluxions of similar circular arcs by *Art. 22*; and consequently the fluxions of the spaces OAM, OBN, are as the squares of the lines OM, ON, by *Art. 21*.

Of various Orders of Fluxions.

When fluents are not augmented by any uniform velocity, it is convenient in many problems to consider how these velocities vary. The rate of this variation is called the fluxion of fluxion, or the second fluxion of the fluent. These second fluxions may also vary in different magnitudes of the fluent, and the rates of these variations is called the third fluxion. This may continue to as many subordinate degrees of fluxions, as the fluent admits of variations.

To give a clear and distinct notion of the different degrees

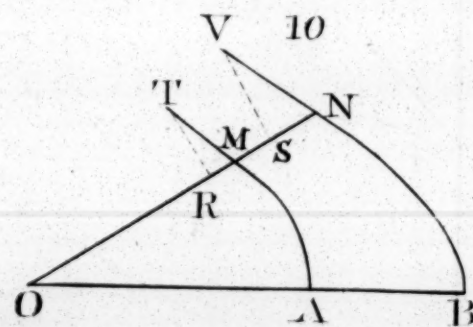
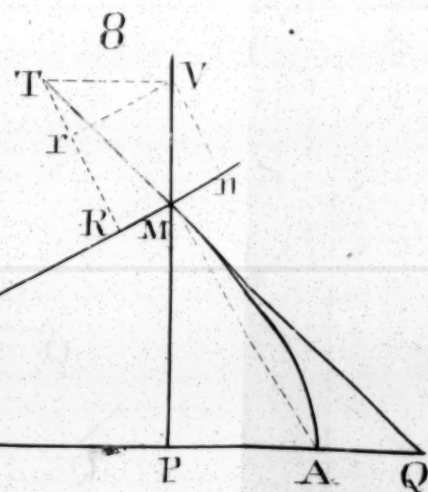
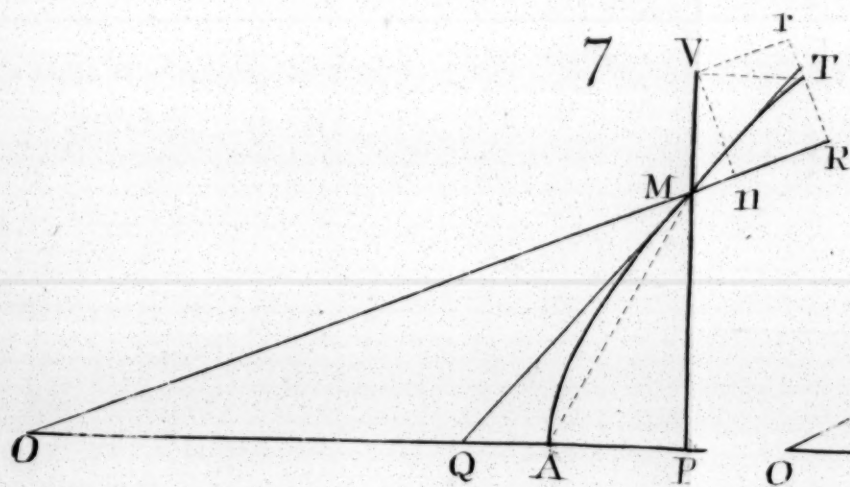
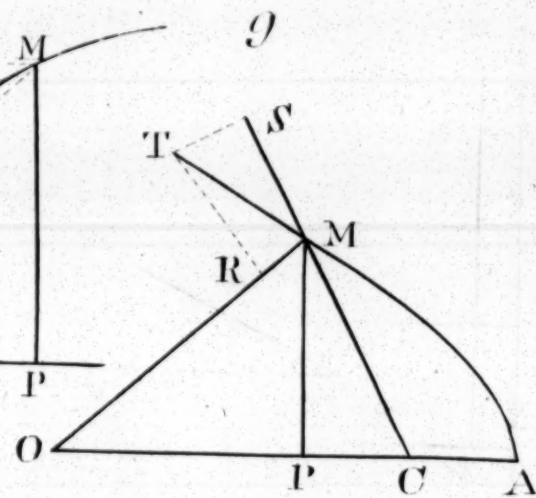
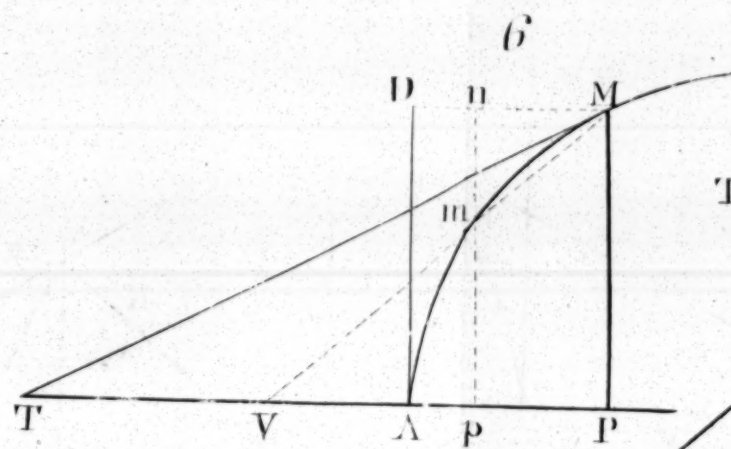
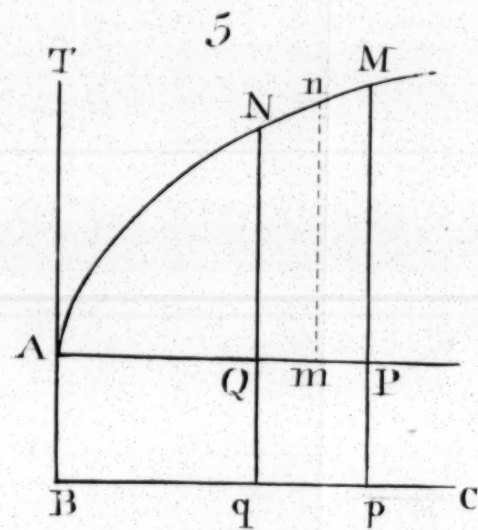
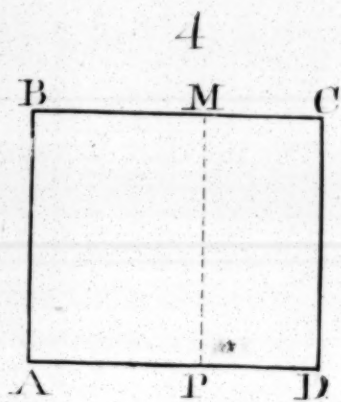
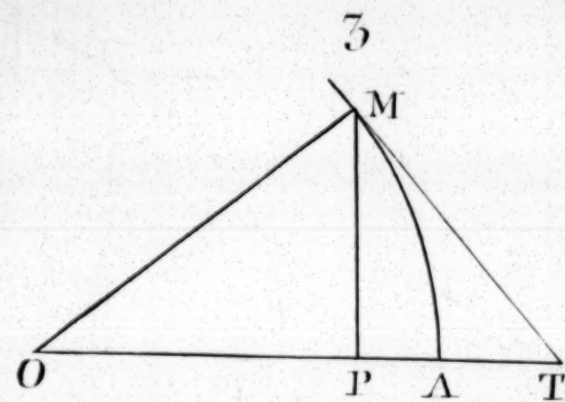
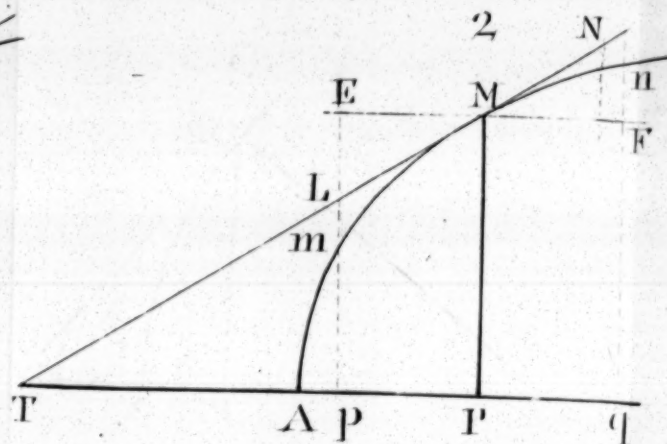
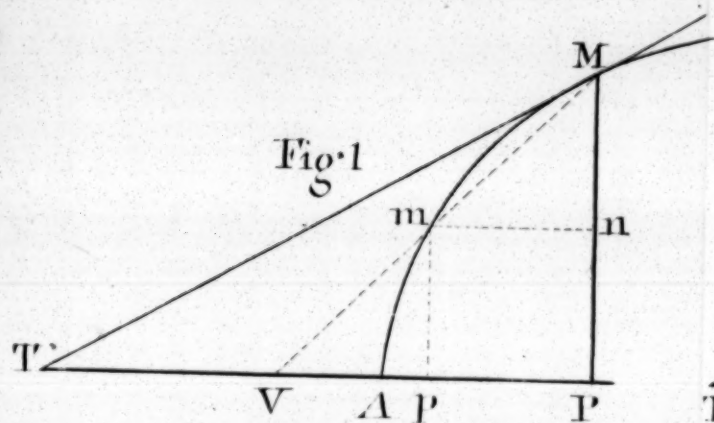
degrees of fluxions, it must be observed, that the fluxions of a curve, abscissa and ordinate, may be expressed by the three sides of a triangle formed by the tangent, subtangent and ordinate, by *Art.* 4. and since the sides of this triangle may vary in different points of the curve, the rates or these variations are called the second fluxions; and as these variations may likewise be expressed by the sides of another triangle, the fluxions of its sides are called the third fluxions.

As one side of a triangle may be constant, without changing the rate of variation; so one of the fluxions may be supposed constant, while the others vary; the ratio of any two may likewise be supposed given, provided regard is had of this supposition in future operations.

The first fluxions, algebraically expressed, being marked with a point, the second and third are likewise marked by two and three points. Thus the first fluxions of x, y, z , being marked thus, $\dot{x}, \dot{y}, \dot{z}$; the second are $\ddot{x}, \ddot{y}, \ddot{z}$; and the third $\dddot{x}, \dddot{y}, \dddot{z}$.

The second and third fluxions are found in the same manner as the first, by considering the former as variable magnitudes; thus, if $x = y^n$, denotes the equation of a curve, then will $\dot{x} = n\dot{y}y^{n-1}$, and if y be supposed constant, then will $\ddot{x} = n \times \overline{n-1} \dot{y}^2 y^{n-2}$ and $\ddot{x} = n \times \overline{n-1} \times \overline{n-2} \dot{y}^3 y^{n-3}$.

From whence, if the index n is any whole positive number, x will have as many subordinate fluxions as are units in n , and no more. For when $n = 1$; then $\dot{x} = \dot{y}$ and $\ddot{x} = 0$; when $n = 2$, $\dot{x} = 2\dot{y}y$, $\ddot{x} = 2\dot{y}^2$ and $\dddot{x} = 0$, and when $n = 3$; $\dot{x} = 3\dot{y}y^2$; $\ddot{x} = 6\dot{y}^2y$ and $\dddot{x} = 6\dot{y}^3$.



SECTION II.

The METHOD of finding the greatest and least values of variable QUANTITIES or MAGNITUDES.

PROBLEM. Plate II. Fig. 11, 12, 13, 14.

26. **T**O find the greatest and least values of a variable quantity or magnitude.

Let the ordinate PM of a curve ME, be expressed by powers of the abscissa AP and constant quantities: it is evident, that when that ordinate becomes either the greatest or least DE; the subtangent TP, becomes infinite or nothing in that case, since the tangent TM becomes parallel to the axis AP, in the first case, and coincides with the ordinate DE, in the second; the fluxion of the ordinate becomes therefore 0 in the first case and that of the abscissa in the second. From whence follows this general

R U L E.

Make the fluxion of the expression or ordinate equal to 0; then the value found of the abscissa or variable quantity will determine the value sought.

As the value of the variable quantity thus found, may give either the greatest or the least quantity, it is necessary to observe, that if any other value of the variable quantity be taken, and the quantity be always less than that found, it will be the greatest, or if they are always greater, the least of all.

E X A M P L E.

27. Let $yy = 2ax - xx$, be the equation of the curve, $AP = x$ and $PM = y$; then $0 = 2ax - 2xx$ or $x = a$; this value wrote into the equation gives $yy = aa$ or $y = a$, and y is the greatest ordinate: for when x is either 0 or $= 2a$, y becomes 0.

N. B. Observe that when y is either the greatest or the least, any of its powers, or the sum of the products

18 *The Method of finding the greatest and least* Sect. II.
ducts of y and constant quantities, are likewise the
greatest or least, as may be easily proved.

P R O B L E M.

28. *To divide a given line a into two such parts x, y ,
that $x^m y^n$ shall be the greatest of all the like products.*

The fluxion $mxy^{n-1}\dot{x} + nyx^{m-1}\dot{y}$ divided by $x^m y^n$,
and made equal to 0 gives $\frac{m\dot{x}}{x} + \frac{n\dot{y}}{y} = 0$, or $my\dot{x} + nx\dot{y} = 0$:
but $a = x + y$, by supposition, and therefore $\dot{x} + \dot{y} = 0$,
or $-\dot{x} = \dot{y}$; hence $my\dot{x} + nx\dot{y} = 0$, gives $my = nx$. Which
shews, that if the given line be divided into two parts
proportional to the indices m, n , the product $x^m y^n$, will
be the greatest of all.

For since $y = a - x$, the product $x^m y^n$ becomes $x^m \times$
 $\overbrace{a - x}^n$; now when x is either 0 or $= a$, that expression
vanishes or is $= 0$; the value found is therefore the
greatest.

Because $my = nx$ and $y = a - x$, we get $ma - mx = nx$,
or $\frac{ma}{m+n} = x$: now if n be negative and less than m ;

then $\frac{ma}{m-n} = x$, and since m is greater than $m-n$, x
must be greater than a ; and this cannot be unless we
suppose $y = a + x$, which makes $x^m y^n$, or its equal $x^m \times$
 $\overbrace{x - a}^{-n}$, to be the least of all, when $my = nx$, because
if $x = a$, or x infinite, the expression is infinite.

P R O B L E M.

29. *To divide a given line a into three parts x, y, z ,
such that the product $x^m y^n z^r$ shall be the greatest of all.*

It is evident, that when z is constant, and $x^m y^n$ the
greatest, the product $x^m y^n z^r$ will also be the greatest;
and when x is constant, and $y^n z^r$ the greatest, the pro-
duct $x^m y^n z^r$, will be likewise the greatest; but when x
is constant we get $ry = nz$, by the last, and when z is
constant $my = nx$. From these two equalities we get m :

$$x :: n :$$

$x::n:y::r:z$; which shews that if the line a be divided into three parts proportional to the indices m, n, r ; the product above will be the greatest of all. For $y=a-x-z$, by supposition, the product $x^m y^n z^r$, or its equal $x^m z^r \times \overbrace{a-x-z}^n$ becomes $=0$, when x is either 0 or $=a-z$; this product is consequently the greatest of all, when $ry=nz$ and $my=nx$.

It may be proved in the same manner, that if a line be divided into ever so many parts proportional to the indices of the powers, then the product of these powers will be the greatest of all. Thus the square of half a line is greater than any rectangle that can be made by any two unequal parts, and the cube of one third of a line greater than the parallelepiped of any three other parts.

N. B. When the variable quantity contained in the equation found by making the fluxion of the expression equal to 0 , has several positive values, the question contains more than one greatest or least, and they are alternately the greatest and least: that is, if the first is a greatest, the second will be the least, and the third the greatest, and so on; or if the first is the least, the second will be the greatest: but the least between the two greatest may sometimes be 0 , and the greatest between two least infinite.

T H E O R E M. *Fig. 15.*

30. *Of all the triangles which have equal bases and equal angles at the vertex, the isocles contains the greatest area.*

Upon the given base AC as a chord, describe the arc ABC of a circle, so as to contain the given angle: then because the line BE , which bisects the chord AC at rightangles, passes thro' the center of the circle, and is therefore the greatest that can be drawn from any point in the arc ABC , by the property of the circle; and triangles which have equal bases are as their altitudes; the isocles triangle ABC is greater than any other triangle of the same base, and the same angle at the vertex.

C 2

31. Hence,

31. Hence, of all rectilineal figures, terminated by the same number of equal sides, that which is inscribed in a circle contains the greatest area. For if the radii are drawn to all the angles, it will be divided into as many isocles triangles as there are sides, and each of which being greater than any other of the same base and angle at the vertex; their sum or the whole figure inscribed in a circle will consequently be greater than any other of the same number of equal sides.

C O R.

32. It is manifest, that of all figures whatever, which have the same boundary, the circle is the greatest, and of all equal figures, the circle has the least boundary. For since circles are always greater than any inscribed figure, and the inscribed figures being greater than any other of the same boundary; the circle will be greater than any other figure of the same boundary, or of all equal figures has the least boundary.

P R O B L E M. Fig. 16.

33. Let a squared timber BD, placed horizontally, be supported by a prop at one end D, and by a strut AB on the other B; to find the position of the brace AC, of a given length, so as to support the timber BD in the best manner possible.

Draw AF perpendicular to BD and BE to AC: then from a known principle in mechanics, if AC expresses the force with which the timber BD is supported in a perpendicular direction, AF will express the force by which it is supported in this oblique direction; therefore the rectangle made by that force and the distance BC of the point C from the point fixed B, expresses the momentum of that force, which should be the greatest possible. Now the rightangled similar triangles AFC, BEC, give $AC : AF :: BC : BE$, or $AF \times BC = AC \times BE$: whence the momentum of this force is proportional to the area of the triangle

ABC,

ABC, whose base is given as well as the angle B at the vertex ; which is the greatest, when $AB=BC$, by *Art.* 30.

C O R. *Fig.* 17.

34. If the squared timber BD be fixed at one end B, and supported by a prop AC, of a given length, it is plain that when AC makes equal angles with the timber AD, and the line BA, it will support the timber AD in the best manner. For if BE be perpendicular to AC, the product $BE \times AC$, expresses the momentum of the force with which the prop supports the timber, which is the greatest possible when $BA=BC$, by *Art.* 30.

C O R.

35. Since the length of the prop AC, has been taken at pleasure, in the two last articles, it may have one more advantageous than any other, which is, when it supports the timber in the middle, supposing the timber fixed at both ends.

P R O B L E M. *Fig.* 18.

36. To find the best angle of a roof in respect to the strength of the principal rafters AD, BD.

Let DC be perpendicular to AB, and AE to BD produced : then as the strength of the rafters decreases as their length increases, and their weights increases as their length increases, their strength decreases likewise on that account ; that is, the strength of rafters decreases in a duplicate ratio, to that of their length. Now as the absolute force of the rafter AD, is to that part by which it acts in the oblique direction, as AD is to AE, the momentum of the strength of the rafter AD is therefore as AE directly, and the square of AD or BD inversely.

By the similarity of the triangles AEB, DCB, we have $BD : BC :: BA : BE$; and since BC as well BA, are given, BE is as BD inversely : whence $AE \times BE^2$, must be the greatest possible. If $AB=a$, $AE=x$; then

22 *The Method of finding the greatest, &c.* Sect. II:
 will $BE^2 = aa - xx$, and $aax - x^3$, the greatest, whose
 fluxion made $= 0$, gives $aa = 3xx$. For when $x = a$ or
 $= 0$, the expression becomes 0. By the similarity of
 triangles we get $CD^2 = \frac{1}{3} aa$, and by extracting the
 square root $CD = \frac{\sqrt{3}}{3} a$ nearly. This answers nearly
 the practical rule used by carpenters, of making the
 height CD one-third of the base AB.

P R O B L E M. *Fig. 19.*

37. *To find a point D, within a given triangle, such
 that the sum of the lines drawn from this point to the three
 angles, may be the least possible.*

Suppose one of these lines AD, as invariable, from
 the center A with the radius AD describe an arc of a
 circle, and from any point a in the tangent at D,
 draw a b perpendicular to BD and a c to CD produ-
 ced: then will Db, be the fluxion of DB, and Dc,
 that of CD, by *Art. 24*: and since the sum of these
 lines is given their fluxions must be equal, and as Da
 is common to the right-angled triangles Dca, Dba,
 they are equal in all respects: whence by adding the
 angle aDb to the right angle ADa, we get the angle
 ADB, and by adding aDc or its equal dDC to the
 right angle ADD, we get the angle ADC equal to
 ADB. It may be proved in the same manner, by
 supposing CD constant, that the angles CDA, CDB,
 are equal consequently, when the three angles at the
 point D, are equal, the sum of the three lines AD, BD,
 CD, will be the least possible.

P R O B L E M. *Fig. 20.*

38. *To find a point within a quadrilateral ABCD,
 so that the sum of the four lines drawn from it to the
 angles, shall be the least possible.*

The intersection E of the two diagonals AC, BD
 will be the point required. For if from any other
 point F, lines are drawn to the opposite angles A, C,
 the

the sum of the lines AF, CF, will always be greater than the diagonal AC: and for the same reason, the sum of the lines drawn from the point F to the opposite angles B, D, will always be greater than the diagonal BD, the intersection E of the diagonals is therefore the point required.

What has been said on this subject in our elements of mathematics, and here together, with what will be said more hereafter, will sufficiently explain this very useful method, in all mechanical subjects.

SECTION III.

Of the RADII of CURVATURE.

WHEN a circle touches a curve so as to have the same fluxion in that point, the radius of this circle is called the *Radius of Curvature* in that point.

PROBLEM. Fig. 21.

39. To find the radius of curvature DM, by means of the lines drawn from the given point C.

From any point T in the tangent at M, and the extremity D of the radius of curvature, draw TR and DE, perpendicular to CM, and from the point C, the line CS to MT: let $DM=r$, $CM=y$, $CS=s$, and $MT=\dot{z}$, $MR=\dot{y}$, $TR=\dot{x}$; then the angular motion of the sides DM, MS, of the right angle DMS, at M, being always equal; the radius DM will be to the radius MS, as the fluxion $\dot{z}$ of the circular arc described by the point M is the fluxion $\dot{s}$ of the circular arc described by the point S in the same time; that is $DM:MS::\dot{z}:\dot{s}$; and the similar triangles TRM, CSM, give $MS:CM::MR(\dot{y}):TM(\dot{z})$; the compound of these two proportions give $r:y::\dot{y}:\dot{s}$, or $r\dot{s}=y\dot{y}$. But $CM(y):CS(s)::DM(r):ME=\frac{rs}{y}$. By writing the value of r , in $r\dot{s}=y\dot{y}$, into that of ME, we get $\dot{s}\times ME=s\dot{y}$

EXAMPLE. *Fig. 22.*

40. Let AM be a common parabola, A its vertex, and C the focus; then by the property of that curve, the tangent at M meets the tangent at A, in the same point S, as the perpendicular drawn from the focus to the tangent MS; the right angled triangles CAS, CSM, are therefore similar; that is, if $CA = a$, we get $a:s::s:y$, or $ay = ss$, whole fluxion is $ay = 2sy$, and as $r\dot{s} = yy$, we get $ar = 2sy$.

CONSTRUCTION.

Draw ME parallel to the axis AC, and equal to twice CM; then the perpendicular ED to ME determines the radius required. For the similar triangles SAC, DEM, give AC (a):CS (s): : ME ($2y$):MD (r) or $ar = 2sy$.

When the point M coincides with the vertex A of the axis, it is plain that CM, CS become equal to CA, that is $s = y = a$, so then $ar = 2sy$, gives $r = 2a$, which shews that the radius of curvature is equal to half the parameter of the axis at the vertex.

EXAMPLE. *Fig. 23.*

41. Let AM be an ellipsis or hyperbola, C the center and Nn the conjugate to the diameter Mm: if $CN = u$, a, b , half the first and second axis, then will $us = ab$, by the properties of these curves, whose fluxion is $-s\dot{u} = u\dot{s}$; and since the sum or difference between the squares of any two conjugate diameters is always constant, that is $yy \pm uu$ is constant, we $\mp yy = u\dot{u}$, the compound of this and $-s\dot{u} = u\dot{s}$ gives $\mp sy\dot{y} = uu\dot{s}$; this and $\dot{s} \times ME = sy$, by *Art. 39.* gives $y \times ME = \pm uu$.

CONSTRUCTION.

In CM take ME towards the center in the ellipsis or from the center in the hyperbola, so that $\pm CM : CN : ME$, then ED perpendicular to ME, will determine the radius required. When

When the point, M coincides with the vertex A or B, the radius of curvature becomes equal to the parameter of half the first or second axis.

OTHERWISE.

42. Because $r\dot{s} = y\dot{y}$, by Art. 39, and $us = ab$, by the property of the curve, or $s = \frac{ab}{u}$, whose fluxion is $-\dot{s} = \frac{ab\dot{u}}{uu}$, this and $r\dot{s} = y\dot{y}$ gives $-\frac{rab\dot{u}}{uu} = y\dot{y}$; but $y\dot{y} \pm u\dot{u} = 0$, therefore $rab = \pm u^3$, by substitution. Now if MD intersects the first axis AC, in K and $MK = c$; then will $ac = bu$, by the nature of the curve. Hence $rab = \pm u^3$ becomes $rb^4 = aac^3$, or if $ap = bb$, we get $rpp = c^3$. Consequently the radius MD, is a fourth continued proportional to the parameter of half the first axis AC, and to the perpendicular MK.

No. 2. Since (Fig. 21) $TM(\ddot{z}) : TR(\ddot{x}) :: CM(y) : CS(s)$ or $s\ddot{z} = y\ddot{x}$, and if we suppose the point C to remove at an infinite distance from the point M, so that the ordinates CM may become parallel and equal, then y being infinite, may be considered as constant the fluxion of $s\ddot{z} = y\ddot{x}$, compared with $r\dot{s} = y\dot{y}$, gives $r\ddot{x} = y\ddot{z}$, or $-y\ddot{z} = r\ddot{x}$ according as $\dot{z}$ or $\dot{x}$ is constant, for the equations, when the ordinates are parallel.

There are different ways of finding the radii of curvatures, but none more easy and so short as this, nor that leads more directly to the application; the many examples of curves that may be described, and the finding the nature of the curve which is continually touched by the radii of curvature, are no more than a useless speculation; we shall therefore take no notice of them here, excepting the following one useful in the theory of pendulums.

DEFINITION. Fig. 24. Plate III.

If a semi-circle ANB, rolls upon a right line BD, the

the curve AMD described by the point A in the circumference, is called a *Cycloid*.

C O R.

43. Hence, when the describing point A, arrives to any other point M; the chord ML, drawn from the point M to the point of contact L, will from the nature of the description be perpendicular to the tangent TM of the curve at M, and this tangent will be the chord of the complement of LM to the semi-circle. It appears also that the base BD is equal to the semi-circumference ANB of the generating circle, the part BL to the arc TM, and the part LD to the arc ML.

P R O B L E M.

44. *To find the length of any arc AM of the cycloid.*

Draw PM perpendicular to the diameter TL, which itself is perpendicular to the base BD, and call LT or AB, a , $TP = x$, the chord $TM = v$, and the arc $AM = z$: then because $TP : TM$ or $TM (v) : LT (a) :: \dot{x} : \dot{z}$, by *Art. 4.* we have $v\dot{z} = a\dot{x}$; and since $LT (a) : TM (v) :: TM (v) : TP (x)$ or $ax = vv$, by the similarity of triangles, whose fluxion is $a\dot{x} = 2v\dot{v}$, whence $(a\dot{x} =) v\dot{z} = 2v\dot{v}$, or $\dot{z} = 2\dot{v}$: and as the fluxion $\dot{z}$ of the arc is always double the fluxion $\dot{v}$ of the chord TM; the arc AM will likewise be double the chord TM, by *Art. 7.* the whole curve AMD is therefore double the diameter AB of the generating circle.

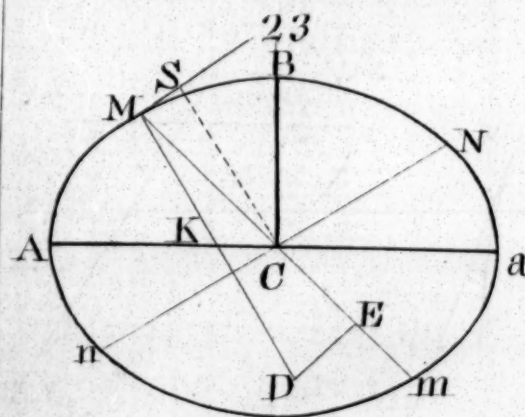
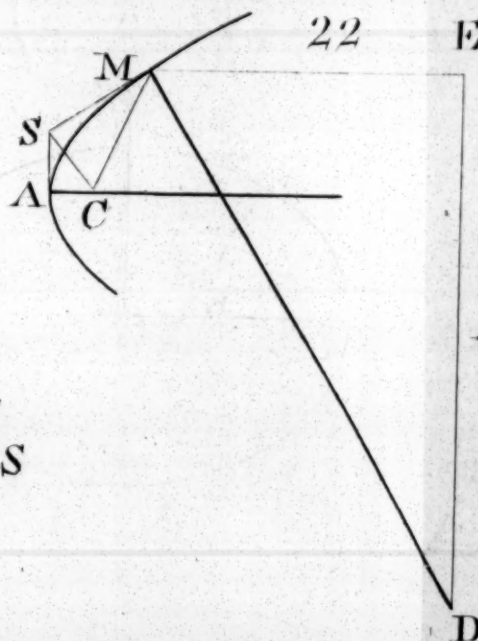
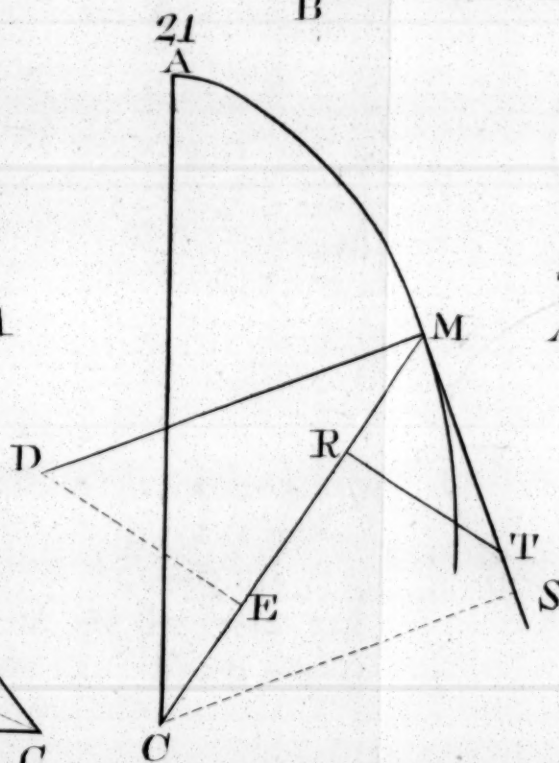
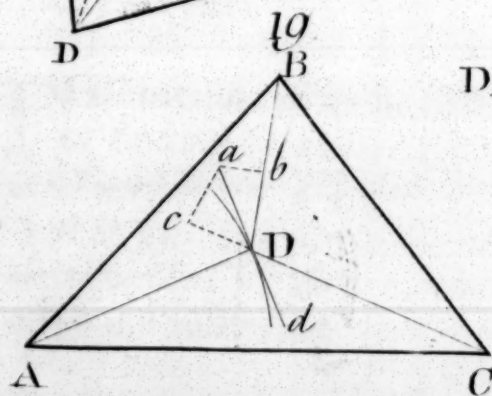
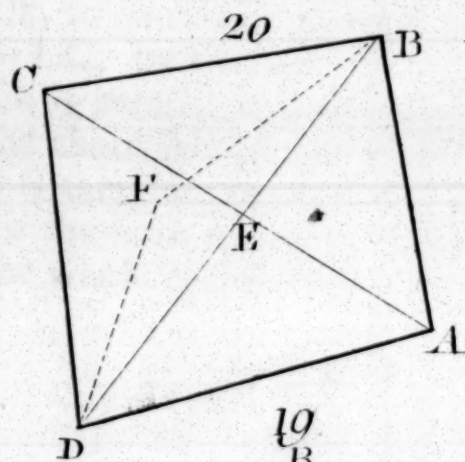
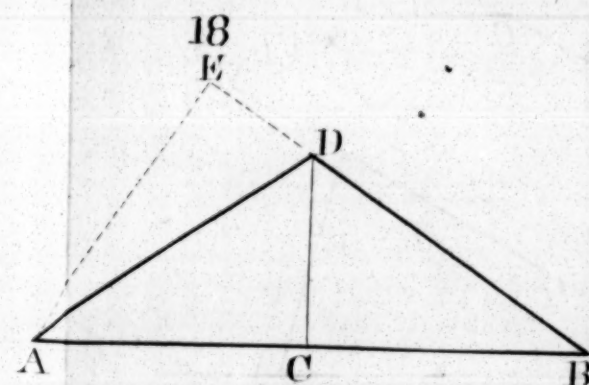
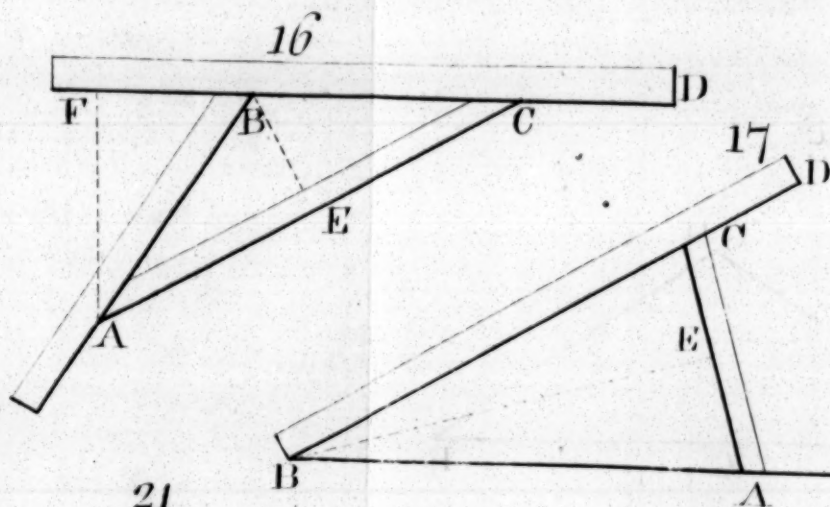
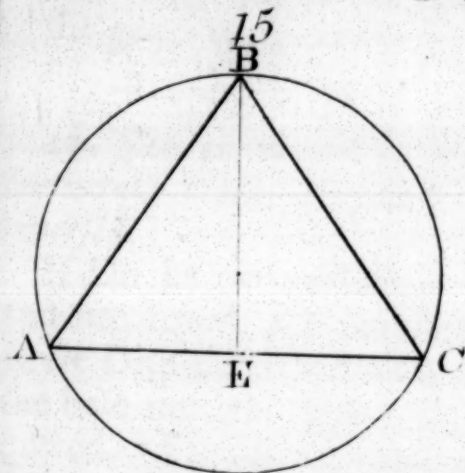
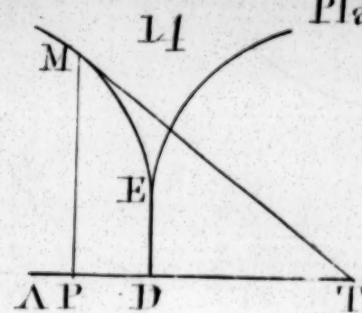
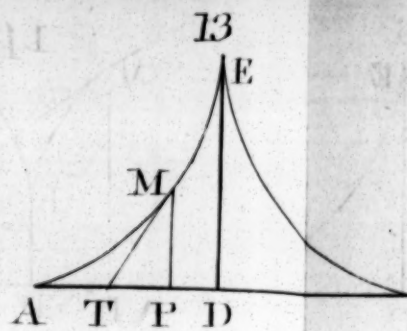
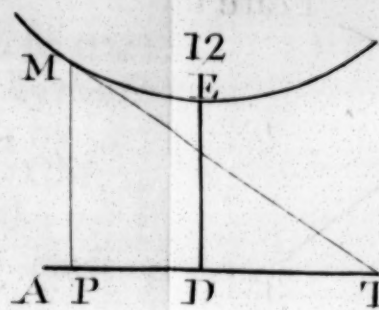
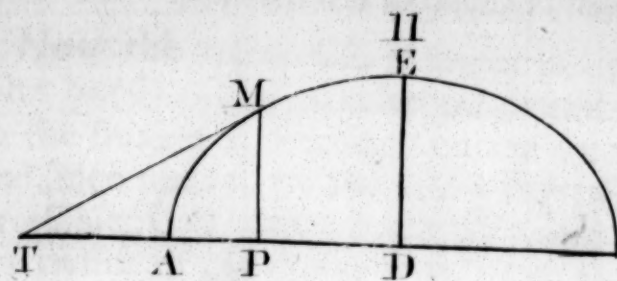
P R O B L E M.

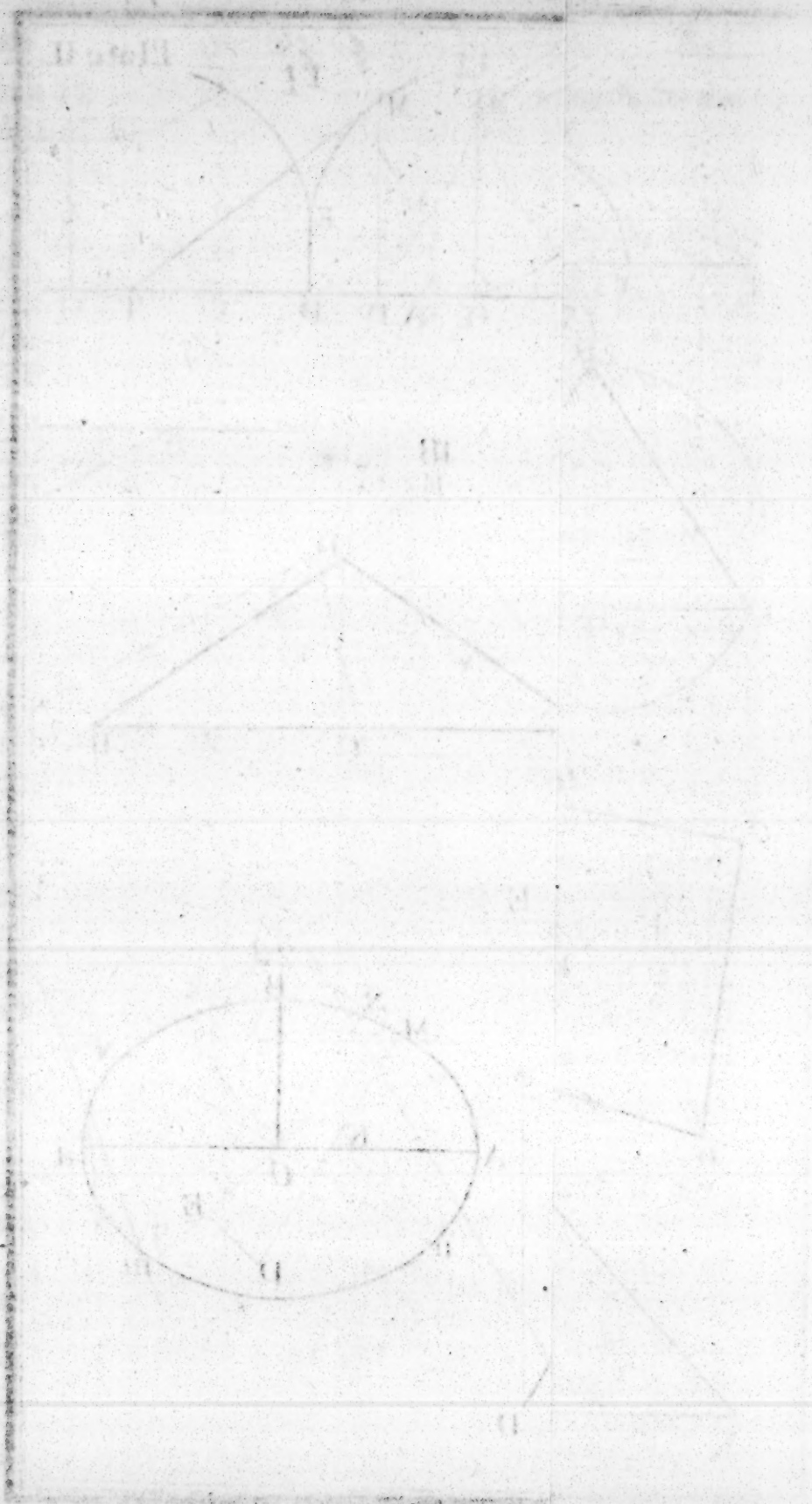
45. *To find the radius of curvature CM at any given point in the cycloid.*

In the base BD, take $bL = y$ for the fluxion of BL or its equal the arc TM, and draw ba perpendicular to CM: then the similar triangles TLM, bLa, give $TL : LM :: bL (y) : ba$: but the radius TL is to the sine LM of the angle LTM, as the fluxion y of the arc is to the fluxion $\dot{v}$ of the cosine TM, by *Art. 4.*

No. 2. Hence $y : ba :: y : \dot{v} = ba$, by equality of ratios:

Now





Now the radius CL is to the radius CM as the fluxion $ba(\dot{v})$ of the circular arc described by the first, is to the fluxion $\dot{z}$ described by the second, by *Art.* 25. and since $\dot{z} = 2\dot{v}$, by the last we get $CL : CM :: \dot{v} : 2\dot{v}$, or $2CL = CM$, and $CL = LM$; and at the vertex A, the radius of curvature is equal to twice the diameter of the generating circle.

THEOREM.

46. *The curve mCD which the radii of curvature continually touch, is another cycloid equal to the first, but reversed.*

If Dn be perpendicular to BD, and equal to AB, and nm equal and parallel to DB, then if the semi-circle Dkn described with the diameter nD, rolls upon the base nm, the point D will describe the curve required. For if the chord Dk be drawn parallel to MC, and Ck joined; then the parallel chords ML, Dk, in equal circles are equal, and since ML is equal to LC, by the last, CD is a parallelogram, and so Ck is parallel and equal to LD. But MC is double the chord ML, or equal to the arc DC of the curve, it is therefore equal to the arc AM, by *Art.* 44; and for the same reason the arcs DM, mC are equal, and the whole cycloid AMD to the whole cycloid DCm.

SECTION IV.

Of FLUENTS.

THE method of finding fluents, is the inverse of that of finding fluxions. For the fluent of ny^{n-1} is y^n , that of $ncr\dot{z}z^{n-1} \times \overline{a + cz^{n-1}}$, is $\overline{a + cz^n}$; and that of $\dot{y}z + y\dot{z}$ is yz ; for these fluents are the very same as those proposed in *Art.* 9, 10, 14. From whence we deduce the following

RULE I.

47. *To find the fluent of a single quantity $\dot{y}y^n$, raised to any power.* Change

Change the fluxion $\dot{y}$ into a fluent, and divide by the index $n+1$ thus increased.

R U L E II.

48. To find the fluent of a binomial $\dot{z}z^{n-1} \times \overline{a+cz^n}$, whose fluxion is in a constant ratio to that which multiplies the binomial

Increase the index n of the binomial by unity, and divide by the product of the index thus increased and the fluxion of the binomial.

E X A M P L E S.

49. The fluent of $\dot{a}x + \dot{b}xx + \dot{c}xx^2 + \dot{d}xx^3$, is $ax + \frac{1}{2}bx^2 + \frac{1}{3}cx^3 + \frac{1}{4}dx^4$, by the first rule.

The fluent of $xx \times \overline{aa+xx}^{\frac{1}{2}}$, is found by adding unity to the index $\frac{1}{2}$, which gives $\frac{1}{2}+1$, or $\frac{3}{2}$, this multiplied by the fluxion $2xx$ of the binomial and divided by the product $3xx$, gives $\frac{1}{3} \times \overline{aa+xx}^{\frac{3}{2}}$, for the fluent required by rule the second.

If $xx \times \overline{aa+xx}^{-\frac{1}{2}}$; then unity added to the index $-\frac{1}{2}$ gives $1-\frac{1}{2}$ or $\frac{1}{2}$, this multiplied by $2xx$ the fluxion of the binomial $aa+xx$, and the fluxion $xx \times \overline{aa+xx}^{\frac{1}{2}}$ divided by the product xx gives $\overline{aa+xx}^{\frac{1}{2}}$, or $\sqrt{aa+xx}$ for the fluent required.

N. B. When a fluent has a constant term, it is generally incompleat, because the constant terms vanish in the fluxion, but may be compleated; for if the fluent should become 0, when the variable quantity is of a certain value, what remains in this case, must be added with a contrary sign. Thus if $\sqrt{aa-xx}$, should become 0, when $x=0$; then because when $x=0$, there remains $\sqrt{aa}$ or a : we get $\sqrt{aa-xx}-a$ for the compleat fluent. But if $\sqrt{aa-xx}$ should vanish when $x=a$, it requires no addition.

P R O-

P R O B L E M. Fig. 25.

50. To express the circular arc AM by the radius CA (r) and the tangent AT (z).

The square $rr + zz$ of CT, being to the square rr of the radius CM, by Art. 25, as the fluxion $\frac{1}{2}r\dot{z}$ of the triangle CAT, to the fluxion $\frac{\frac{1}{2}\dot{z}r^3}{rr + zz}$ of the sector

CAM, which divided by $\frac{1}{2}r$ gives $\frac{r\dot{z}}{rr + zz}$ for the flu-

xion of the arc AM, by Art. 21. Now this fluxion reduced to an infinite series by a continual division, supposing $r=1$, gives $\dot{z} - \dot{z}z^2 + \dot{z}z^4 - \dot{z}z^6 +$, &c. whose fluent, by the first rule, is $z - \frac{1}{3}z^3 + \frac{1}{5}z^5 - \frac{1}{7}z^7 + \frac{1}{9}z^9 -$, &c. Or if $A=z$, $B=zzA$, $C=zzB$, $D=zzC$, &c. This fluent becomes $A - \frac{1}{3}B + \frac{1}{5}C - \frac{1}{7}D + \frac{1}{9}E -$, &c.

As this series converges but very slowly unless z is very small, we shall give the following

L E M M A. Fig. 26.

51. The tangents $AT=a$, and $AB=b$, of any two arcs AF, AM, being given, to find the tangent $FG=t$, of their difference FM.

Draw GD perpendicular to CA, intersecting CB in E, make $CA=1$, and $CB=x$: then the similar triangles CAB, GFE, give $CA:AB::GF:EF=bt$, or, $CE=1-bt$, and $CA:CB::GF:GE=tx$, and by the parallels GD, TA, we have $CE:CB::GE:TB$, that is, $1-bt:x::tx:a-b$, or $txx=a-b \times 1-bt$, or because $xx=1+bb$, we get $t+bbt=a-b-abt+bbt$, and therefore $t=\frac{a-b}{1+ab}$, by reduction.

C O R. Fig. 27.

52. Hence, if from any arc AE, another arc AB be taken as often as can be done, and the point F be nearest to the point E; by calling $AE=A$, $AB=B$,
BE

$BE=C$, $CE=D$, $DE=E$, $FE=F$, and their respective tangents, a, b, c, d, e ; and supposing $a=1$, $b=\frac{1}{3}$, the other tangents are as follows,

$$A-B=C, c=\frac{2}{3}.$$

$$C-B=D, d=\frac{7}{17}.$$

$$D-B=E, e=\frac{9}{48}.$$

$$B-E=F, f=\frac{1}{239}.$$

Now, if the three first equalities of the arcs are added together, then will $A-3B=E$, and this equality subtracted from the fourth, gives $4B-F=A$; that is, four times the arc B , whose tangent is $\frac{1}{3}$, less the arc F whose tangent is $\frac{1}{239}$, is equal to an arc of 45 degrees.

In the same manner other tangents may be found, so as their sums or difference shall be equal to a given arc. Thus if $c=\frac{1}{7}$, $d=\frac{3}{79}$; then will $2D+5C=45$ degrees, or if $a=\frac{1}{12}\sqrt{3}$, and $b=\frac{1}{37}\sqrt{3}$, then $3A+2B=30$ degrees.

EXAMPLE.

53. Let $x=\frac{1}{239}$, then by reducing the series $A-\frac{1}{3}B+\frac{1}{3}C-\frac{1}{7}D+\frac{1}{9}E$, &c. into decimals, and setting all the terms of the same sign under each other, we get,

$$+0.00418,41004,18410,04184,101=A.$$

$$256,47231,442=\frac{1}{3}C:$$

$$4=\frac{1}{9}D.$$

$$+0.00418,41004,18665,51415,574.$$

$$-0.00000,00244,16591,78708,380=\frac{1}{3}B.$$

$$320,713=\frac{1}{7}D:$$

$$-0.00000,00244,16591,79029,093.$$

Hence we get $0.00418,40760,02074,72386,454$
= the arc F .

And

And if we suppose $x = \frac{1}{3} = .2$; the series above gives

$$+0.20006,40568,88888,88888,889 = A + \frac{1}{3}C;$$

$$63015,38461,538 = \frac{1}{9}E.$$

$$77,10117,647 = \frac{1}{3}G.$$

$$9986,438 = \frac{1}{17}I.$$

$$13,422 = \frac{1}{31}L.$$

$$18 = \frac{1}{25}N.$$

$$+0.20006,40569,51981,47467,952.$$

$$-0.00266,66666,66666,66666,666 = \frac{1}{3}B.$$

$$18285,71428,57142,857 = \frac{1}{7}D.$$

$$18,61816,18181,818 = \frac{1}{11}F.$$

$$2184,33333,333 = \frac{1}{15}H.$$

$$2,75941,053 = \frac{1}{19}K.$$

$$364,722 = \frac{1}{23}M.$$

$$498 = \frac{1}{27}O.$$

$$-0.00266,84971,02100,71630,947.$$

Hence $0.19739,55598,49880,75837,005 =$ to the arc B.

Now because four times the arc B, less the arc F, is equal to 45 degrees, 16 times the arc B less 4 times the arc F, will be equal to half the circumference, which therefore is $3.14159,26535,89793,23846,264$, true in all its places:

C O R.

54. Since unity is to 3.14159 , &c. as the radius a of any circle is to $a \times 3.14159$ the semi-circumference of that radius; and if D expresses any number of degrees of that circle; then 180 degrees will be to D degrees as the semi-circumference $a \times 3.14159$ is to $.01745329 \times aD$, that arc expressed by parts of the radius. Consequently the number of degrees being given, the arc itself is given.

C O R.

55. If the diameter of a circle be 7, to find its circumference, say, 1 is to 3.14159 as 7 is to 21.99 , or

22 nearly, which is Archimedes's proportion; or if the diameter be 113, then its circumference will be 354.99967 or 355 exceeding near, which is that of Snellius.

P R O B L E M. Fig. 28.

56. To find the area of an equilateral hyperbolic space BCMP, terminated by two perpendiculars to the asymptote AP.

Let $AB=1$, $BC=a$, and $BP=x$, then will $PM=\frac{a}{1+x}$, by the nature of the curve, and $\frac{ax}{1+x}$ the fluxion of the space BM, by Art. 16. which being reduced into an infinite series by a continual division, and the fluent found by Art. 47, gives $a \times$ by $x - \frac{1}{2}x^2 + \frac{1}{3}x^3 - \frac{1}{4}x^4 + \frac{1}{5}x^5 -$, &c.

But if $Bp=x$, then will $pm=\frac{a}{1-x}$, and $\frac{-xa}{1-x}$ the fluxion of the space pC, whose fluent being found as before is $-a \times$ by $x + \frac{1}{2}x^2 + \frac{1}{3}x^3 + \frac{1}{4}x^4 + \frac{1}{5}x^5 +$, &c.

C O R.

57. Since $\frac{xa}{1+x}$ is the fluxion of the hyperbolic space BM or of the logarithm of AP, and $\frac{-xa}{1-x}$ that of Ap, by the nature of the curve; it is manifest, that the fluxion of the logarithm of any number, is equal to the fluxion of that number divided by the same number.

C O R.

58. From the nature of logarithms, the difference between the logarithms $a \times$ by $x - \frac{1}{2}x^2 + \frac{1}{3}x^3 -$, &c. and $-a \times$ by $x + \frac{1}{2}x^2 + \frac{1}{3}x^3 +$, &c. which is, $2a \times$ by $x + \frac{1}{3}x^3 + \frac{1}{5}x^5 + \frac{1}{7}x^7 +$, &c. will be that of their quotient $\frac{1+x}{1-x}$. Or if $A=1$, $B=xxA$, $C=xxxB$, $D=xxx C$, and so on, this difference will be $2ax \times$ by $A + \frac{1}{3}B + \frac{1}{5}C + \frac{1}{7}D +$, &c.

E X-

E X A M P L E.

Let $\frac{1+x}{1-x} = 2$, then will $x = \frac{1}{3}$, and this value of x , being wrote into the last series, when reduced into decimals, gives

$$\begin{array}{r}
 1.00000,00000 \\
 3703,70370 \\
 246,91358 \\
 19,59632 \\
 1,69351 \\
 15396 \\
 1447 \\
 139 \\
 15 \\
 \hline
 1.03972,07708
 \end{array}$$

This multiplied by $2ax$ or $\frac{2a}{3}$, gives $a \times .69314, 71805$, for the logarithm of 2.

E X A M P L E.

Let $\frac{1+x}{1-x} = \frac{5}{4}$; then will $x = \frac{1}{9}$, and this value wrote into the series gives

$$\begin{array}{r}
 1.00000,00000 \\
 - 411,52263 \\
 3,04832 \\
 2688 \\
 26 \\
 \hline
 1.00414,59809
 \end{array}$$

This multiplied by $2ax$ or $\frac{2a}{9}$, gives $a \times .22314, 35513$ for the logarithm of $\frac{5}{4}$. And as this number multiplied by 4 gives 5, by adding twice the logarithm of 2, we get $a \times$ by 1.60943,79124, for the logarithm

garithm of 5 : and if the logarithm of 2 be added to that of 5, we shall have $a \times$ by 2.30258,50929 for the logarithm of 10.

C O R.

59. It is plain, that there may be as many sorts of logarithms as there are different values taken for a ; when $a=1$, they are called hyperbolic. When any particular system of logarithms is given, the value of a may be found; thus in the common tables the logarithm of 10 is unity of 100 or 1000; its is 2 or 3: we have then no more to do than to make the logarithm of 10 found above equal to unity, or to divide unity by 2.3025, &c. which gives .43429,4482 for the value of a in that system.

P R O B L E M. Fig. 29.

60. To express the area of the hyperbolic sector CAM, by parts of the tangent At.

Let CN be an asymptote, and CM intersect the tangent AT at A in t. If $CA=a$, $AT=b$, and $At=z$; then the square of Ct is to the square of CM as the fluxion $\frac{1}{2}a\dot{z}$ of the triangle CAt is to the fluxion of the sector CAM, by Art. 25; but, by the nature of the hyperbola, the square of Ct is to the square of CM, as the difference $bb-zz$, between the squares of the tangents AT and At, to the square bb of AT:

therefore $bb-zz : bb :: \frac{1}{2}a\dot{z} : \frac{abb\dot{z}}{2bb-2zz} =$ to the fluxion of the sector CAM, whose fluent is the hyperbolic logarithm of $\frac{b+z}{b-z}$, multiplied by $\frac{1}{4}ab$. For the fluxion of the logarithm $b+z$ is $\frac{\dot{z}}{b+z}$, by Art. 57, and that of $b-z$, $\frac{-\dot{z}}{b-z}$, by the same: This last subtracted from the first, gives $\frac{2b\dot{z}}{bb-zz}$, which multiplied by $\frac{1}{4}ab$, gives the fluxion proposed.

E X.

E X A M P L E.

Let $b=2z$; then $\frac{b+z}{b-z} = 3$, whose tabular logarithm is .4771213, which multiplied by 2.302585, *Art.* 58, gives 1.0986123, for the hyperbolic logarithm of that number; and this multiplied by $\frac{1}{4} ab$, gives .274653 $\times ab$, for the area of the sector CAM.

P R O B L E M.

61. To express the same area CAM, by parts of the ordinate PM to the first axis CA.

Let the ordinate PM produced, meet the asymptote in N, and the tangent at M, the axis in V; if $PM = y$ and the rest as before, then CA will be to CV as the fluxion $\frac{1}{2} ay$ of the triangle CAM, is to the fluxion of the sector CAM, by *Art.* 23, but $CP:CA::CA:CV$, by the property of the curve, as also $CP:CA::PN:AT$; and the difference between the squares of PN, and PM, is equal to the square of AT, or $bb+yy=PN^2$; therefore by equality of ratios $\sqrt{bb+yy}:b::\frac{1}{2} ay:\frac{aby}{2\sqrt{bb+yy}}$ equal to the fluxion of the sector. If $z=\sqrt{bb+yy}$, the hyperbolic logarithm of $\frac{y+z}{b}$, multiplied by $\frac{1}{2} ab$, will be the fluent.

For the fluxion of the logarithm $\frac{y+z}{b}$ is $\frac{\dot{y}+\dot{z}}{y+z}$ by *Art.* 57; and the fluxion of $zz=bb+yy$ is $z\dot{z}=y\dot{y}$: hence $\dot{y}+\dot{z}=\dot{y}+\frac{y\dot{y}}{z}$ or $y+\overline{z}\times\frac{\dot{y}}{z}$, which divided by $y+z$, gives $\frac{\dot{y}}{z}$ or $\frac{\dot{y}}{\sqrt{bb+yy}}$, whose product multiplied by $\frac{1}{2} ab$ is equal to the fluxion proposed.

E X A M P L E.

If $y = \frac{1}{2}b\sqrt{3}$; then $\frac{y+z}{b} = \frac{1}{2}\sqrt{3+\sqrt{7}} = 2.1889$ nearly, whose tabular logarithm is .3402259, and this multiplied by 2.302585, *Art.* 58, gives .391699 $\times ab$, for the area of the sector CAM.

L E M M A. *Fig. 30.*

62. If from the center C of an ellipsis, and with half the greater axis CA as radius, a quadrant AND of a circle be described, and thro' the intersections of the same ordinate the lines CN, CM, are produced so as to meet the tangent AT in T, t; and from the point t a line tQ be drawn parallel to CT, and Am perpendicular to it, then will AQ be equal to CB half the second axis, and Am equal to the ordinate PM of the ellipsis.

By the parallels we have PM : PN ($:: At : AT$) :: AQ : AC, and by the property of the ellipsis, PM : PN :: CB : CD or CA. Therefore AQ : AC :: CB : AC, by equality of ratios, and $AQ = CB$.

The right-angled similar triangles AmQ, CPN, give AQ : CN or CB : CA :: Am : PN, and by the property of the ellipsis CB : CA :: PM : PN; therefore Am : PN :: PM : PN, by equality of ratios, and $Am = PM$.

P R O B L E M.

63. To express the area of an elliptical sector CAM by parts of the tangent At.

Let $CA = a$, $CB = b$, $At = z$; then the square of Ct is to the square of CM, as the fluxion $\frac{1}{2}az$ of the triangle CAAt is to the fluxion of the sector CAM, by *Art.* 25. But Ct : CM :: CT : CN or CA, and by the right-angled similar triangles CAT, QAt, we have CT : CA :: Qt : QA or CB; hence Ct : CM :: Qt : CB, by equality of ratios. Therefore $bb + zz :: bb ::$

$bb :: \frac{1}{2}a\dot{z} :: \frac{abb\dot{z}}{2bb + 2zz}$ equal to the fluxion of the sector CAM; whose fluent is a circular arc described by the radius CB (b) and tangent is At (z) multiplied by $\frac{1}{2}a$, *Art. 50.*

E X A M P L E.

Let $2z = b$; that is, the radius being 100000, the tangent will be 50000, which answers to an angle of 26 degrees and 34 minutes, or $26\frac{1}{3}$ degrees nearly; this multiplied by $.017453 \times b$, *Art. 54*, gives $.46358 \times b$, and this by $\frac{1}{2}a$ gives $.23179 \times ab$ for the required area.

P R O B L E M.

64. To express the same sector CAM, by parts of the ordinate PM (y) to the axis CA.

Let the tangent at M meet CA produced in V; then CA is to CV, or CP to CA, by the property of the ellipsis, as the fluxion $\frac{1}{2}a\dot{y}$ of the triangle CAM, is to the fluxion of the sector CAM, by *Art. 23*, and by the similar triangles AmQ, CPN, CP : CN or CA :: Qm : QA; or because AQ = b , Qm = $\sqrt{bb - yy}$,

we have $\sqrt{bb - yy} : b :: \frac{1}{2}a\dot{y} : \frac{aby}{2\sqrt{bb - yy}}$ equal to the fluxion of the sector; whose fluent is a circular arc described with the radius is b and sine y , multiplied by $\frac{1}{2}a$, by *Art. 22.*

E X A M P L E.

Let $3y = 2b$; then if the radius is 10000, the sine y will be 66666, which answers to an arc of 41 degrees and 48 minutes nearly, or $41\frac{4}{5}$: this being multiplied by $.01745 \times b$, *Art. 54*, gives $.72941 \times b$, and therefore $.3647 \times ab$, will express the area required.

This method of finding fluents, by means of tables of sines and logarithms, saves much trouble, when they depend on the quadratures of the conic section; for this could no otherwise be done than by infinite series,

series, which in most cases converge but slowly. If the arcs or logarithms are required to a greater degree of exactness, the rule in *Art.* 381, of my Elements of Mathematics, must be applied.

We shall give a few general problems for finding the fluents of binomial expressions, and shew when they may be found exactly or reduced to the quadrature of the conic sections.

L E M M A.

65. Suppose $z^{8n}v^{s+1}$, whose fluxion $8nz^{8n-1}v^{s+1} + s+1, \dot{v}z^{8n}v^s$, may be reduced into this form, $8nz^sv + s+1, \dot{v}z \times$ by $z^{8n-1}v^s$.

Now, if $v=c+fz^n$, then $\dot{v}=nfz^{n-1}$, and the values of v and $\dot{v}$, wrote into the last fluxions, gives $8nc\dot{z} + 8nfz\dot{z} + s+1, n\dot{z}fz^n \times$ by $z^{8n-1}v^s$, or $8e + 8 + s+1, fz^n \times n\dot{z}z^{8n-1}v^s$.

P R O B L E M.

66. To find the fluent of $\dot{z}z^{r-1} \times e + fz^{n1}$, when r is greater than unity.

Suppose $v^{s+1} \times$ by $az^{rn-n} + bz^{rn-2n} + cz^{rn-3n} + dz^{rn-4n} \dots \pm xF$ to be the fluent sought, F being the fluent of the remainder and x its coefficient, both to be determined.

The fluxion of this fluent, being disposed into a proper order, by making $r+s=q$, gives $qfa + \frac{r-1, ca}{q-1, bf} z^{-n} + \frac{r-2, eb}{q-2, fc} z^{-2n} + \frac{r-3, ec}{q-3, fd} z^{-3n} + \frac{r-4, ed}{q-4, dx} z^{-4n} \times$ by $n\dot{z}z^{r-1}v^s$.

For, if instead of 8 we write $r-1, r-2, r-3, \&c.$ into the fluxion $8e + 8 + s+1, fz^n \times n\dot{z}z^{8n-1}v^s$, we shall get the fluxions of the several terms of the supposed fluent.

Now,

Now, the coefficient qfa of the first term being made equal to unity, and those of the other terms, each equal to 0; the supposed fluent will come that required.

$$qfa = 1$$

$$a = \frac{1}{qf}$$

$$\overline{r-1}, ea + \overline{q-1}, fb = 0$$

$$-b = \frac{r-1}{q-1} \times \frac{ea}{f}$$

$$\overline{r-2}, eb + \overline{q-2}, fc = 0$$

$$c = \frac{r-2}{q-2} \times \frac{eb}{f}$$

$$\overline{r-3}, ec + \overline{q-3}, fd = 0$$

$$-d = \frac{r-3}{q-3} \times \frac{ec}{f}$$

And because $x = \overline{r-4}, edn$ and $\dot{F} = \dot{x}z^{rn-4n-1}v^s$; when $\dot{x}z^{rn-4n}v^{s+1}$, is the last term; by the same reason $x = negQ$ and $\dot{F} = \dot{x}z^{gn-1}v^s$, when $Qz^{gn-1}v^{s+1}$ is the last term.

The coefficients above being wrote into the fluent, gives $\frac{v^{s+1}z^{rn-n}}{nqf} - \frac{r-1}{q-1} \times \frac{Ae}{fz^n} + \frac{r-2}{q-2} \times \frac{Be}{fz^n} - \frac{r-3}{q-3} \times \frac{Ce}{fz^n} + \frac{r-4}{q-4} \times \frac{De}{fz^n} - \dots \pm negQF$.

The letters A, B, C, D, &c. express each its preceding term, and the remainder F must always have a contrary sign to that of the last term of the series; the coefficient or letter Q is always equal to that of the last term of the fluent.

C O R.

67. Hence, if r be any whole positive number, and q a fraction, or a whole positive number greater than r , or else when s is a whole negative number greater than unity, the fluent will be finite and contain r terms; but when r is a fraction, the fluent depends upon that of F , and in this case g must be the least fraction of r .

C O R.

C O R.

68. When r is negative, the fluent of $\dot{z}z^{-rn-1} \times e+fz^n$ may be found by the former. For if the term z^{-rn-1} , without the radical sign, be multiplied by z^n , and that $e+fz^n$ divided by z^n , we get $\dot{z}z^{sn-rn-1} \times e z^{-n} + f^s$, without changing its value; and as $\frac{sn-rn}{-n} = r-s$, this value wrote into that of $q=r+s$, found above, instead of r , gives $q=r$, and writing q for r in the numerator, r for q in the denominator, as well as $-n$, f , e , z^{-n} instead of n , e , f , z^n , respectively, the last fluent becomes $\frac{z^{n-qn}v^s+1}{-ner} \times \frac{q-1}{r-1} \times \frac{A}{e} fz^n - \frac{q-2}{r-2} \times \frac{B}{e} fz^n + \frac{q-3}{r-3} \times \frac{C}{e} fz^n - \dots \pm gnf QF.$

Observe we have here $v=ez^{-n}+f$ or $v=\frac{e+fz^n}{z^n}$ and $q=r-s$.

C O R.

69. Hence, if q or its equal $r-s$ is a whole number, and r a fraction or any whole positive number greater than q , the fluent will be finite, and consist of q or $r-s$ terms; and when q is a fraction, g must be the least positive fraction of r , and the series must be continued to $r-s$ terms.

E X A M P L E.

70. Let $\frac{\dot{z}z^{rn-1}}{e+fz^n}$ be the fluxion, and r any whole positive number; then will F be the fluent of $\frac{\dot{z}z^{rn-1}}{e+fz^n}$, which is the hyperbolic logarithm of $\frac{e+fz^n}{e}$, divided by fn , by *Art. 57*; and since s is here -1 , the general fluent

ent, in *Art.* 36, becomes in this case $\frac{z^{rn-n}}{nqf} - \frac{r-1}{q-1} \times$
 $\frac{Ae}{fz^n} + \frac{r-2}{q-2} \times \frac{Be}{fz^n} - \frac{r-3}{q-3} \times \frac{Ce}{fz^n} \dots - \frac{e}{f} QF$; and
 the series must be continued to $r-1$ terms.

When r is negative, the fluent of $\frac{z^{rn-n-1}}{e+fz^n}$ depends
 on that of $\frac{z^{rn-1}}{e+fz^n}$, which is the hyperbolic logarithm
 of $\frac{ez^{-n}+f}{f}$, by *Art.* 57, divided by $-ne$; and the

general fluent in *Art.* 68. becomes $\frac{z^{rn-qn}}{-ner} + \frac{q-1}{r-1} \times$
 $\frac{A}{e} fz^n - \frac{q-2}{r-2} \times \frac{B}{e} fz^n - \frac{q-3}{r-3} \times \frac{C}{e} fz^n \dots - \frac{f}{e} F$.

This fluent must be continued to r terms.

PROBLEM.

71. To find the fluent of $z^{rn-1} \times e+fz^n$ in a differ-
 ent manner.

Suppose $v=e+fz^n$ and $z^{rn}v^s \times$ by $a+b\frac{z^n}{v} \times c\frac{z^{2n}}{vv}$
 $+d\frac{z^{3n}}{v^3} + \&c.$ to be the fluent sought. The flu-
 xion of $z^{rn}v^s$, will be $rnz^{rn-1}v^s + sz^{rn}\dot{v}v^{s-1}$, or writ-
 ing the values nfz^{n-1} of $\dot{v}$, becomes $rnz^{rn-1}v^s +$
 $nfsz^{rn+n-1}v^{s-1}$; or which is the same, $r + fs\frac{z^n}{v} \times$
 by $nz^{rn-1}v^s$. Now, writing $r+1, r+2, r+3, \&c.$
 for r , and at the same time $s-1, s-2, s-3, \&c.$ for
 s , this last fluxion will become that of the first, se-
 cond, third, $\&c.$ term of the supposed fluent, each
 being multiplied by its respective coefficient $a, b, c, d,$
 and placed in proper order, gives

E

ar +

$$ar + \frac{\overline{s-0}, af}{r+1, b} \left\{ \frac{z^n}{v} + \frac{\overline{s-1}, bf}{r+2, c} \left\{ \frac{z^{2n}}{vv} + \frac{\overline{s-2}, cf}{r+3, d} \right\} \frac{z^{3n}}{v^3} \right\} \times \text{by } n z^{rn-1} v^s.$$

Now if the coefficient ar of the first term be made equal to unity, that of the proposed fluxion, and each of the rest equal to o , we get

$$ar = 1,$$

$$a = \frac{1}{r}$$

$$\overline{s-0}, af + \overline{r+1}, b = o \quad -b = \frac{s-0}{r+1} \times af$$

$$\overline{s-1}, bf + \overline{r+2}, c = o \quad c = \frac{s-1}{r+2} \times bf$$

$$\overline{s-2}, cf + \overline{r+3}, d = o \quad -d = \frac{s-2}{r+3} \times cf$$

Therefore $\frac{z^{rn} v^s}{rn} - \frac{s-0}{r+1} \times \frac{A}{v} f z^n + \frac{s-1}{r+2} \times \frac{B}{v} f z^n - \frac{s-2}{r+3} \times \frac{C}{v} f z^n + \&c.$ will be the fluent required: the letters A, B, C, &c. express as usual, each the term that preceeds it.

C O R.

72. Hence, when s is any whole positive number, and r a broken or a whole positive number, the fluent will be finite and expressed by $s+1$ terms.

And when e and f are both positive, this series will always converge; but when e is positive and f negative, the former series will always converge, and therefore should be used, when it is not finite or can be reduced to the quadrature of the conic section.

There is besides a method whereby this last series may be brought to converge faster, as follows: Let

$\frac{z}{1+zz}$ be the fluxion proposed, which being reduced into a series by a continued division, and the fluent taken,

taken, gives $z - \frac{1}{3}z^3 + \frac{1}{5}z^5 - \frac{1}{7}z^7 + \&c.$ and there will remain $\frac{z^{2r}}{1+z^2}$, t expresses the number of the preceding terms, the fluent of this remainder added to the preceding terms, will make up the fluent of the fluxion proposed $\frac{z}{1+z^2}$.

Now, the fluent of $\frac{z^{2r}}{1+z^2}$, by the general form above, is $\pm \frac{z^{2r}}{2rv} + \frac{1}{r+1} \times \frac{z^2}{v} A + \frac{2}{r+2} \times \frac{z^2}{v} B + \frac{3}{r+3} \times \frac{z^2}{v} C + \&c.$ because $n=2, f=1=e, 2r-1=2t$, or $2r=2t+1$.

But if $z=1$, the first series becomes $1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \frac{1}{9} \&c.$ and the sum of the 10 first terms or of 5 of this $\frac{2}{3} + \frac{2}{35} + \frac{2}{99} + \frac{2}{192} + \frac{2}{323}$, gives .76045,99049; and as t is 10 in this case, we get $r=\frac{21}{2}$, we have also $v=2$; and hence the series of the remainder becomes $\frac{1}{42} + \frac{1}{23}A + \frac{2}{25}B + \frac{3}{27}C + \frac{4}{29}D + \frac{5}{31}E + \&c.$ which reduced into decimals gives

.02380,95283
103,51967
8,28158
92018
12692
2047
372
75
16
4

.02493,82585

The sum of these two numbers gives .78539,81634, for an arc of a circle of 45 degrees true in all its places: whereas the first series would require more than 1000000000 terms.

P R O B L E M.

73. To find the fluent of $z z^{rn-1} v^s \times$ by $a + b z^n + c z^{2n} + d z^{3n} +$, &c. whence v becomes o .

If A expresses the fluent of $z z^{rn-1} v^s$ and B that of $z z^{rn+n-1} v^s$: then will $ernA + r + s + 1, n f B = z^{rn} v^s + 1$ by Art. 65; because 8 there is r here; if $r + s = q$, when v or its equal $e - f z^n$ becomes equal to o , we get

$$B = \frac{r+o}{q+1} \times \frac{e}{f} A, \text{ because } f \text{ is here negative. By writing}$$

$r+1, r+2, r+3$, into the fluxions of A and B , they become those of the succeeding terms C, D ;

$$\text{thus } B = \frac{r+o}{q+1} \times \frac{e}{f} A \text{ becomes } C = \frac{r+1}{q+2} \times \frac{e}{f} B, D =$$

$$\frac{r+2}{q+3} \times \frac{e}{f} C. \text{ Hence if we write the value of } B \text{ into}$$

$$\text{that of } C, \text{ we get } C = \frac{r+o, r+1}{q+1, q+2} \times \frac{ee}{ff} A; \text{ and this va-}$$

$$\text{lue of } C \text{ into that of } D, D = \frac{r+o, r+1, r+2}{q+1, q+2, q+3} \times \frac{e^3}{f^3} A.$$

Consequently, these values multiplied by their respec-

tive coefficients gives $A \times$ by $a + \frac{r+o}{q+1} \times \frac{be}{f} + \frac{r+o, r+1}{q+1, q+2} \times \frac{cee}{ff} + \frac{r+o, r+1, r+2}{q+1, q+2, q+3} \times \frac{de^3}{f^3}$ for the fluent re-

quired. Observe, 1°. That the points between the factors $r+o, q+1$, signify multiplication: that is, as

$r+o \times r+1 \times r+2$, and $q+1 \times q+2$, &c. 2°. All but one of the coefficients a, b, c, d , may be o : thus if

$$a=b=c=o; \text{ then } A \times \frac{r+o, r+1, r+2}{q+1, q+2, q+3} \times \frac{de^3}{f^3}, \text{ will be}$$

the fluent.

3°. If θ expresses any whole positive number, the fluent of $z z^{rn-1} v^s \times d z \theta^n$, will have as many factors $r+o, r+1, r+2$, and $q+1, q+2, q+3$, of both sorts as there are units in θ , multiplied by $\frac{de\theta}{f\theta}$.

P R O

P R O B L E M. Fig. 31.

74. Let the curve LMN be such that any ordinate PM=y, whose distance AP from a given point A is z, may be expressed by the equation $y = A + zB + \frac{z \cdot z - 1}{1 \cdot 2} C + \frac{z \cdot z - 1 \cdot z - 2}{1 \cdot 2 \cdot 3} D +, \&c.$ to find the coefficients A, B, C, D, &c.

In the base AP, take the equal parts AB, BC, CD, DE, &c. which being each to AP as unity to z, and draw ordinates thro' these points of divisions, which call, a, b, c, d, e, &c. ; then by making z equal to 0, 1, 2, 3, 4, &c. successively, in the equation the ordinate y becomes equal to a, b, c, d; this gives the following equalities.

$$\begin{array}{ll} a=A & A=a \\ b=A+B & B=b-a \\ c=A+2B+C & \text{hence } C=c-2b+a \\ d=A+3B+3C+D & D=d-3c+3b-a \\ e=A+4B+6C+4D+E & E=e-4d+6c-4b+a. \end{array}$$

Consequently if these ordinates are given, the coefficients A, B, C, are given also.

P R O B L E M.

75. The equation of the curve LMN being $y = A + zB + \frac{z \cdot z - 1}{1 \cdot 2} C +, \&c.$ to find the area ALMN.

If y and its equal be multiplied by z, and the fluent of each term taken as usual, we get z x by $A + \frac{1}{2}Bz + \frac{1}{2}C \times \frac{z \cdot z}{3} - \frac{z}{2} + \frac{1}{6}D \times \frac{z^3}{4} - \frac{3z^2}{3} + \frac{2z}{2} + \frac{1}{24}E \times \frac{z^4}{5} - \frac{6z^3}{4} + \frac{11z^2}{3} - \frac{6z}{2} +, \&c.$ For the area required, the continuation is plain : for the numerators are formed by continually multiplying the terms of the series z,

$z-1$

$z-1$, $z-2$, $z-3$, &c. the divisors of the several powers of z , by unity more than their indices and the divisors of the coefficients A , B , C , the products of the numbers $1, 2, 3, 4, 5, 6$, after the second term.

E X A M P L E S.

76. Hence, if $z=2$, we get $z \times \overline{A+B+\frac{1}{6}C}$, nearly for the area.

If $z=3$, then $z \times$ by $A+\frac{3}{2}B+\frac{3}{4}C+\frac{1}{8}D$.

If $z=4$, $z \times$ by $A+2B+\frac{5}{3}C+\frac{2}{3}D+\frac{7}{90}E$. Or because $A=a$, $B=b-a$, $C=c-2b+a$, $D=d-3c+3b-a$, $E=e-4d+6c-4b+a$, by *Art.* 74, we get $A+B+\frac{1}{6}C = \frac{c+4b+a}{6}$. $A+\frac{3}{2}B+\frac{3}{4}C+\frac{1}{8}D = \frac{d+3c+3b+a}{8}$:

and $A+2B+\frac{5}{3}C+\frac{2}{3}D+\frac{7}{90}E = \frac{7e+32d+12c+32b+7a}{90}$.

Hence, if A expresses the sum of the first and last ordinate, B the sum of the second and last but one, C that of the two next to these, and D , E , &c. always the sum of the two next, till you come to the two middle ones, or if the number of terms is odd, the last letters denote the middle term. The following table, where the number of ordinates is expressed by Roman figures, will give the areas nearly, z expressing always the base or distance between the first and last ordinates.

$$\text{III. } \frac{A+4B}{6}z$$

$$\text{IV. } \frac{A+3B}{8}z$$

$$\text{V. } \frac{7A+32B+12C}{90}z$$

$$\text{VI. } \frac{19A+75B+50C}{288}z$$

$$\text{VII. } \frac{41A+216B+27C+272D}{840}z$$

VIII.

$$\text{VIII. } \frac{751A + 3572B + 1323C + 2989D}{17280} z$$

$$\text{IX. } \frac{989A + 5888B - 928C + 10496D - 4540E}{28350} z;$$

To apply this method to some examples, we shall suppose $\frac{z}{1+z}$ to be the fluxion whose fluent is required when $z=1$. Suppose the base z to be divided into six equal parts, viz. $\frac{1}{6}, \frac{2}{6}, \frac{3}{6}, \frac{4}{6}, \frac{5}{6}, \frac{6}{6}$; then writing o and these parts for z into $\frac{1}{1+z}$, we get the seven following ordinates $1, \frac{6}{7}, \frac{6}{8}, \frac{6}{9}, \frac{6}{10}, \frac{6}{11}, \frac{6}{12}$; hence the form $\frac{41A + 216B + 27C + 272D}{840}$, for seven ordinates, gives $A = 1 + \frac{6}{12} = \frac{3}{2}$, $B = \frac{6}{7} + \frac{6}{11} = \frac{108}{77}$, $C = \frac{6}{8} + \frac{6}{10} = \frac{27}{20}$, $D = \frac{6}{9} = \frac{2}{3}$; these values multiplied by their respective coefficients, and reduced into decimals, give

$$\begin{array}{r} 61.50000 \\ 302.96078 \\ 36.45000 \\ 181.33333 \\ \hline 582.24411 \end{array}$$

and the sum divided by 840 gives .693147 for the area required, which is the hyperbolic logarithm of 2, true in all its places.

Suppose the fluxion $\frac{z}{1+zz}$, whose fluent is required, when $z=1$; divide the base z into five equal parts, as $\frac{1}{5}, \frac{2}{5}, \frac{3}{5}, \frac{4}{5}, 1$; by writing o and these values into $\frac{1}{1+zz}$, we get these six ordinates $1, \frac{25}{26}, \frac{25}{29}, \frac{25}{34}, \frac{25}{41}, \frac{1}{2}$; now by

by the area of six ordinates $\frac{19A + 75B + 50C}{288}$, we get

$A = 1 + \frac{1}{2} = \frac{3}{2}$, $B = \frac{25}{26} + \frac{25}{41} = \frac{1675}{1066}$, $C = \frac{25}{29} + \frac{25}{34} = \frac{1575}{986}$, and these values multiplied by their respective coefficients give

$$\begin{array}{r} 28.5000 \\ 117.8470 \\ 79.8600 \\ \hline 226.2070 \end{array}$$

This sum divided by the divisor 288, gives .7854 for an arc of a circle of 45 degrees, whose radius is unity. Observe that when the base z is more or less than unity, the area found as above must be multiplied by it, as it is marked in the tables: observe also, that the more ordinates made use of, the nearer the area will be; and this method can be applied when all others fail, that is, when the ordinates are expressed by *furd* quantities, as will be seen hereafter.

*A GENERAL METHOD of finding the ROOTS
of Equations by Approximation.*

MANY attempts have been made by authors to find the roots of equations; but as their operations are in general very intricate and tedious, and besides no general rule has been given by any so as to answer all cases, excepting *Daniel Bernoullie*, in the *Memoirs at Petersburgh*; but as he has given no direct demonstration, and, I think, is mistaken in some cases, the reader will be pleased to find here a very concise demonstration, as also a general rule, expressed in fewer words, than could hardly have been expected in so extensive a subject.

P R O-

P R O B L E M.

77. To find the least value of x in the general equation $1 = ax + bx^2 + cx^3 + dx^4 + \dots$, &c.

If we divide both sides by x , we get $\frac{1}{x} = a + bx + cxx + dx^3 + \dots$, &c. and if we suppose x so small that all the terms multiplied by it may be neglected as inconsiderable, we get $x = \frac{1}{a} = \frac{1}{A}$ nearly. To get a nearer value of x , write $\frac{1}{A}$ for x in bx , and by neglecting the higher powers, we get $\frac{1}{x} = a + \frac{b}{A}$, or $x = \frac{A}{aA + b} = \frac{A}{B}$. To carry on the approximation, write this last value of x into bx and its square, or which is nearly the same, and more convenient, the product $\frac{1}{A} \times \frac{A}{B} = \frac{1}{B}$, of the preceding values, into cx^2 , and neglecting the higher powers, we get $\frac{1}{x} = a + \frac{bA + c}{B}$, or $x = \frac{B}{aB + bA + c} = \frac{B}{C}$. Writing again this last value of x into bx , the product $\frac{B}{C} \times \frac{A}{B} = \frac{A}{C}$, of the two last values of x into cx^2 , and the product $\frac{B}{C} \times \frac{A}{B} \times \frac{1}{A} = \frac{1}{C}$, of the three last values of x into dx^3 , by neglecting the higher powers, if there are any more; then $\frac{1}{x} = a + \frac{bA + cA + d}{C}$, or $x = \frac{C}{aC + bB + cA + d} = \frac{C}{D}$. By continuing thus to write the last value found for x , the product of the two last for x^2 , the product of the three last for x^3 , and the product of the four last for x^4 , and repeating these operations, the value of x may be approximated to any degree of exactness. The continuation of this rule is evident by this table.

F

 $I = I$

$$1=1$$

$$A=a$$

$$B=aA+b$$

$$C=aB+bA+c$$

$$D=aC+bB+cA+d$$

$$E=aD+bC+cB+dA.$$

G E N E R A L R U L E.

78. Find a series, whose first term being unity, the second, the first coefficient, then to find any other term, multiply constantly the coefficients each by the respective term of the series, in a contrary order; that is, the first a by the last term, the second b , by the last but one, the third c by the last but two, &c. and the sum of all these products will be the next term of the series; by continuing thus the series at pleasure, the more terms the better, then the last term but one divided by the last, will be the root required nearly.

E X A M P L E I.

Let $1=x+2x^2+3x^3+4x^4$; then will $1:1$; the next term $1 \times 1 + 1 \times 2 = 3$; the next $1 \times 3 + 2 \times 1 + 3 \times 1 = 8$; the next to this $8 \times 1 + 3 \times 2 + 3 \times 1 + 4 \times 1 = 21$; and the next to this $21 \times 1 + 8 \times 2 + 3 \times 3 + 4 \times 1 = 50$, by continuing thus we get $1.1.3.8.21.50.128.323.813.2043$; and the last term but one 813 divided by the last 2043 , gives $x=.39729$ nearly, or $x=.3973$. To know how near the value of x has been found, find the several powers of x , and write them into the proposed equation, and you will find how far the approximation has been carried on. Thus here I find $1=10001$.

E X A M P L E II.

Let $1=-2x+5x^2-4x^3+x^4$, be the equation, then will 1 and -2 be the two first terms, the third $-2 \times -2 + 5 = 9$, the fourth $-2 \times 9 + 5 \times -2 - 4 \times$

$$1=$$

$1 = -32$, the fifth $-2 \times -32 + 9 \times 5 - 4 \times -2 + 1 = 118$; hence the series will be $1. - 2. 9. - 32. 118. - 434. 1595$, and $x = \frac{-4 \cdot 3 \cdot 4}{1 \cdot 3 \cdot 9 \cdot 3} = -.272$ nearly.

If a positive root is required, suppose $x = 2 + y$; then the equation will be changed into this $1 = 2y + 5y^2 + 4y^3 + y^4$, and $1:2.5.24.82.306.1123$; whence $y = .272$, and $x = 2.272$, nearly.

When a term of the equation is wanting, write 0 in its place, and proceed as before; thus if $1 = 2x + x^3$; write $1 = 2x + 0 + x^3$; then 1.2, will be the two first terms, $2 \times 2 + 0 \times 1 = 4$ the third, $2 \times 4 + 0 \times 2 + 1 = 9$ the fourth, $2 \times 9 + 0 \times 4 + 2 = 20$, and so 1.2.4.9.20.44.97.214.472; hence $x = \frac{2 \cdot 1 \cdot 4}{4 \cdot 7 \cdot 2} = .4534$ nearly; since when this value is wrote into the equation it gives $1 = 10000$.

If $1 = 0 + 3x^2 + 2x^3 + x^4$, then because the root found is alternately greater and less, by supposing $x = .5 - y$, the equation will be changed into this $1 = 80y - 120y^2 + 48y^3 + 16y^4$; and 1.80.6280.492848. Hence $y = .012742$ and $x = .487257$, true in all its places.

When the equation is not equal to unity, observe the following

GENERAL RULE.

Find the nearest whole number to the root, which may be done by inspection or trials; suppose that number more or less a variable quantity, with an indetermined coefficient equal to the variable quantity of the equation, by which the equation is changed into another, then if this indetermined coefficient be made equal to the known term, the equation will by division be reduced to unity, and always converge.

EXAMPLE III.

Let $2 = x + x^2 + x^3$, the equation, as x is less than unity, suppose $x = 1 - y$; then will $1 = 6y - 4y^2 + y^3$; and 1.6.32.169.892.4708: now 892 divided by 4708,

F 2

gives

gives $y=.18946$, and $x=.810535$, true in all its places.

E X A M P L E IV.

Let $2=-x-x^2+x^3$; if $3-ay=x$, the equation becomes $13=20ay-8aay^2+a^3y^3$. Hence if $a=13$, according to the rule we get $1=20y-104y^2+169y^3$, by division; and $1.20.296\ 4009.52776, 688608.8960977.116523452$: therefore $y=.0769$, and $3-13y=x=2.0003$ nearly; the true root being 2.

Let $14=2x+x^3$, whose root is 2 nearly; by making $2+ay=x$, this equation will be changed to this, $2=14ay+6aay^2+a^3y^3$. Hence $a=2$, according to the rule, and $1=14y+12y^2+4y^3$. Hence, $1.14.208.3084.45728$, and $y=.067442$; therefore $x=2.134884$, true in all its places.

Lastly, let $90=x+xx+x^3$, whose root is 4 nearly; let $x=4+ay$, this will change the equation to this, $6=57ay+13aay^2+a^3y^3$; and if $a=6$, it becomes $1=57y+78y^2+36y^3$, by division: and $1.57.3327.194121.11326455$; so that $y=.017138725$, and $x=4.10283235$.

This method is also excellent, in numbers, when the root is required to a great many places. Thus to find the square root of 26, which is 5 nearly. Suppose it to be $5+y$: then $26=25+10y+y^2$, or $1=10y+y^2$. Hence $1.10.101.1020.10301.104030$, and $y=\frac{1.0301}{104030}$; therefore $5+y=5.09901951360$, true in all its places. The square root of 487 is 22 and 3 over: if $22+3y$, be the root sought, then $487=484+132y+9y^2$, or $1=44y+9y^2$, by reduction. Hence $1.44.1939.85448.3765529$, and $22+3y=22.06807, 64908$, its root required.

If the root of a small number be required to a great many places, multiply that number by such a square number so as the product may be a square within unity, then the root of the product divided by the root of the multiplier, gives the root required. Thus to find
the

the square root of 3, multiply it by 16 the square of 4, which gives 48; then suppose $7-y$ to be the root: hence $48=49-14y+y^2$, or $1=14y-y^2$, by five steps only we get 1.73205,08075,689, for the root required.

To find the square root of 2, multiply it by 144 the square of 12, and supposing $17-y$ to be the root; then will $1=34y-y^2$, from whence by 4 terms the root will be 1.41921,35627.

To find the cube root of 20, which is 3 less 7; let $3-7y$ be the root, then will $1=27y-63y^2+49y^3$, and 1.27.666.16330.400275.9811269; hence $7y=.2855823$, and $3-7y=2.7144177$, the root required.

Lastly, the cube root of 2 is found by multiplying it by 64 the cube of 4, and supposing it to be $5+3y$, which gives $1=75y+45y^2+9y^3$; hence 1.75.5670.428634, and $3y=.0396842$, and $5+3y$ divided by 4, gives 1.25992105. The rule holds good in any other power.

The greatest root of an equation may be found in the same manner; for suppose $x^n=ax^{n-1}+bx^{n-2}+cx^{n-3}---+d$; to be the equation, and if we suppose

$$x=\frac{1}{y}, \text{ we get } x^n=\frac{1}{y^n}, x^{n-1}=\frac{y}{y^n}, x^{n-2}=\frac{y^2}{y^n}, x^{n-3}=\frac{y^3}{y^n}.$$

Hence $1=ay+by^2+cy^3---dy^n$, which gives this

GENERAL RULE.

79. Let the greatest power of the unknown quantity be on one side with unity for its coefficient, and all the other terms in their order on the other; find a series of numbers as before; then the last term divided by its antecedent gives the root nearly.

EXAMPLE.

Let $x^4=5x^3+4x^2+2x+2$; then 1.5.29.167.963.5551.31999, and $x=\sqrt[4]{\frac{1.999}{5.551}}=5.7645$ nearly.

Again,

Again, let $4x^4 = 3x^2 + x + 1$; make $x = \frac{1}{y}$; then $4 = 3y + y^2 + y^3$; or if $y = 1 - z$, we get $1 = 8z - 4z^2 + z^3$, and 1.8.60.449.3360.25144: hence $z = \frac{3360}{25144} =$ and $1 - z = y = \frac{21784}{25144}$. Therefore $x = \frac{25144}{21784} = 1.15424$, nearly: for writing this value into the equation it gives $1 = 00001$.

OF EXPONENTIAL QUANTITIES.

When a quantity has a variable index, it is called *Exponential*, as a^x , y^x , &c.

80. To find the fluxion of a^x , y^x , or $\overline{a+y}^x$.

I. Let $a^x = z$, taking the logarithms on both sides, which denote by the quantities with an l before them; then $xla = lz$, whose fluxion is $\dot{x}la = \frac{\dot{z}}{z}$, by *Art. 57* and multiplying both sides, the first by a^x and the second by z its equal, we get $\dot{x}la \times a^x = \dot{z}$ for the fluxion required.

II. Let $y^x = z$, then by logarithms $xly = lz$, whose fluxion is $\dot{x}ly + \frac{\dot{xy}}{y} = \frac{\dot{z}}{z}$, or because $y^x = z$; this fluxion becomes $\dot{x}ly \times y^x + xy\dot{y}x^{-1} = \dot{z}$, the fluxion required.

III. Lastly, let $\overline{a+y}^x = z$, then $xla + y = lz$, whose fluxion by logarithm is $\dot{x}la + y + \frac{\dot{xy}}{a+y} = \frac{\dot{z}}{z}$; or $\dot{x}la + y + \frac{\dot{xy}}{a+y} \times$ by $\overline{a+y}^x = \dot{z}$. From hence this rule,

Take the logarithms of the exponentials and multiply the fluxion of this logarithm by the exponential quantity.

As these kind of quantities occur but seldom in computations, we shall dwell no longer upon this subject; and only observe, that if the letter a in a^x denotes the number 2.71828, whose hyperbolic logarithm is unity; then because $la = 1$, the fluxion $\dot{x}la \times a^x$, of a^x , becomes $\dot{x}a^x$; on the contrary, the fluent of

$\dot{x}a^x$

Δa^x will in this case be a^x . This remark will be useful hereafter.

SECTION V.

Of MOTION in General considered in a non-resisting Medium.

MECHANICS, of which the laws of motion are the foundation, in regard to the necessities in life, is without doubt the most useful branch of all mathematics, and at the same time the most extensive. It is become more valuable since it has been applied to astronomy: for which reason the greatest mathematicians have made it their chief study preferably to all other branches: each making use of what was discovered before, endeavoured to improve it, and to extend its limits, in such a manner that this inestimable science is arrived to such a high degree of perfection, that it seems to be impossible to add any thing new worth the public notice.

Galileus discovered the laws of motion of bodies which fall in a non-resisting medium, and applied it to the curve described by a body projected with a given velocity and in a given direction.

M. Descartes, observing the force by which a stone, turned round in a sling, endeavours to recede from the hand, opened a new field in regard to the laws of motion, much more extensive than *Galileus* had done. *Mr. Huggens*, who called this force *centrifugal*, applied himself strongly to the discovery of its laws: he after many trials, gave at last thirteen theorems, but concealed the demonstrations.

As soon as *Sir Isaac Newton* had knowledge of it, his extraordinary penetration led him to the discovery of the laws of motion, by which a moving body is drawn, or compelled out of its direction so as to describe any curve whatever; he named this force *centripetal*: but not satisfied of having demonstrated these laws

laws in a non-resisting medium, he applied them likewise in any medium whatever: another discovery much more extensive than all that was done before him upon this subject.

Kepler having found, by a great many astronomical observations, that the orbits of the planets were ellipses, and that the sun was placed in one of the foci as center, that their periodic times were to each other as the areas described by the lines drawn from the center of the body to the center of the forces or focus; and that the squares of the periodic times were to each other as the cubes of their mean distances: the incomparable *Sir Isaac Newton* gave mathematical demonstrations of these laws; and he found besides, that the centripetal forces which retain the body in an ellipsis were in the inverse ratio of the squares of their distances from the focus, and that the same law which retains one body in its orbit, retains likewise all others in theirs. From thence followed the universal force of gravity: it was upon this principle that he composed his system of the world, which certainly will render his name immortal.

Since his principles on natural philosophy have been published, the greatest mathematicians of Europe were occupied to unfold some parts of it; for he treats every thing in a very short and succinct manner, and gives synthetical demonstrations, which renders most of them very obscure: on the other hand, that book contains almost every part of the mathematics and natural philosophy; so that there remains very little hopes of any more discoveries: therefore all that could be done since, was either to demonstrate some parts in an easier manner, or to extend them farther than the author has done.

The intent of this section is, to reduce all what is dispersed in different parts of *Sir Isaac Newton's* principle concerning motion, under one view, by a general system of motion, in as clear and easy a manner

ner as the subject will admit of. All what is given depends on two general theorems only, well known by those who have written upon motion; the rest follows as necessary consequences.

After having given the first principles of motion in a non-resisting medium, we endeavour to unfold the ratios of the forces which impel or attract bodies, that is, the centripetal forces, which tend towards a point, or whose directions are parallel; and to render this problem more conspicuous, examples are given in the conic sections of both cases.

The centripetal forces, which tend towards one of the foci, being always as the squares of the distances inversely, the times as the areas described by the line drawn from the focus to the center of the body, and the squares of the periodic times, as the cubes of the mean distances, and if these laws agree with *Kepler's* observations, in regard to the planets, it was reasonable to conclude that the celestial bodies describe ellipses in their revolutions about the sun; which occupies one of the foci. Nevertheless this uniformity of the laws observed by *Kepler*, and those deduced from the theory, some mathematicians have doubted, whether there might not be some other curve that had the same property as the ellipsis. Tho' it appeared to be impossible, yet we gave the inverse of this problem; that is, the centripetal force being as the squares of the distances inversely, to find the curve described by the body, which is found again to be a conic section, that is an ellipsis, for it cannot be a parabola or hyperbola, because they open or recede from the axis continually.

It has been found by the late observations on the transit of Venus, that the horizontal parallax of the sun was 9.732 seconds, from whence the mean distance from the earth to the sun is found, and expressed by semidiameters of the earth, and thence the distances of the other planets are found by proportion.

G

Now,

Now, knowing the apparant diameters of the planets by observation, as well as their periodic times, it is easy to find their true diameters and solid contents, as well as the densities of those which have satellites; as for those that have none, we have no method to find theirs.

Tho' our method is general, and may be applied to all the different laws of which the centripetal force may be endowed with, yet we have given no other examples than when they are as the squares of the distances inversely, as being the only one true in nature, for the sake of him who endeavours to instruct himself: for which reason we have here given such questions only as are either very useful or instructive; and we think to have not a little obliged the learner, to save his time and the trouble of reading a large volume.

T H E O R E M. *Fig 32.*

80. *The fluxion of the velocity generated by a power varying according to any law, is as the rectangle made by the fluxion of the time elapsed, and the measure of that power at that time.*

Let AM be such a curve, that when the abscissa AP denotes the time, the corresponding ordinate PM shall denote the measure of the power: then because the effects are proportional to their adequate causes, the instantaneous velocities are as the measures of the powers that generates them; and since the area AMP contains all the measures of the power during the time AP, that area will also express the sum of all the velocities generated in that time. Consequently, the fluxion of that area is equal to the fluxion of the velocity: therefore the fluxion of the velocity is equal to the rectangle made by the fluxion of the time AP, and the measure PM of the power at that time, by *Art. 16.*

T H E O-

T H E O R E M.

81. *The fluxion of the space described by a motion varying according to any law, is as the rectangle made by the fluxion of the time elapsed, and the velocity at that time.*

If the ordinate PM denotes the velocity at the end of the time AP; then by the same argument as in the last theorem, the area AMP will denote the sum of all the velocities, as well as the total space described by these velocities, during the time AP; and therefore the fluxion of that area, or the rectangle made by the fluxion of the time AP and the velocity PM, will express the fluxion of the space described during that time.

C O R.

82. Hence, if p expresses the measure of the variable power, v the velocity generated, in the time t and z the space described; then will $\dot{v} = pi$ by *Art.* 80, and $\dot{z} = vi$, by *Art.* 81. If we expell i , by means of these two equations, we get $p\dot{z} = v\dot{v}$, and if i be supposed constant whilst $\dot{z}$ and v vary; the fluxion of $\dot{z} = vi$ will be $\ddot{z} = \dot{v}i$, or $\ddot{z} = p\ddot{t}$, by writing pi instead of its equal $\dot{v}$.

E X A M P L E.

83. Let the motion be uniform; then will the velocity v be uniform, and the fluent of $z = vt$ will be $z = tv$. From whence it follows, 1°. That the spaces described uniformly at the same time are as the velocities. 2°. The spaces described uniformly with the same velocity, are as the times. 3°. Lastly, spaces described uniformly with unequal velocities, and in different times, are as the velocities and times conjointly.

E X A M P L E.

84. If the generating power be constant, such as gravity near the surface of the earth the fluent of $\dot{v} = pi$, will be $v = pt$, that of $p\dot{z} = v\dot{v}$, $2pz = vv$: and

expelling v we get $2z = ptt$; in the same manner, by expelling p , we get $2z = vt$. Hence,

1°. The velocities acquired by a falling body freely from a state of rest, and being acted upon by gravity only, are by $v = pt$ as the times elapsed.

2°. The spaces described in the descent, are by $2zp = vv$, as the squares of the velocities, or by $2z = ptt$, as the squares of the times; or else by $2z = vt$, as the times and velocity conjointly.

3°. Lastly, The space described uniformly, with the velocity acquired by a body falling freely from a state of rest, is double the space descended thro' by the body. For since $z = tv$, by *Art.* 83, when the motion is uniform, and $2z = tv$, when accelerated, and as v and t are the same; the space z described uniformly is double the space z descended thro' by the accelerated motion in the same time.

R E M A R K.

As it is not readily conceived, how the measures of powers, velocities and time, can be compared together, it is necessary to observe that the time denotes the order in which the other quantities are generated or produced; thus one hour, minute or second, being taken for unity, then that unite is to any number of unites as one number to any other number, and the spaces described by falling bodies, as well as the velocities acquired, are in certain ratio's expressed by numbers, and therefore are as one line to another line; and thus we are enabled to compare these several quantities together, in as clear and distinct a manner as if they were of the same kind, as will appear by what follows.

C O R.

85. If in the equation $2z = ptt$, *Art.* 84, we suppose the time t to be one second, then will $2z = p$. Which shews that the force of gravity p is measured by twice the space fallen thro' freely from a state of rest, in a second: hence when the space fallen thro' is given, the

the time during the fall is also given, and the contrary: Thus the space fallen thro' in a second is found by experiments, to be 16 feet nearly; then $p=32$ feet; and if the space z be 100 feet, then $2z=ptt$, gives $100=16tt$, whose square root is $10=4t$, or $2.5=t$. Again, suppose t to be 6 seconds, then $z=16 \times 36$, or $z=576$ feet.

C O R.

86. Again, the equation $2pz=vv$; gives $\sqrt{2pz}=v$; or if b be the height fallen thro' in a second, then $p=2b$; and $2\sqrt{bz}=v$: Which shews, that twice the square root of the product of the space fallen thro' in a second, and the space fallen thro' in any time, will express the number of feet described uniformly by the velocity acquired in the fall, in a second. For since the squares of the velocities are as the spaces fallen thro', *Art.* 84. N°. 2. $\sqrt{b}$ will be to $\sqrt{z}$, as the velocity acquired by falling thro' the height b to the velocity acquired thro' the height z : and the velocity $\sqrt{b}$ is to the velocity $\sqrt{z}$, as $2b$ twice the space described uniformly in a second is to the space $2\sqrt{bz}$ described uniformly at the same time, *Art.* 48. N°. 3.

E X A M P L E.

87. To find the number of feet described uniformly per second, by a velocity acquired in falling thro' a space of a 100 feet: then $z=100$, and $b=16$, we get $bz=1600$ feet, whose square root 40 taken twice gives 80 for the space described.

Of CENTRAL FORCES.

WHEN a body is continually attracted towards a point fixed, by a force acting according to any given law, that force is called *centripetal*; but when a body continually recedes from a given point, that force is called *centrifugal*.

Many

Many objections have been made at first against the word *attraction*, as if it implied an occult quality; whereas it means no more than the force by which the bodies are retained in their orbits. Whether it proceeds from the pression of some other matter, or from an inherent quality of matter, which causes revolving bodies to tend towards a point fixed, is not material, for it is certain, that there is such a power in nature, otherwise bodies would move in a right line, tangent to the curve described, by the projectile force; and therefore it is indifferent by what name it is called, provided its meaning is understood. The French and German mathematicians, in order to avoid this word, have substituted that of *solicited*, which is far from conveying any true idea of its meaning, and seems to be coined merely for the sake of contradiction; and, what must astonish an English reader, several of our best authors have made use of this expression.

GENERAL PROBLEM. *Fig. 33.*

88. *If a body projected from a given point A in a given direction AE, with a given velocity, be attracted by a force which acts according to any law, towards a point fixed C; to find the nature of the curve described by the body.*

From any point M in the curve let a tangent be drawn, and from the point fixed C, the perpendicular CS to it. From any point T in the tangent, draw TR perpendicular to the ray CM: then if P denotes the centripetal force at M, c the given velocity at A, v the variable velocity at M, t the time of describing the arc AM; and we call $AM=z$, $CM=y$, $CS=s$; then $MT=\dot{z}$, $MR=\dot{y}$ $TR=\dot{x}$, *Art. 4*; and $\dot{z}=vt$, *Art. 82*. Now by the similarity of triangles we have CM to MS or TM($\dot{z}$) to MR($\dot{y}$) as the centripetal force P in the direction CM to its effect $\frac{\dot{y}}{\dot{z}}P$, in the direction TM of the tangent; this force $\frac{\dot{y}}{\dot{z}}P$ being

wrote

wrote for p into $\dot{z}=pi$, *Art. 82*, gives $P\dot{y}i=\dot{v}\dot{z}$, or expelling i , by means of the equation $\dot{z}=vt$, *Art. 82*. we get $P\dot{y}=vv$.

Again, CM is to CS, as the force P in the direction CM is to its effect $\frac{s}{y}P$ in the direction CS perpendicular to the tangent; and by the similitude of triangles TM:MR::CM:MS= $\frac{y\dot{y}}{\dot{z}}$; and the radius MS is to the radius MT as the fluxion $\dot{s}$ of the circular arc described by the point S is to the fluxion $\frac{\dot{s}\dot{z}^2}{y\dot{y}}$, of the circular arc described by the point T: but this last space is described by the force $\frac{s}{y}P$, in the direction CS: if therefore we write $\frac{\dot{s}\dot{z}^2}{y\dot{y}}$ for $\dot{z}$ and $\frac{s}{y}P$ for p , into the equation $\dot{z}=pi$, *Art. 82*, we get $\frac{\dot{s}\dot{z}^2}{y\dot{y}}=i\dot{t}\times\frac{s}{y}P$, or because $\dot{z}=vt$, by expelling $i\dot{t}$, we shall have $P\dot{y}s=vv\dot{s}$.

By means of these two equations, we shall determine the nature of the curve when the law of the centripetal force P is given.

N.B. We used the equation $\dot{z}=pi$, to find $P\dot{y}s=vv\dot{s}$, because the space $\frac{\dot{s}\dot{z}^2}{y\dot{y}}$ described by the force in the direction CS perpendicular to the tangent, is a second fluxion in respect to the space $\dot{z}$ described in the direction of the tangent; besides, if any other expression had been used, the equality of the fluxions would not have been preserved. Observe likewise, that when y increases while the velocity v decreases, the first equation must be $-P\dot{y}=v\dot{v}$.

C O R.

C O R.

89. Since we have always $\pm P\dot{y} = v\dot{v}$, whatever line or curve the body describes, it is manifest, that the velocities generated by the same centripetal force P , are always equal at the same distance y from the point fixed C ; and these generated by different centripetal forces, are as these forces at equal distances.

C O R.

90. Hence if P be expelled by means of the equations $-P\dot{y} = v\dot{v}$ and $P\dot{y}s = vv\dot{s}$, we get $-v\dot{s} = s\dot{v}$, or $v\dot{s} + \dot{v}s = 0$, whose fluent is $vs = A$. If $CE = d$, be perpendicular to the tangent at A ; then when v becomes c , s becomes d ; and hence $cd = A$, or $vs = cd$. Which shews, that in any concave curve towards the point fixed C , the velocity v is every where inversely as the perpendicular drawn to the tangent, let the centripetal forces be what they will.

C O R.

91. If v be expelled by means of the equations $\dot{z} = vi$ and $vs = cd$, we get $\dot{z}s = cdi$; and as $CM:CS::TM::TR$, or $\dot{z}s = y\dot{x}$; hence $y\dot{x} = cdi$; and since $y\dot{x}$ is the fluxion of the area $2CMA$, we have $2CMA = cdi$. Consequently, in any concave curve towards the point fixed C , the times are always as the areas described by the ray CM drawn to the point of contact, let the centripetal force be what it will.

C O R.

92. If the center C be removed at an infinite distance from the vertex A , so that the rays CA , CM , may become parallel and equal; then if the sin of the angle CAE be k and the radius unity, then $CA(y):CE(d)::1:k$, or $d = ky$; hence $s\dot{z} = y\dot{x}$ found above becomes $ks\dot{z} = d\dot{x}$; whose fluxion is $k\dot{s}\dot{z} = d\dot{x}$, by supposing $\dot{z}$ constant: now the values of s and $\dot{s}$ wrote into the equations $P\dot{y}s = vv\dot{s}$, $vs = cd$; give $P\dot{y}\dot{x} = vv\dot{x}$ and $vx = ck\dot{z}$; and as $y\dot{x} = c(d)\dot{i}$, by the last, and $d = ky$, we get $\dot{x} = ck\dot{i}$, or $x = ckt$.

The

The same otherwise. Fig. 34.

Let the ordinate PM be perpendicular to the base AB, draw RS perpendicular to the tangent MT, call $PM=y$, $AP=x$, and the rest as before; then the force P in the direction MP is to the effect in the direction RS perpendicular to the tangent, as MR to RS, or as $MT(\dot{z})$ to $RT(\ddot{z})$, by the similarity of triangles; which therefore is $\frac{\dot{x}}{\dot{z}}P$: and if r expresses the radius of curvature CM: then the radius CM is to the radius MT, as the fluxion $\dot{z}$ of the arc described by the first radius is to $\frac{\dot{z}^2}{r}$ the fluxion of the arc described by the second, or the fluxion of the space described by the power $\frac{\dot{x}}{\dot{z}}P$. Whence $\frac{\dot{x}}{\dot{z}}P \times ii = \frac{\dot{z}^2}{r}$, by Art. 82, or because $\dot{z}=vi$, we get $Pr\dot{x}=vv\dot{z}$. Now since $r=\frac{y\dot{z}}{\dot{x}}$, by Art. 42, No. 2, when $\dot{z}$ is constant the last equation becomes $Pj\dot{x}=vv\ddot{x}$, the same as above: we have likewise $-r=\frac{y\dot{z}^2}{\dot{x}\dot{z}}$, by the same Art. 42, No. 2, when $\dot{x}$ is constant and $\dot{z}\dot{z}=y\ddot{y}$, the equation $Pr\dot{x}=vv\dot{z}$; becomes $P\dot{z}^2+vv\ddot{y}=0$, in this case.

C O R. Fig. 35.

94. Hence if P be expelled by the equations $-Pj=v\dot{v}$ and $Pj\dot{x}=vv\ddot{x}$, we shall have $v\ddot{x}+\dot{v}\dot{x}=0$, whose fluent is $v\dot{x}=A\dot{z}$, because $\dot{z}$ was supposed constant: but the cosine k of the angle EAP, is to the radius 1, as $\dot{x}:\dot{z}$ or $k\dot{z}=\dot{x}$ at A; and as $v=c$, in that point, we get $A=ck$; and the equation $v\dot{x}=A\dot{z}$, becomes $v\dot{x}=ck\dot{z}$, the same as above: from this equation and $\dot{z}=vi$, we get $\dot{x}=cki$ or $x=ckt$.

H

COR.

C O R.

95. Hence it appears from the equation $\frac{vx}{z} = ck$; that the velocity v reduced into the horizontal direction AK, is constant and expressed by ck , and in the direction of the tangent AE, by constant velocity c .

E X A M P L E. *Fig. 36.*

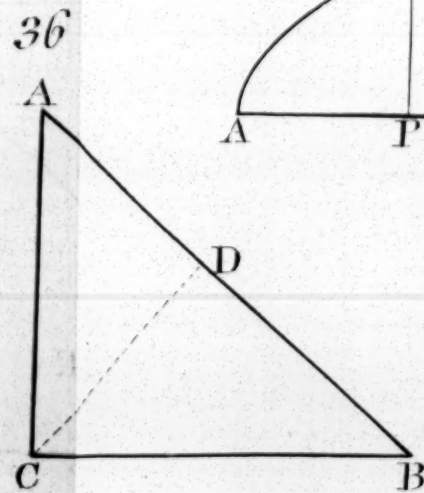
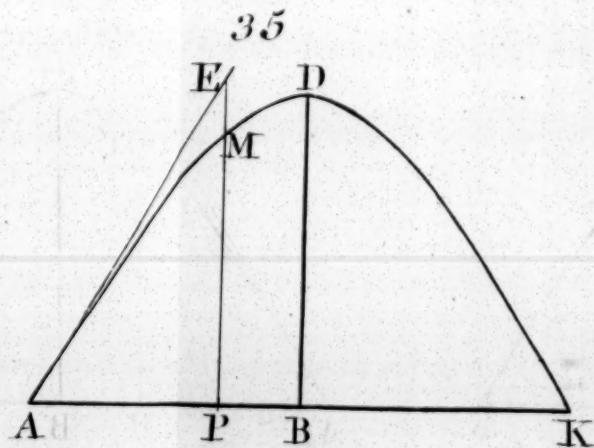
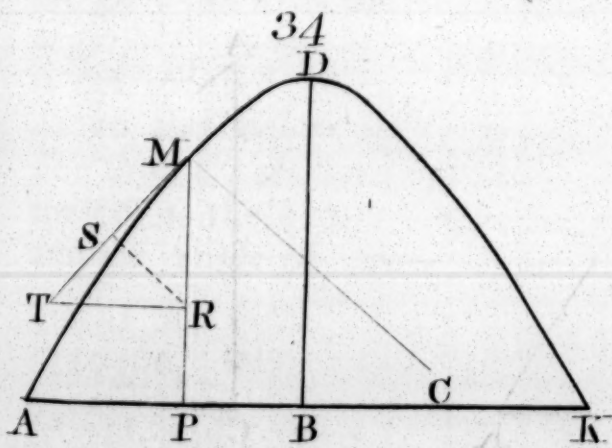
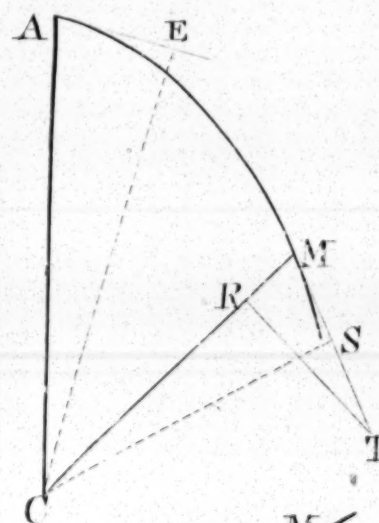
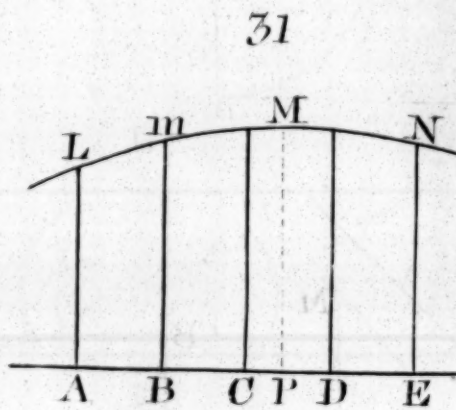
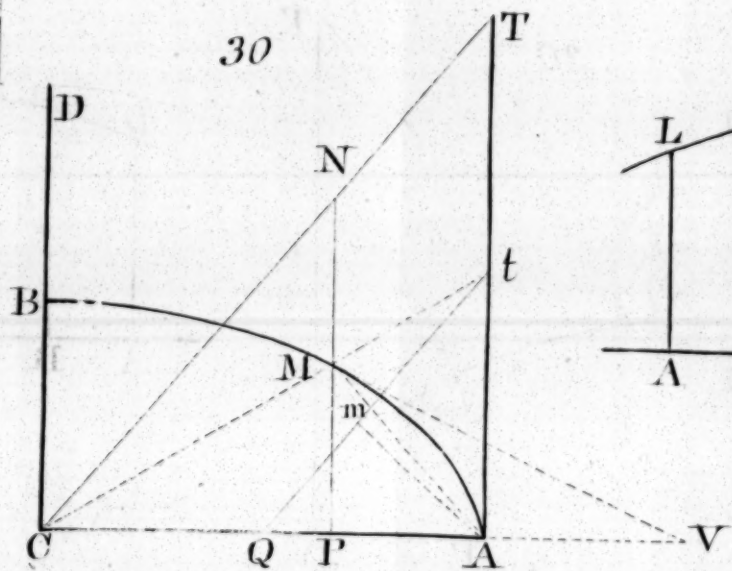
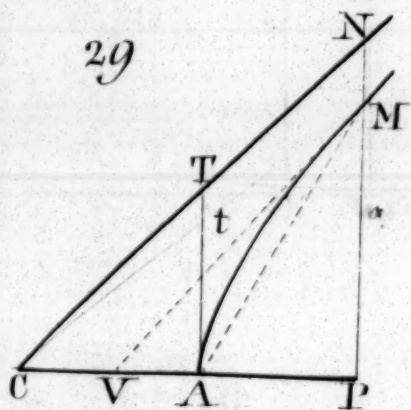
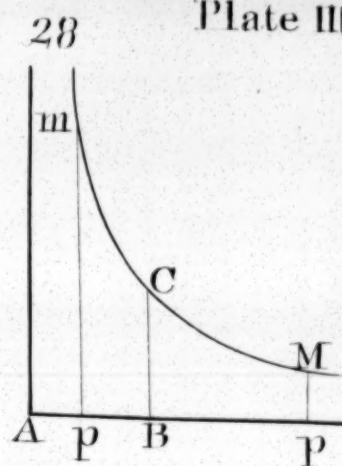
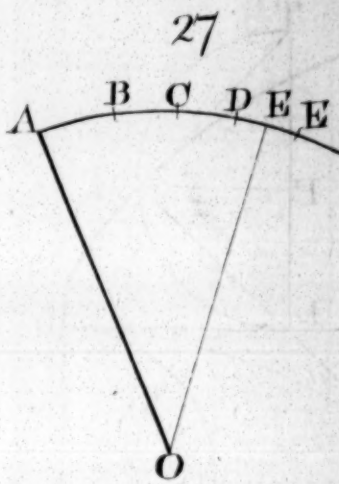
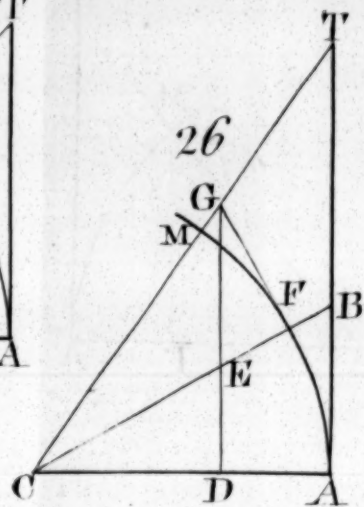
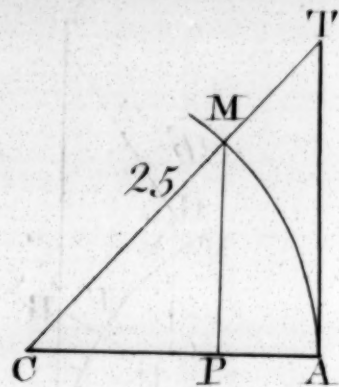
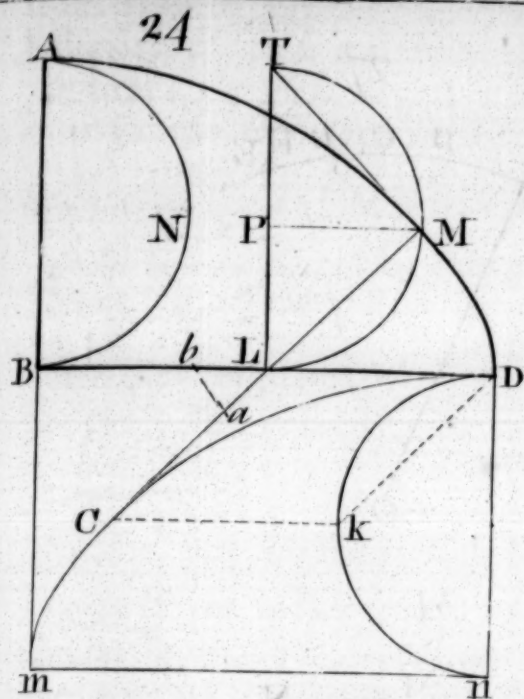
96. Let a body descend upon an inclined plane AB, freely from a state of rest by the force of gravity: then the velocity acquired in the descent is equal to the velocity acquired by falling thro' the height AC of that plane; for the tangent is here at right angles, and by *Art. 95.* the velocity equal to c that by which the body is projected upwards. And since spaces described by the same velocity are as the times, the time of the descent is to the time of the fall as AB is to AC: or if CD be perpendicular to AB; then as AB: AC:: AC: AD, the time of the descent is to the time of the fall as AC is to AD.

C O R. *Fig. 37.*

97. Hence, if from the extremity A of a vertical diameter AE, any chords AB, AC, AD, are drawn in a semicircle, the body will descend thro' either of these chords in the same time that it would fall thro' the diameter AE, by the last; and the velocity acquired in descending thro' any number of contiguous inclined planes, is always equal to the velocity acquired by falling thro' the vertical height of all these planes.

E X A M P L E. *Fig. 35.*

98. Let the force P be constant and its directions parallel; by producing the ordinate PM, so as to meet the tangent AE in E, and calling l the sine of the angle PAE, the rest the same as before; then whilst the body moves uniformly in the direction AE from A to B, with the given velocity c , it will descend
from





from E to M, by the force of gravity, and since spaces fallen thro' from a state of rest are as the squares of the times, by *Art.* 84, if d denotes the height fallen thro' in the first second, and $x = ckt$, by *Art.* 94, we have $1 : \frac{xx}{cckk} :: d : EM = \frac{dxx}{cckk}$; and $k : l :: x : \frac{lx}{k} = PE$.

Hence $PE - EM = PM$, or $y = \frac{lx}{k} - \frac{dxx}{cckk}$. Or if b denotes the height thro' which a falling body acquires the projectile velocity c , or $4db = cc$, *Art.* 85; we get $y = \frac{lx}{k} - \frac{xx}{4bkk}$.

C O R.

99. If we make $y = 0$, we get $AK = x = 4blk$, for the horizontal range; and if $y = 0$, $AB = x = 2blk$, for the abscissa corresponding to the greatest ordinate $BD = bl$. Hence,

1°. The horizontal ranges $4blk$, with the same projectile velocity $2b$, are as the sines $2lk$ of an angle double the angle of elevation.

2°. The horizontal ranges, with the same angle of elevation, are as the projectile forces.

3°. The abscissa AB corresponding to the greatest ordinate, is half the horizontal range AK .

C O R.

100. It is manifest, that the horizontal range is the greatest, with the same projectile force, when the angle of elevation is 45 degrees; since the product lk is the greatest when $l = k$; and all ranges equally above and below 45 degrees are equal; since the sine in one case becomes the cosine in the other, and the cosine the sine. Thus let $b = 5280$ feet, $d = 16$, and the angle of elevation 45 degrees: then $l = k$, $2lk = 1$ or $k = \frac{1}{2}\sqrt{2}$. Hence $AB = 5280$ feet, $AK = 10560$, $BD = 2640$, $c = 2\sqrt{db} = 581.3$ feet per second, and $t = \frac{x}{ck} = 25.68$ seconds.

Geometrical Construction. Fig. 38.

On the base AK, erect AL perpendicular and equal to b , describe the semi-circumference of a circle AEL, intersecting the direction AT in E; thro' the point E draw ED perpendicular to AK; then $AK = 4AD$ will be the range, and BF drawn perpendicular to the middle of AK and equal to DE, will be the greatest height of the projectile.

For the angles ALE, DAE, which insist on the same arc being equal, the right-angled triangles LEA, ADE are similar; therefore the radius is to the sine l as LA(b) is to AE= bl , and the radius is to the cosine k as AE(bl) is to AD= blk ; hence $4AD = 4blk = AK$. Lastly, the radius is to the sine l as AE(bl) is to DE= $BF = bl$. When the range AK and the height AL are given, the perpendicular DE determines the angle of elevation DAE.

If the object be placed above or below the level of the battery, and q the tangent of the angle it makes with the level: then as $1:q::x:y=qx$; the equation $y = \frac{lx}{k} - \frac{xx}{4bkk}$, becomes $\pm q = \frac{l}{k} - \frac{x}{4bkk}$; or if a be the tangent and s the secant of the angle of elevation, we get $a = \frac{b}{k}$, $s = \frac{1}{k}$, and $x = a \mp q \times \frac{2b}{ss}$

E X A M P L E. *Fig. 39.*

101. Let DMA be half a cycloid, BD its base placed horizontally, AB the diameter of its generating circle; it is required to find the time of descent thro' any arc AM.

Draw MP perpendicular to AB, and join the chord AN; then as the arc AM of the cycloid is double the chord AN, by *Art. 44*, and the velocity acquired in the descent thro' the arch, MA is equal to that acquired in the fall thro' the vertical height AP, *Art. 97*; but this velocity is as the square root of AP,

AP, *Art.* 84, N°. 2; or because AP:AN::AN:AB, the square root of AP is as the chord AN, that is as the arc AM. Since the velocity is as the space described, the time of descending thro' any arc of the cycloid, terminated by the lowest point A, is consequently equal, let that arc be what it will.

C O R.

102. Hence, if $AP=y$, $AB=a$, and d the height fallen thro' in a second; then $-2dy=v\dot{v}$, *Art.* 88, whose fluent is $A-4dy=vv$; and when $y=a$, v is $=0$, and $A=4ad$; therefore $2\sqrt{ad-dy}=v$: but $\dot{z}=2\sqrt{ay}$, by *Art.* 44, or $\dot{z}=\frac{a\dot{y}}{\sqrt{ay}}$; these values of

v and $\dot{z}$ wrote into $\dot{z}=vi$, *Art.* 82, gives $\frac{a\dot{y}}{2\sqrt{ay-yy}}$ $=i\sqrt{ad}$. If c denotes an arc of a circle whose diameter is a and the sine $\sqrt{ay-yy}$; the fluent of this fluxion will be $c=t\sqrt{ad}$, *Art.* 4, N°. 2, and when $y=0$, c expresses half the circumference and t the time of descent thro' half the cycloid DMA.

C O R.

103. Hence, if the time of the descent thro' half the cycloid be half a second, then $c=t\sqrt{ad}$ becomes $c=\frac{1}{2}\sqrt{ad}$, or $4cc=ad$; or $2a:2tt::aa:4cc$. Consequently, when the vibrations are performed in a second, the length $2a$ of the pendulum is to the force of gravity $2d$ as the square of the diameter a is to the square of the circumference $2c$.

C O R.

104. Since $c=t\sqrt{ad}$ gives $cc=ttad$ or $c\sqrt{\frac{a}{d}}=at$,

and $\sqrt{\frac{a}{d}}$ denotes the time of descent thro' the diameter AB, *Art.* 84, N°. 2; it follows, that the time of descending

descending thro' the diameter BA, is to the time of one vibration in a cycloid as the diameter a is to the circumference $2c$.

C O R.

105. Because $c = t\sqrt{ad}$, gives $a \times \frac{cc}{ac} = dt$, and the ratio of the diameter to its circumference is constant, as well as d in the same latitude: the times of vibrations by pendulums of different lengths, are in a subduplicate ratio to that of their lengths.

C O R.

106. If a and b are the lengths of two pendulums which perform their vibrations in the time t and T ; then by the last $\sqrt{a} : \sqrt{b} :: t : T$, or $T\sqrt{a} = t\sqrt{b}$; and if we suppose the first performs n number of vibration whilst the second performs but one; then $nt = T$, and $T\sqrt{a} = t\sqrt{b}$, becomes $n\sqrt{a} = \sqrt{b}$, or $na = b$. Consequently, if one of these lengths is given the other is given likewise.

C O R.

107. Lastly, since 3.14159, &c. denotes the circumference of a circle, whose diameter is unity, *Art.* 53, the equation $a \times \frac{cc}{aa} = d$, (*Art.* 104.) t being unity gives $9.869a = d$, when the vibrations are performed in a second. Hence, if the length of the pendulum is given, the space fallen thro' in a second is given likewise and the contrary. Dr. Hally has found by experiment, that the length of such a pendulum is 39.125 inches in the latitude of London; whence, if $2a = 39.125$, then $d = 193.0623$ inches, or $d = 16.088$ feet; for the height fallen thro' in a second.

Mr. de Mairand has found that length to be 440.9778 lines at Paris; so that if $2a = 440.567$, we get $d = 2173.9778$ lines, or $d = 15.097$ feet for that space.

The

The length of such a pendulum is commonly supposed to be 39.2 inches at *London*, which is a mistake, as will be shewn hereafter.

E X A M P L E. *Fig. 40.*

108. Let the body revolve in a parabola AM, the point fixed, C the focus and A the vertex, to find the force P. If $CA=a$, $CM=y$; then $CT=s=\sqrt{ay}$, by the nature of the curve; whence the equation $vs=cd$, in *Art. 92*, becomes $ac=v\sqrt{ay}$, because d is here a ; or $vv=\frac{acc}{y}$, whose fluxion is $-v\dot{v}=\frac{acc\dot{y}}{2yy}$, and since $-P\dot{y}=v\dot{v}$, by *Art. 88*, we get $P=\frac{acc}{2yy}$, and at the vertex A where $y=a$, we get $P=\frac{cc}{2a}$

E X A M P L E. *Fig. 41.*

109. Let the body revolve in an ellipsis AM about one of its foci C; to find the centripetal force P. From the other focus F draw FM, and let a, b , denote half the first and second axis; then $CT=s=\frac{by}{\sqrt{2ay-yy}}$, by the nature of the curve. Whence $vs=cd$ becomes $bvy=dc\sqrt{2ay-yy}$, or $vv=\frac{ddcc}{bb}\times\frac{2a-y}{y}$, whose fluxion is $-v\dot{v}=\frac{addcc\dot{y}}{bbyy}$ equal to $-P\dot{y}$, or $P=\frac{addcc}{bbyy}$. Which shews, that the centripetal force of a body revolving in an ellipsis or parabola, about the focus, is as the square of the distance inversely.

E X A M P L E. *Fig. 42.*

110. Let the point fixed C be the center of an ellipsis; then as the sum of the squares of two semi-conjugate diameters CM and CN is constant, suppose $=rr$; we have $CN=\sqrt{rr-yy}$, and $CN\times CT=ab=s\sqrt{rr-yy}$, by the property of the curve: therefore the

the equation $vs = cd$, becomes $bv = c\sqrt{rr - yy}$, because $d = a$ here. Hence $vv = \frac{cc}{bb} \times rr - yy$, whose fluxion is $-v\dot{v} = \frac{cc\dot{y}}{bb} = P\dot{y}$ or $P = \frac{ccy}{bb}$; the force P is therefore as the distance directly.

C O R.

111. If $a = b$, then the elipsis becomes a circle and $P = \frac{cc}{a}$. Which shews, that the centripetal force is as the square of the velocity directly and the radius inversely; and the velocity, as the square root of the rectangle made by the force and radius, likewise the radius, as the square of the velocity directly and the force inversely.

P R O B L E M. Fig. 41.

112. If the centripetal force be as the square of the distance inversely, to find the nature of the curve described.

Because $-P\dot{y} = v\dot{v}$, by Art. 88, and $P = \frac{I}{yy}$ by supposition, we get $-\frac{\dot{y}}{yy} = v\dot{v}$, whose fluent is $A + \frac{2}{y} = vv$. Now if the point G be taken in the vertical line CA produced, such that when CM (y) becomes CG or $2a$, the velocity v becomes 0; then will $A + \frac{2}{y} = 0$, or $-A = \frac{I}{a}$, and so $\frac{2}{y} - \frac{I}{a} = vv$; if $CA = d$, and the tangent at A be at right angles to CA , when y becomes $=d$, then v becomes $=c$, and the last equation $\frac{2}{d} - \frac{I}{a} = cc$. Now these values of vv and cc being wrote into $vvss = ccdd$, the square of $vs = cd$, gives $ss \times \frac{2a - y}{y} = 2ad - dd$, or if $2ad - dd = bb$, we get

$by =$

$by = \sqrt{2ay - yy}$, which is the equation of an ellipsis in respect to the foci; a, b , are half the first and second axes.

P R O B L E M.

113. To determine the periodic time of a body revolving in an ellipsis about one of its foci.

Let T express that time, a, b , half the first and second axis, and $r = 3.14159$, &c. that is, let r denote the circumference of a circle whose diameter is unity; then because $2CAM = cdt$, by *Art.* 90, and when the area CAM becomes that of the whole ellipsis, it is $=rab$, then t becomes T , and therefore $2abr = cdT$.

C O R.

114. Hence if a, b, d , become equal, the ellipsis degenerates into a circle, and $2ra = cT$; and as $aP = cc$ by *Art.* 111; the square $4rraa = ccTT$ of $2ra = cT$ gives $4arr = PTT$. Which shews, 1°. That the centripetal forces in circles, are as the radius directly, and the square of the periodic time inversely.

2°. The radius is as the centripetal force and the square of the periodic time conjointly.

3°. That in the same circle, the centripetal force, is as the square of the periodic time inversely.

4°. When the periodic time remains the same, the centripetal force is as the radius.

C O R.

115. If a body revolves in such a circle that the centripetal force be equal to that of gravity near the surface of the earth. Let d denote the space fallen thro' in a second, then $2d = P$, *Art.* 85, and $aP = cc$, gives $c = \sqrt{2ad}$; and $2ar = cT$, becomes $2ar = T\sqrt{2ad}$. N. B. The velocity is expressed by the number of feet described uniformly in a second and the time in seconds.

I

C O R.

C O R.

116. If in CA produced to H, so that the centripetal force be equal to the force of gravity, at the distance CH (b), the force of gravity being unity, we have $P = \frac{bb}{yy}$, and by following the computation of

Art. 112, we get $bb \times \frac{2a-y}{ay} = vv$, and $bb \times \frac{2a-d}{ad} = cc$,

or because $2ad - dd = bb$; we shall find $bb = cd\sqrt{a}$; and this value of cd wrote into $2arb = cdT$ gives $2ra^{\frac{3}{2}} = bT$, or $4rra^3 = bbTT$. Consequently, the squares of the periodic times of bodies revolving in different ellipses, about the same center or focus, are as the cubes of their mean distances. For $4rr$ and bb remain here the same.

The famous astronomer Kepler was the first who discovered this law by observations; that the times were as the areas described by the radius drawn from the focus to the body, and that the orbit of all heavenly bodies were elliptical, and that the planets were attracted by the sun, by forces that are as the squares of the distances inversely.

C O R.

117. When $y = a$ in the ellipsis, $bb \times \frac{2a-y}{ay} = vv$,

becomes $b = v\sqrt{a}$; and as $bb = cd\sqrt{a}$, Art. 116; when the ellipsis becomes a circle, we get $b = a$, and $b = c\sqrt{a}$ in the circle. It follows, that the velocity of a body revolving in an ellipsis at the mean distance is equal to the velocity of a body describing a circle at the same distance.

C O R.

118. Lastly, from $P = \frac{bb}{yy}$, it appears that the centripetal

tripetal forces of bodies revolving in ellipsis at equal distances and about different centers, are as the squares *bb* of the distances CH, from the center C to the point H, where the centripetal force becomes equal to the force of gravity near the surface of the earth.

Before we can apply the foregoing theory to astronomy, it is necessary to know the mean distance of the earth from the sun, and this distance can be obtained from no other means than by the sun's parallax; but this, by all the care and industry of the astronomers, has not yet been discovered to any degree of exactness. But to ascertain something, we shall make use of the result of all the last observations on the transit of Venus, given in the Transactions, which makes that parallax to be $9.732''$: and as the tangent of that angle is 47181,06744, the radius being 100000,00000,00000; then this tangent is to the radius as the semi-diameter of the earth (1) is to the distance required 21194. As the semi-diameter of the earth is 3959 English miles, that distance will be 83,907,046 miles. We shall however take 20000 semi-diameters for that distance in the following computations.

PERIODIC TIMES.

Saturn	Jupiter	Mars	Earth
10759.275,	4332.514,	686.9785,	365.2565,
Venus	Mercury	Moon	
244.6176,	87.9692,	27.31.	

These times have been accurately observed by most astronomers, and are expressed in days and decimals.

Since the squares of the periodic times are as the cubes of the mean distances, *Art.* 116, we have

Mean Distances from the Sun to

Saturn	Jupiter	Mars	Earth	Venus	Mercury
190752,	104017,	30472,	20000,	14463,	7742.

These distances are expressed in semi-diameters of the earth; and the distance of the moon from the earth has been found to be 60.11 of the same measure.

Diameters of the Planets.

Sun	Saturn	Jupiter	Mars	Earth	Venus
92.000,	7.260,	9.160,	0.523,	1.000,	1.033,
		Mercury	Moon		
		0.368,	0.270.		

These diameters are proportional to their apparent ones, which have been determined by observation. These bodies are nearly similar, their contents are then as the cubes of their diameters.

Solid Contents of the

Sun	Saturn	Jupiter	Mars	Earth	Venus
778688,	383.76,	769.53,	0.1426,	1.000,	1.1023,
		Mercury	Moon		
		0.510,	0.020.		

Hence, the contents of the planets are in no proportion to their distances from the sun, as might have been expected, if astronomers made no mistake. The sun is nearly a million of times bigger than the earth, and 670 times bigger than all the planets. Therefore the common center of gravity of the sun and the planets must fall within the surface of the sun; which accounts for the sun's motion being so slow, in comparison to that of the planets, as to appear to be at rest.

We have found $2ra^{\frac{3}{2}} = bT$ in *Art.* 116. To adapt this equation to practice, it must be observed, that when the earth is supposed in the center of the motion, the distance is then equal to the semi diameter of the earth or unity, by supposition; and as the moon revolves about the earth in 27.32 days, and its mean distance

distance is 60.11; we have $a^{\frac{3}{2}} = 466.04$; which gives this proportion $27.32 : 466.04 :: bT ; a^{\frac{3}{2}}$, or $bT = 0.0586a^{\frac{3}{2}}$.

From whence we may compare the force of gravity at equal distances, or the quantity of matter, of such bodies as have satellites, as follows:

If the sun be in the center, then having the distance of the earth $a = 20000$, and the periodic time $T = 365.2565$; we get $b = 453.83$, or $bb = 205961.67$, for the quantity of matter in the sun, when that of the earth is unity.

If Saturn be in the center, from Hugen's satellite, whose periodic time is 15.945 days, and its distance 135.65 from the center of Saturn, we get $b = 5.58$, or $bb = 31.14$. But if Jupiter be in the center, the periodic time of the farthest satellite is 16.6889 days, and distance 231.74; therefore $b = 12.25$ or $bb = 153.26$. Hence,

Quantity of Matter in the

Sun	Saturn	Jupiter	Earth
20596167,	3114,	15326,	100.

The quantity of matter in bodies that have no satellites, cannot be found by any method we know of.

As equal bodies placed on the surfaces of the planets, are attracted by forces, that are as the quantity of matter in the attracting body directly, and the squares of the distances inversely; if then the last numbers are divided by the squares of the semi-diameters, we shall get for the

Attractions of equal Bodies placed on the Surface of the

Sun	Saturn	Jupiter	Earth
97335,	2363,	7306,	1000.

And if D expresses the density of a body placed in the center of attraction, then as the quantity of matter

ter is as the density and magnitude or cube of the radius conjointly: we have $bb = Dy^3$, or as $Pyy = bb$, by *Art.* 118, we get $P = Dy$. If therefore the last numbers are divided by their respective diameters, we get

Density of the

Sun	Saturn	Jupiter	Earth
2116,	657,	285,	1000.

Which shews that the density of the sun is something more than twice that of the earth, and those of Saturn and Jupiter less.

We have given some of the principal applications of the general laws of motion laid down in this section, and which are particularly useful in astronomy; as to those generally given, when the centripetal force is as any other power of the distances, have no foundation in nature, we have purposely omitted them, as a mere speculation.

Problems of the Greatest and Least QUANTITIES.

L E M M A. *Fig.* 43.

119. *If from any two given points D, K, in a line AC, the indefinite lines DE, KT, are drawn at right angles to AC; and if a, m, are two constant quantities: It is required to draw lines from the given points A, C, to a point L in DE, so as $a \times AL - m \times KN$ shall be the least possible.*

Let the angle DAE be such that the radius is to its sine, in the compound ratio of a to m and of CD to CK, then the point E will be the required one. From the point D and any other L in DE, the lines DF, LP, are drawn at right-angles to AE; then because, 1°, $CD : CK :: DL : KN$, or, 2°, $a \times CD : m \times CK :: a \times DL : m \times KN$: and the right angled similar triangles EPL, EFD, and DFA, give, 3°, $DL : FP :: (LE : PE ::) AD : DF$; but, 4°, $AD : DF :: a \times CD : m \times CK$ by construction: therefore (No. 2, 3,

4)

4) $DL : FP :: a \times DL : m \times KN$, by equality of ratios, or $a \times FP = m \times KN$. Hence, $a \times AL - m \times KN$ becomes equal to $a \times \overline{AL - FP}$, or if upon AE we take $AM = AL$, then $a \times \overline{AL - m \times KN}$, or its equal $a \times \overline{AL - FP}$, becomes $a \times \overline{AF + PM}$: Which will be the least of all when PM, vanishes, because AF is constant; and this will happen when the point L falls upon the point E.

This lemma includes both M'Laurin's given in *Art.* 572 and 578; whereby all the problems called isoperimetrical are solved in the most clear and elegant manner, without making use of any higher order of fluxions than the first.

C O R.

120. Hence, if $CK = d$, $CD = y$, $DE = x$, and $AE = z$, then $DA : DF$ or $ay : md :: AE(z) : DE(x)$ or $axy = mdz$, when $a \times AL - m \times KN$ is the least possible. It must be observed that the radius DA must always be greater than the sine DF, or ay greater than md , of the angle DAE, otherwise the problem would be impossible.

C O R.

121. If the point K coincides with the point D, or which is the same, when CD and CL become parallel, CK becomes equal to CD; that is, $y = d$, and $KN = DL$; therefore $ax = mz$, when $a \times AE - m \times DE$ is the least possible, and this will be when the radius AD is to the sine DF as a is to m or in a constant ratio.

C O R.

122. It is manifest, that when $a \times AL - m \times KN$ is the least possible $\frac{1}{m} AL - \frac{1}{a} KN$ is likewise the least possible, since both expressions give the same construction; or from that either the greatest or least expression, divided or multiplied by the same

same constant quantity am , does not change the question.

C O R.

123. It is equally manifest, that when the first term $a \times AL$ or $\frac{1}{m}AL$ of the expression is constant or given, the second mKN or $\frac{1}{a}KN$ will be the greatest possible, when $a \times AL - m \times KN$, or $\frac{1}{m}AL - \frac{1}{a}KN$ is the least possible. Since $m \times KN$ or $\frac{1}{a}KN$ may increase from 0, but can never exceed $a \times AL$ or $\frac{1}{m}AL$.

C O R.

124. On the contrary, if the second term $m \times KN$ or $\frac{1}{a}KN$ is constant or given, the first $a \times AL$ or $\frac{1}{m}AL$ will be the least possible, when $a \times AL - m \times KN$ or $\frac{1}{m}AL - \frac{1}{a}KN$ is the least possible. Since the first term may increase ever so much, but can never become less or equal to the second.

P R O B L E M. *Fig. 39.*

125. To find the line DMA, drawn thro' two given points D, A, such that a body shall describe it in the shortest time possible, beginning to move from a state of rest, and being acted upon by the force of gravity only.

Draw the horizontal line DB and AB perpendicular to it; and from any point M, the line MP parallel to DB: then if $AB = a$, $AP = y$, the arc $DM = z$, $PM = x$,

PM = x , c , a constant velocity, and v the velocity at M: Hence $\frac{\dot{z}}{v}$ expresses the fluxion of the time in describing the arc DM by *Art.* 84, and $\frac{\dot{z}}{v} - \frac{\dot{x}}{c}$ the least possible, when $v\dot{z} = c\dot{x}$, by *Art.* 121. Now, when the point M coincides with the lowest point A, where $\dot{x} = \dot{z}$, we have $v = c$; so c denotes the velocity of the body at A: But the velocity at M is to the velocity at A, as the square roots of AB and AP, by *Art.* 84, N^o. 2, that is, $\sqrt{a-y} : \sqrt{a} :: v : c$. Hence $\dot{z}\sqrt{a-y} = \dot{x}\sqrt{a}$, this squared and $\dot{z}^2 - \dot{y}^2$ wrote for its equal $\dot{x}^2$, gives $a\dot{y}^2 = y\dot{z}^2$ or $\dot{z} = \frac{a\dot{y}}{\sqrt{ay}}$, whose fluent is $z = 2\sqrt{ay} + p$; which is the equation of the common cycloid by *Art.* 44; BD its base and AB the diameter of its generating circle, when $p = 0$.

P R O B L E M. *Fig.* 39.

126. If the length of the curve AMD be given, or remains the same, to find its nature such that the area ABD shall be the greatest possible.

If AM = z , PM = y , AP = x , and a a constant quantity; then will $a\dot{x} = y\dot{z}$, when $a\dot{z} - y\dot{x}$, is the least possible, by *Art.* 121; and as the fluent of $a\dot{z}$ is given by supposition, that of $y\dot{x}$ of the area AMP, will be the greatest possible, by *Art.* 123; but $a\dot{x} = y\dot{z}$ is the equation of the circle, by *Art.* 4. N^o. 2; since the sine y is to the radius a , as the fluxion $\dot{x}$ of the cosine is to the fluxion of the arc AM.

It is no less evident, that if the fluent of $y\dot{x}$, or the area ABD, is given or constant, the rectangle az or because a is given, the arc z will be the least possible, by *Art.* 124.

P R O B L E M.

127. If the surface described by the arc AM, about the axis AP be given; to find the nature of the curve such, that the solid described by the area AMP, in the rotation, shall be the greatest possible.

It is evident, that when $ayz - yyx$ is the least possible, we have $ax = yz$, by Art. 121; since the radius will be to the sine of the angle made by the ordinate PM and the tangent at M as ay to yy , or as a to y . Consequently the curve is again a circle. As the fluent of ayz denotes the surface described by the arc AM, when a is the circumference of a circle whose radius is unity, which being always the same, the fluent of yyx or the solid will be the greatest possible, by Art. 123.

C O R.

128. Hence, of all solids of the same boundary the sphere is the greatest, and of all equal solids the sphere has the least boundary. This has likewise been proved in Art. 32.

G E N E R A L C O R O L L A R Y.

129. If the radius be to the sine of the angle at M, or which is the same, if z be to x as $Ay + By^2 + Cy^3$, is to $a + by + cy^2$, and $z \times by$ as $Ay + By^2 + Cy^3$, is the least possible; we have $z \times by$ as $Ay + By^2 + Cy^3$ is to $a + by + cy^2$, by Art. 121; therefore, if the fluent of $z \times by$, $a + by + cy^2$ be given, that of $x \times by$, $Ay + by^2 + Cy^3$, will be the greatest possible by Art. 123; and on the contrary, when the fluent of $x \times by$, $Ay + By^2 + Cy^3$ is given, that of $z \times by$, $a + by + cy^2$, will be the least possible.

Observe that the coefficients a, b, c, A, B, C , are supposed constant but indeterminated; and that any one or more of them may be unity or 0.

To

To shew some particular example of this general expression, let the length of the arc AM, and the surface described by it about AP be given: to find the nature of the curve, such that the solid described by the area AMP in that rotation, shall be the greatest possible: if we make $A=C=c=0$, and $B=1$, we shall have $z\dot{y}\dot{y}=a\dot{x}+b\dot{y}\dot{x}$ for the equation of the curve required. And the sum of the length of the arc AM, and the surface described by it, will be the least possible, when the solid is given.

SECTION VI.

The MOTION of FLUIDS *and the* RESISTANCE of BODIES.

LEMMA. Fig. 44.

130. **T**HE force by which a fluid strikes a plane in a perpendicular direction, is to the force by which it strikes it in an oblique one, as the square of the radius is to the square of the sine of the incident angle.

With the radius CA describe a quadrant DBA of a circle, draw the radius CB at pleasure, and from the points D, B, the lines DL, BP, perpendicular to CB, CA; then if the radius CD expresses the absolute force of the fluid in a perpendicular direction, DL will express its effect against the inclined surface CB; or because of the equal right-angled triangles CLD, CPB, DL is equal to CP, so that the absolute force CD in a perpendicular direction, is to the relative force in an oblique direction CB, as CA is to CP: Now the number of particles striking the oblique plane CB, at the same time as CA is CP: the absolute force that strikes a plane CB, in a perpendicular direction, is therefore to the part that strikes it in an oblique direction, as the square of the

K 2

radius

radius CA is to the square of the sine CP of the incident angle DCL or CBP.

C O R.

131. Hence unequal planes differently inclined are acted upon by forces which are in the compound ratio of theses planes and the squares of the sines of the incident angles. For let P, Q, R, denote the forces which act upon the planes CA, CB, CE, respectively; then will $CA^2 : CP^2 :: P : Q$, by the last, and the forces acting upon inequal planes CB, CE, equally inclined, are as these planes, that is, $CB : CE :: Q : R$, and the compound of these two proportions gives $CB \times CA^2 : CE \times CP^2 :: P : R$.

P R O B L E M.

132. *If a machine be moved by means of a water-wheel, to find the velocity of that wheel, such that the machine shall produce the greatest effect possible.*

Let a denote the uniform velocity of the stream, x the velocity required: then will $a - x$ be the velocity by which the stream acts upon the wheel; and as the force of the stream is as the square of the velocity, by *Art. 131*, $\overline{a - x}^2$, represents the force acting upon the wheel, and the velocity x multiplied by the force gives $x \times \overline{a - x}^2$ for the momentum of that force, by mechanics; which must be the greatest possible: therefore $x \times \overline{a - x}^2 - 2xx \times \overline{a - x} = 0$: hence, $x = \frac{1}{3}a$, which shews that when the velocity of the wheel is one third of the velocity of the stream, the machine will produce the greatest effect possible.

C O R

133. Hence, if the value of x be wrote into the momentum $x \times \overline{a - x}^2$, we get $\frac{4}{27}a^3$ for the greatest momentum possible; and $\overline{a - x}^2$ gives $\frac{4}{9}aa$, for the weight thereof in this case; that is, if the machine

chine be so loaded as to remain just at rest, then by loading it by $\frac{1}{9}$ of that weight, it will produce the greatest effect possible; setting aside all friction or impediment.

L E M M A. Fig. 45.

134. *The sines and cosines of any two arcs of a circle being given, to find the sine and cosine of an arc equal to their sum.*

Let BE, MP, be the sines of the arcs BM, MA; BL the sine of the sum BMA, draw ED, EQ perpendicular to BL and CA: then because the angles CEQ, DEB are equal the right-angled triangles CQE, BDE are similar to the triangle CPM; and if the radius CM be unity, we have, 1°. $CM:CP::BE:BD=CP \times BE$, and $CM:CE::PM:QE=CE \times PM$. Now since the sum of BD and QE is equal to the sine BL, we have $BL=CP \times BE + CE \times PM$; that is, *the sum of the rectangles made by the sine of one arc and the cosine of the other, is equal to the sine of the arc equal to their sum.*

2°. $CM:CP::CE:CQ=CP \times CE$, and $CM:PM::BE:DE=PM \times BE$; and as the difference between CQ, DE, is equal to the cosine CL: we have $CL=\pm CP \times CE \mp PM \times BE$, according as the arc AB is less or greater than a quadrant. Hence, *the difference between the rectangles made by the cosines and by the sines of any two arcs, is equal to the cosine of an arc equal to their sum.*

P R O B L E M. Fig. 46.

135. *The angles which the wind makes with both the axis and sails of a wind mill, being given; to find the force with which the mill turns.*

Let CA be the direction of the axis, CQ that of the sails, supposed to be flat, and CM that of the wind.

wind. From the center C describe an arc AMF; draw MQ, AL perpendicular to CF and MP to CA: then because the force of the wind against the sails, being as the square of the sine QM of the incident angle FCM, by *Art.* 130; if we take CD equal to QM perpendicular to CQ, and draw DE perpendicular to CA; then will CD be to DE as the force of the wind to its part which turns the mill. Now, if the sine PM= x , the sine QM= y , and the radius unity, then will $CP = \sqrt{1 - xx}$, $CQ = \sqrt{1 - yy}$ and $CL = \sqrt{1 - yy} \times \sqrt{1 - xx - xy}$, by *Art.* 134, N^o. 2: and because CM is to QM, as the absolute force of the wind is to its effect against the sails, if the absolute force be unity, we get y^2 for the force in the direction QM or DC; and because of the similar triangles ACL, CED, we have $AC(1) : CL(\sqrt{1 - yy} \times \sqrt{1 - xx - xy}) :: QM \text{ or } CD(yy) : DE = y^2 \sqrt{1 - yy} \times \sqrt{1 - xx - xy}^3$, for the force in the direction perpendicular to the axis CA.

C O R.

136. Since then the force by which a wind-mill is turned, depends upon the sines of the angles made by the direction of the wind and position of the sails in respect to the axis, it is therefore necessary to find what these angles ought to be, to produce the greatest effect possible. But as there are two variable quantities, one must be supposed constant, and then the other; for conveniency we shall suppose first the sine y of the angle MCQ constant. Now it is plain that the force $y^2 \sqrt{1 - yy} \sqrt{1 - xx - xy}^3$, increases while x decreases; for $\sqrt{1 - xx}$ increases and $-xy^3$ decreases; and therefore the force will be the greatest possible when $x = 0$, or when the direction CM of the wind coincides with the axis CA, let the sails make what angle they will with the axis. This agrees with the common practice.

Now,

Now, the direction of the wind being supposed to coincide with the axis, the next inquiry is what angle the sails are to make with the axis: when $x=0$, the expression of the force becomes $y^2 \sqrt{1-yy}$, which ought to be the greatest possible; therefore its fluxion being made $=0$, gives $2y-2y^3=y^3$, or $y=\sqrt[3]{\frac{2}{3}}$; and the angle corresponding to this sine is $54^\circ, 44'$.

C O R.

137. If now, AC represents the direction of a ship; and CD the rudder, the rest as before, we shall draw the same conclusions as we have done in regard to the wind-mill, viz.

1°. If the direction of the ship is in that of the current; the rudder CD ought to make an angle of $54^\circ, 44'$ with the ship, to make it turn round with the greatest force possible.

2°. When the direction of the wind is in that of the ship, the sails ought to be at right-angles to them. Since then the sine QM (y) becomes equal to the radius CA; and its square, which expresses the force, will be the greatest possible.

3°. To find the angles made by the ship, wind and sails; we shall suppose y constant; and if s be the cosine of the angle QCM; then the force $yyx \sqrt{1-yy} + y^3 \sqrt{1-xx}$ becomes $yyxs + y^3 \sqrt{1-xx}$, whose fluxion being made $=0$, gives $s \sqrt{1-xx} = xy$, the square of which is $ss - ssxx = y^2 x^2$, or $ss = xx \times ss + yy$, or because $ss + yy = 1$, $ss = xx$ and $s = x$, or $s = -x$, which shews that the sine of one angle must be equal to the cosine of the other.

This gives two different cases; when the direction of the wind is between that of the ship and that of the sails, as in this figure the angle ACQ must be a right one.

But

But when the direction CM of the wind is on one side, and the sails CF on the other to that of the ship CA, as in *Fig. 47*; the sine MP(x) becomes negative, and equal to the cosine of the angle FCA.

Now, as in these two different dispositions one must be preferable to the other, we are then to find which of them, makes $yyx\sqrt{1-yy} + y^3\sqrt{1-xx}$ the greatest.

In the first case, when the angle ACF, *Fig. 46.* is a right angle, the sines x , y are both positive, and $x = \sqrt{1-yy}$, by the last, and hence the expression of the force becomes y^2 . In the second x becomes negative, and that of the force is $-yyx\sqrt{1-xx} + y^3\sqrt{1-xx}$, which by substituting the former value x becomes $2y^4 - y^2$. Hence, because y must be less than unity, the first expression y^2 of the force is greater than the second $2y^4 - y^2$; which shews that when the direction of the wind is between that of the ship and the sails, and the angle ACF is a right one, the ship sails with the greatest velocity possible; and this case $yy + xx = 1$, wherein y may be greater equal to, and less than x . And when $yy = \frac{2}{3}$, then will $xx = \frac{1}{3}$; and the angle FCA, 54° , $41'$.

The ship cannot always sail in the best manner; by making the fluxion of the force $2y^4 - y^2$ equal to 0, we get $\frac{1}{4}yy$ and so $xx = \frac{3}{4}$, gives the best position in this case.

The Chevalier *Renau* wrote first upon this subject in the year 1689; the famous *John Bernoullie* wrote against him in 1714, pretending that the Chevalier had been mistaken in some cases. But, according to what has been said here, both of them seem to have been mistaken in some cases; however, their works are so full of mathematical learning as to deserve being perused by the curious.

PROBLEM

P R O B L E M. *Fig. 48.*

138. *Let a crooked tube ABCD of any diameter be filled with water to the height LK; to find the time and velocity of the descent and ascent of the water.*

Let d denote the mean length of the tube filled with water, x the length of descent in one branch or ascent in the other, v the velocity, and t the time; then will $x = vt$, by *Art. 81*; and if m, n express the sines of the angles ABE, DCF of inclination of the branches, the radius being unity: then the whole length d is to x as the weight of the water (1) is to $\frac{x}{d}$ the weight in the direction of the branches, and the radius is to the sine m or n as the weight $\frac{x}{d}$ in the direction of the branches to their effects $\frac{mx}{d}, \frac{nx}{d}$, in the vertical direction. Whence the force of acceleration, equal to the difference between the quantities of water in both branches, will be $\frac{m+n}{d}x$, and $\frac{m+n}{d}x \times x = v\dot{v}$, by *Art. 82*; whose fluent is $\frac{m+n}{d} \times A - xx = d\dot{v}v$. Now if x becomes a when $v = 0$; then $A = aa$, and if $aa - xx = yy$, the fluent becomes $m + n \times yy = d\dot{v}v$, or $y \sqrt{m+n} = nvd^{\frac{1}{2}}$, this value of v wrote into $x = vt$, gives $xd^{\frac{1}{2}} = yi \sqrt{m+n}$. Whence if z denotes an arc of a circle whose radius is a and sine y , we get $zy = ax$, by *Art. 4. N^o. 2*, and so $zd^{\frac{1}{2}} = at \sqrt{m+n}$, whose fluent is $zd^{\frac{1}{2}} = at \sqrt{m+n}$; and when z denotes the semicircumference, the ratio of z to a is constant, and the time t that of the whole descent, which therefore is always the same, let that descent be what it will.

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C O R.

C O R.

139. Since the oscillations of pendulums are in a subduplicate ratio to that of their lengths, by *Art.* 105, a pendulum, whose length is $\frac{d}{m+n}$, will perform its vibrations in the same time, that the water ascends and descends in the tube.

P R O B L E M. *Fig.* 49.

140. *The density of the air being every where as the force of compression, to find its density at any given height from the surface of the earth.*

Let BMT be the logarithmical curve, AV its asymptote, AB the surface of the earth, and PM any ordinate. If $AB=n$, $PM=y$, $PV=u$, and a a constant quantity or the subtangent; then will $ay=yu$, by the property of the curve; and since the density is proportional to the pressure, yu or it equal ay , will express the fluxion of the density, and therefore ay the density required.

C O R.

141. If $AP=x$, then as the fluent of $ay=yu$, is $a \times \log. y = u$, by *Art.* 57; and when $y=n$, then $a \times \log. n = AV$; by subtracting the first equality from the last we get $a \times \log. n - a \times \log. y = AV - PV = x$, or $x = a \log. \frac{n}{y}$. Therefore the density n at the surface being given, the height x which answers to any other density y may be determined.

C O R.

142. Hence if d expresses the density of quicksilver, and b the height of a column of quicksilver equal to the weight of the atmosphere, standing upon the same base; then as an expresses the weight of a column of air at the surface, we have $an=db$, and therefore,

therefore, the specific gravity of the air and quicksilver being given together, with the height b of the mercury, the height of a column of air uniformly dense with that at the surface, which is equal to the weight of the atmosphere, will also be given.

C O R.

143 Hence, the weight of the atmosphere increases or decreases as the mercury rises or falls in the glass; which never rises above thirty-one inches, nor falls lower than twenty-eight: these two numbers denote the extremes of the atmosphere's weight. Mr. Cotes affirms the quicksilver to be 11900 heavier than air, or 14000 to 1; and Sir Isaac Newton 11890; whence if we suppose $b=31$, $d=11900$ and $n=1$, we get $an=db=31 \times 11900=368900$ inches; or $a=30742$ feet.

We have here considered gravity as uniform; but when the particles are attracted like other bodies, the problem will be different, in the manner following: Let C represent the center of the earth, $CA=r$

the radius, and the rest as before: then will $\frac{-\dot{x}y}{r+x)^2}$ be the fluxion of the density or pressure at P by supposition, and if in CP we take $\div CP:CA:CQ=z$;

then $r+x=\frac{rr}{z}$, whose fluxion is $-\dot{x}=\frac{rr\dot{z}}{zz}$: Now the values of $\dot{x}$ and $r+x$ being wrote into the fluxion of

the pressure, gives $\frac{y\dot{z}}{rr}$. Whence, if the curve BNR

be such that the ordinate QN be always equal to the corresponding one, PM, the area CQNR divided by the square of the radius r , will express the pressure at P. Because the density PM at P is as the pressure above it, its fluxion $\dot{y}$ is in a constant ratio to the fluxion

$\frac{-\dot{x}y}{(r+x)^2}$ or to its equal $\frac{y\dot{z}}{rr}$: and since $\dot{y}$ is to $\frac{y\dot{z}}{rr}$ as $\frac{\dot{y}}{y}$ to

$$\frac{\dot{z}}{rr}$$

$\frac{z}{rr}$ by division; the logarithm of y will be to z in a given ratio; the curve BNR is therefore a logarithmic reversed, and equal to BMT; and since AB is a common ordinate to them both, they must also have the same subtangent a : consequently $ay = yz$, and $\frac{ay}{rr} = \frac{yz}{rr}$, whose fluent $\frac{ay}{y}$ gives $\frac{an}{rr}$ for the pressure at the surface A, and $ay = yz$ gives $aly = z$ and $aln = r$. Therefore $aln - aly$ or $al\frac{n}{y} = r - z$: but as $z = \frac{rr}{r+x}$, we get $al\frac{n}{y} = \frac{rx}{r+x}$: Now since x is but very small in comparison to the radius r of the earth, the attraction can cause no sensible error in the distance AP.

P R O B L E M. Fig. 50.

144. *The density and elastic force of a medium being given; to find the velocity and time of a pulse produced by a tremulous body in that medium.*

Suppose the medium placed in AC, to be moved from C towards A by the pulse, whose extreme distance is AC, so that when the pulse ceases, the elasticity of the medium causes the particles of the medium to return to their former situations, and when the pulse moves backwards, it will move the medium on the other side of the point C, at an equal distance CB to the former CA, and then ceases again; so that the pulse will continually move the particles of the medium forwards from C to A, and backwards from C to B, as long as they continue: but while the medium placed in AC is moved from C to A by the pulse, the medium placed in CB, will expand itself by its elastic force from B to A, and when the medium in BC is moved from C to B by the pulse,

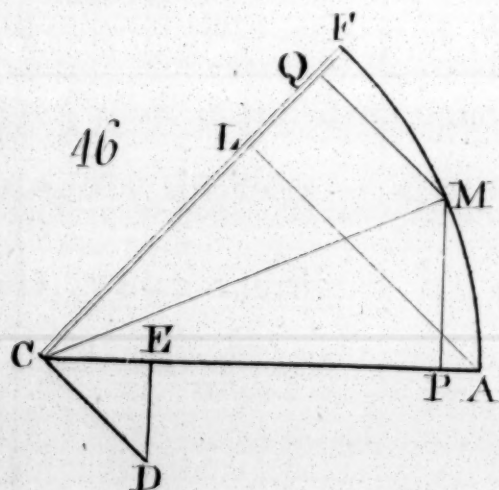
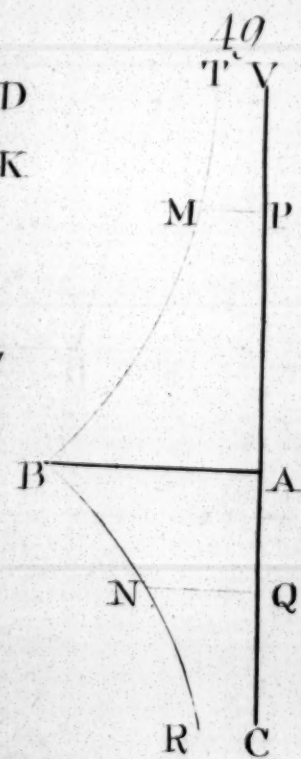
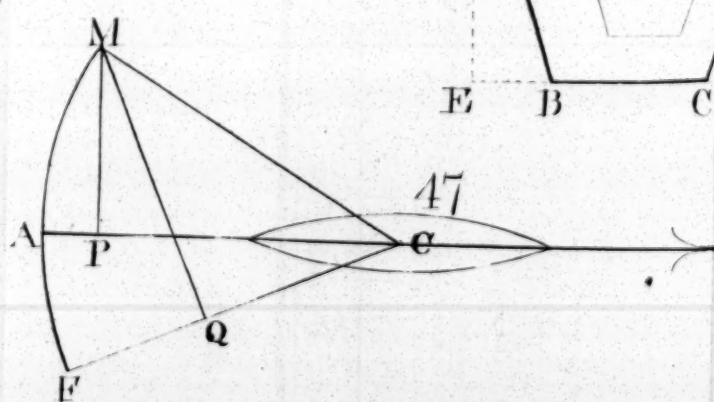
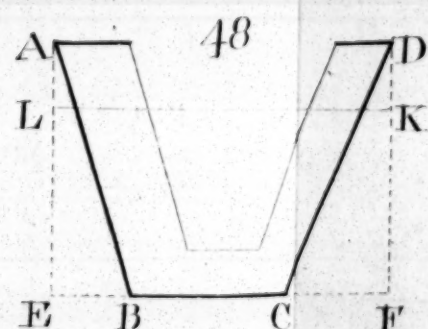
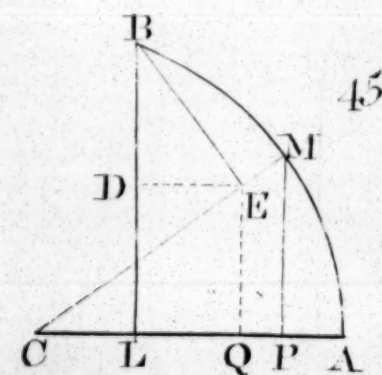
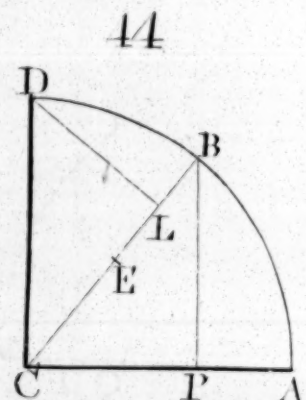
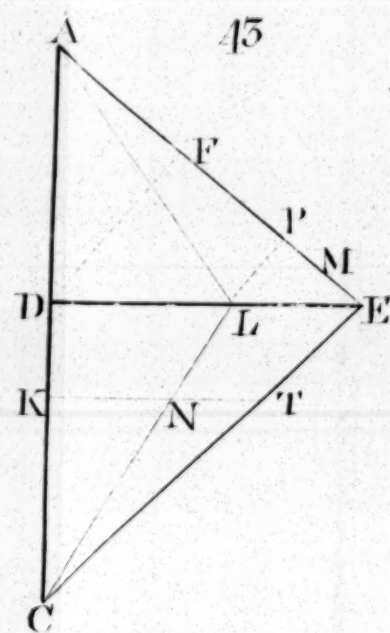
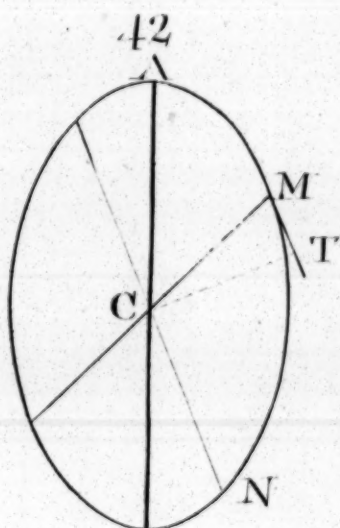
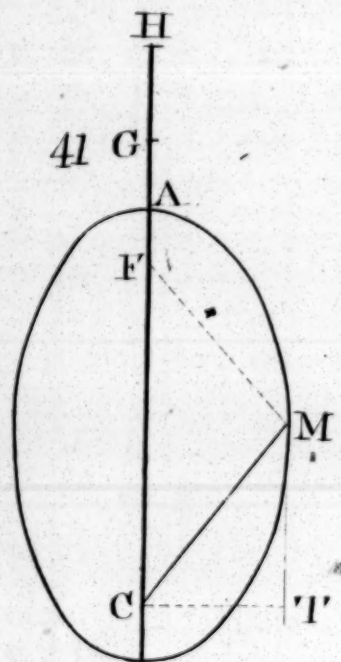
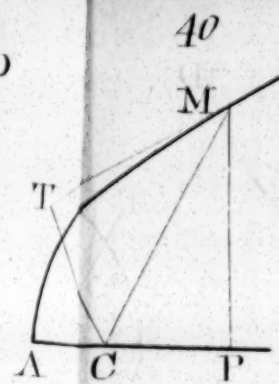
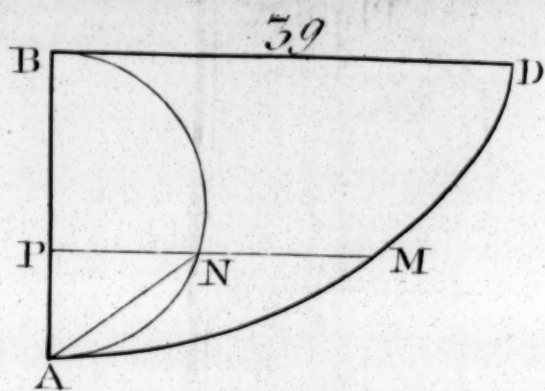
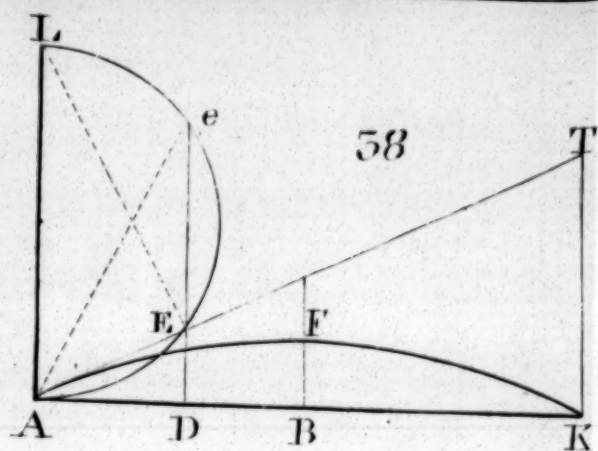
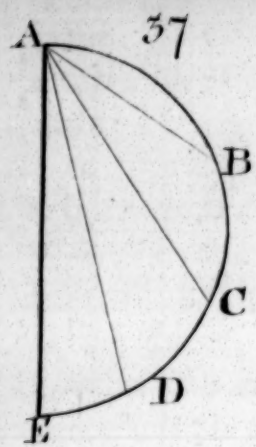


Fig. 1

Fig. 2

Fig. 3

Fig. 4

Fig. 5

Fig. 6

Fig. 7

Fig. 8

Fig. 9

Fig. 10

Fig. 11

Fig. 12

Fig. 13

Fig. 14

Fig. 15

Fig. 16

Fig. 17

Fig. 18

pulse, that in AC expands itself from A to B by its elastic force: and for the same reason, when the medium in CA is moved by the pulse from C to any intermediate point P, the medium in CB expands itself from B to P; the accelerative force in any point P is therefore as the difference between the forces of expansions from A to P and from B to P; or which is the same, between the forces of the pulse that moves the medium from C to P, and the elastic force which dilates the medium BC from B to P.

If A denotes the height of the medium of an uniform density, equal to that in which the pulses are made, $CA=a$, $CP=x$, v the velocity, and t the time of motion from C to P; then as the elastic force is as the density and altitude of the medium by supposition, the rectangle aA will express the elasticity of the medium in its natural state, and this force is increased in the motion from C to P reciprocally as AP the space into which the medium AC is reduced; that is, the force at P in its return from A to P will

be as $\frac{aA}{a-x}$, and for the same reason, the force of expansion of the medium in BC, to BP, will be as $\frac{aA}{a+x}$; and their difference $\frac{aA}{a-x} - \frac{aA}{a+x}$ or $\frac{2axA}{aa-xx}$, will express the accelerative force at P: but on account of the extreme shortness of the pulses, xx may be neglected in the denominator without any sensible error, and therefore $\frac{2x}{a}A$ expresses the accelerative

force, which multiplied by $\dot{x}$, gives $\frac{Axx\dot{x}}{a} = v\dot{v}$,

by Art. 82, whose fluent is $B - \frac{2Axx}{a} = vv$; but at

the

the point A, where $v=0$ and $x=a$, we get $B=2Aa$, and so $\frac{2A}{a} \times aa - xx = vv$.

Now, if from the center C, a circle be described thro' the points A and B, and the ordinate PM drawn, and we make $PM=y$, the arc $AM=z$; then $yy=aa-xx$, and $\frac{2A}{a}yy=vv$, or $y \sqrt{\frac{2A}{a}}=v$, and this value of v wrote into $\dot{x}=vi$, *Art.* 81, gives $\dot{x}=i y \sqrt{\frac{2A}{a}}$, or because $\dot{z}y=a\dot{x}$, by *Art.* 4, N°. 2, we get $\dot{z}=i \sqrt{2Aa}$, whose fluent is $z=t \sqrt{2aA}$; and when z becomes the semicircumference AMB, t expresses the time of a pulse going from C to A and returning from A to C, which is therefore constant, let the extent of the pulses be what it will.

C O R.

145. Since $\frac{2A}{a}yy=vv$; when $y=a$, at the point C, we get $2aA=vv$, or $\sqrt{2aA}=v$. Now, if b expresses the height thro' which a body falls in a second then because the times of falling bodies are as the square roots of the spaces fallen thro', by *Art.* 84, N°. 2, we have $\sqrt{b}$ to $\sqrt{a}$, as one second is to $\frac{\sqrt{a}}{b}$ the time fallen thro' a expressed in seconds; and the space $\sqrt{2aA}$ divided by the time $\frac{\sqrt{a}}{b}$, gives $\sqrt{2bA}=v$ for the space moved over uniformly in a second, with the velocity v .

C O R.

146. It is manifest that velocity of pulses, or the spaces uniformly described, in a second, are as the square root of the altitude A, in the same latitude; or the

the velocity v is also in a subduplicate of the elastic force directly and density inversely, since Aa expresses the elastic force and a the density, by *Art.* 144.

147. The number of pulses propagated, is always equal to the number of vibrations of the tremulous body and no more. For when the particle at C is moved by the pulse from C to A , it will return again to C by its elastic force, where it will remain unless it receives a new motion either from the impulse of the tremulous body, or from the pulses propagated by the body. As soon, therefore, as the pulses cease to be propagated, the motion likewise ceases. It may be observed, that the velocity here determined is not absolutely

true, since in $\frac{2aAxx}{aa-xx} = -v\dot{v}$, we have neglected xx

in the denominator; and the true fluent is $2aA \times b \log. \frac{aa-xx}{aa}$, and $2aA \log. aa = vv$, or $\sqrt{2aA \times \log. a} = v$; the value of v found before is however near enough in all cases where the motion is not very intense or quick.

R E M A R K.

The foregoing principles may be applied to the motion of sound, since it is excited by the tremulous pulses of bodies, as experience shews. As the altitude A has been found to be 30742 feet, by *Art.* 143, and $b=16.1$, by *Art.* 107; therefore $32.2 \times 30742 = 989892.4$, whose square root gives 994.93 feet for the distance the sound moves over in a second.

Sir *Isaac Newton* finds 979 feet for that distance, because he takes the altitude A something less than we have done. But as this does not agree with the experiments, he adds 109 feet for the thickness of the particles of the air; and as this is not yet sufficient, he supposed the motion of sound increased in the ratio of 20 to 21, by the vapours floating in the air; then by means of these additions he finds that distance to be 1142 feet.

Bat

But as these suppositions are not grounded on any facts, I rather imagine that the altitude A has not been rightly determined, and for several reasons.

1°. The density of the air is supposed to decrease from the surface upwards in the ratio of the ordinates of the logarithmic curve, but the particles near the surface being heated by the reflection of the sun's rays, whereby they are more rarified than higher up.

2°. The air being loaded with vapours more or less, must needs vary that altitude.

3°. The air being so very light in comparison to other bodies, it is extremely difficult to determine its density.

It appears therefore to me more expedient to determine the altitude A from the known velocity of the sound, than this from the other, which may be done as follows: divide the square 1304164 of 1142 feet by $2b(32.2)$, which gives $A=40502$ feet for the true medium altitude of the air, or $A=486024$ inches.

And because $an=db$, by *Art.* 142, the altitude $a(486024)$ of the atmosphere is to the altitude $b(31)$ of the quicksilver, so is the specific gravity $d(1400)$ of the specific gravity n of the air, and therefore $n=\frac{434000}{486024}$, or $\frac{7}{8}$ nearly; whereas it is marked $\frac{5}{4}$ in the tables of Mr. Cotes, and generally supposed unity.

As the velocity of sound has been determined by a great number of experiments, which were made by several eminent men, such as Dr. Hally and Derham, and different times and distances, we have great reason to believe that this velocity has been much more accurately determined, than the specific gravity by hydrostatics, because of the great difficulty there is in weighing so small a volume of so very light a substance.

What confirms this determination beyond all exception, are the several accurate experiments made at

Minorca,

Minorca, by General *Williamson*, and several other officers of artillery, in order to determine the best length of guns; and we shall prove hereafter, in *Art.* 182, that the specific gravity of air cannot be unity, much less greater, if these experiments can be depended upon, of which there is no doubt, because I had it from most of these gentlemen, besides the General himself.

P R O B L E M. *Fig.* 51.

148. *The dimensions of a cylindric vessel, ABCD, constantly full of water, being given, as well as the diameter LM of the orifice in the bottom, to find the quantity of water that will run out of it in a given time.*

Let z express the quantity of water in the vessel, v the distance of its center of gravity from the bottom; then zv will express its momentum against the orifice, by hydrostatics, whose fluxion is $\dot{z}v + z\dot{v}$. Now if the height $AB = x$, the surface $BC = m$, that of the orifice $LM = n$; then will $mx = z$, by supposition, and the fluxion $m\dot{x}$ multiplied by the distance x gives $mxx\dot{x}$ for the fluxion of the momentum likewise. Therefore $\dot{z}v + z\dot{v} = mxx\dot{x}$. But when the surface BC descends to EF , it is plain that the quantity of water run out at the orifice must be equal to that contained in the space $EBCF$; and because the velocities of equal quantities of water thro' different openings, at the same time, are as these openings inversely, by hydraulics: if $a + x$ denotes the height thro' which a falling body acquires the velocity of the water at the orifice, then will $mm : nn ::$

$a + x : \frac{nn}{mm} \times a + x =$ to the height thro' which a fal-

ling body would acquire the velocity of the water passing thro' the section BC . And since the momentum of the pressure in any other section, is equal to

M

the

the difference of the momentums at BC and LM, the fluxion $\frac{nn\dot{x}}{mm}$ of the height $\frac{nn}{mm} \times a + x$. Subtracted from

the fluxion $\dot{x}$ of $a + x$, gives $\dot{x} - \frac{nn\dot{x}}{mm}$ for the height

thro' which a falling body would acquire the velocity of the descending water, and this multi-

plied by the section BC(m) gives $m\dot{x} - \frac{nn\dot{x}}{m}$ for the

fluxion of the quantity run out in the time of the descent thro' BA, and since $\dot{z}$ expresses likewise the

fluxion of that quantity, we have $\dot{z} = m\dot{x} - \frac{nn\dot{x}}{m}$; now

this value of $\dot{z}$ and $m\dot{x}$ that of z found above being

wrote into $\dot{z}v + z\dot{v} = m\dot{x}x$, gives $v\dot{x} - \frac{nn}{mm}v\dot{x} + \dot{z}v = x\dot{x}$,

after having divided by m , or if we make $1 - \frac{nn}{mm} = r$,

we shall have $v\dot{x}r + \dot{z}v = x\dot{x}$.

To find the fluent of this equation, multiply both sides by x^{r-1} , then will $rv\dot{x}x^{r-1} + \dot{z}vx^r = \dot{x}x^r$, whose fluent is $vx^r = \frac{x^r + 1}{1+r}$, or dividing by x^r and restoring

the value of r , we get $v = \frac{mmx}{2mm - nn}$.

When the vessel is not kept constantly full, that is, when it is not supplied by fresh water as it runs out, the descent of water will no more be uniform as before, but accelerated, and therefore, the height thro' which a falling body acquires the velocity, will be less at BC in proportion to that at the orifice; that is,

we must from $2\dot{x}$ subtract $\frac{nn}{mm}\dot{x}$, in order to get

$2\dot{x} -$

$2\dot{x} - \frac{nn}{mm}\dot{x}$ for the medium pressure, and hence $2m\dot{x} - \frac{nn}{m}\dot{x} = \dot{z}$; by following the same steps as above, we shall have $v = \frac{mmx}{3mm - nn}$, in this case.

This is the demonstration we promised in our elements of mathematics, page 273, it differs very little from what Sir *Isaac Newton* says in his *Principia*, Book II. first Cor. after Prop. 36, but answers exactly his correction in that Prop. *Daniel Bernoullie* has treated this subject very copiously in his *Hydrodynamica*, but what he says in page 37 is quite a mistake.

C O R.

149. Hence if KP, be perpendicular to the middle P, of the orifice, and equal to twice the height thro' which a falling body acquires the velocity of the water issuing thro' the hole, and the curve BL, CM, be such that $KI : KP :: PL^4 : IB^4$; the velocity at I will be to the velocity at P, as KI is to KP: by *Art. 559* of our *Elem.* and by the revolution of this curve about the axis KP, forms the cataract BLMC of the water, whose property is such that its content expresses the quantity of water which presses upon the orifice in the same manner as if the rest was congealed into ice. And therefore the other part of the water contained in the solid described by MCD, expresses the quantity which presses upon the ring described by MD, in the rotation.

PROBLEM. Fig. 52.

150. If in the middle of the orifice LM, there was placed a small circle ab , described about the center d , to find the quantity of water acb , supported by that circle.

Imagine the figure 51 to revolve about CD, so that MD and CD, in this figure represent the radius ad , and altitude of the solid acd , supported by the small circle ab ; and call $KI=a$, $IP=x$, $IC=b$, $PM=y$, $MD=z$; then will $y=b-z$; and by the property of the curve, we have $a:a+x::y^4:b^4$, or $\frac{ab^4}{y^4}=a+x$, or because $y=b-z$, $x=\frac{ab^4}{(b-z)^4}-a$,

whose fluxion is $\dot{x}=\frac{4a\dot{z}b^4}{(b-z)^5}$; therefore the fluxion of

the solid described by MCD, will be as $\dot{x}xz$, that is as $\frac{4a\dot{z}z^2b^4}{(b-z)^5}$, whose fluent is $\frac{abb}{3} \times \frac{4bz^3-z^4}{(b-z)^4}$.

Since the fluxion of this fluent is the same as the proposed one. And as xzz expresses the cylinder of the same base and altitude as the solid acb ; and

$x=\frac{ab^4}{(b-z)^4}-a$, that cylinder will be $azz \times \frac{b^4}{(b-z)^4}-1$:

consequently the solid is to the cylinder as $\frac{abb}{3} \times$

$\frac{4bz^3-z^4}{(b-z)^4}$ to $azz \times \frac{b^4}{(b-z)^4}-1$; or as $\frac{bb}{3} \times \frac{4bz^3-z^4}{4bz^3-z^4}$

to $zz \times \frac{b^4-b-z^4}{b^4-b-z^4}$: but $b^4-b-z^4=4b^3z^3-6bbz^4+4bz^5-z^6$; and so the former proportion becomes that of $\frac{4}{3}b^3-\frac{1}{3}bbz$ to $4b^3-6bbz+4bz^2-z^3$, when divided by z^3 .

Hence, if z becomes so small in comparison to b , that all the terms multiplied by that quantity may be neglected

glected, without any sensible error, the solid acb will be to the cylinder of the same base and altitude as $\frac{4}{3}b^3$ to $4b^3$, or as $\frac{1}{3}$ is to unity.

Sir *Isaac Newton*, finds this last ratio to be as $\frac{1}{2}$ to unity, in *Cor. 9.* after *Prop. 36. Book II.* What induced him to think so (for he did not compute it as we have done) was, the solid acb must be less than a semi-spheroid, because the angle c is acute, whereas in the spheroid it must be a right one, and the angle at a must also be acute; and as one half is an arithmetical mean between the cylinder and cone, he concludes, that this solid is one half of the cylinder. The reader may easily perceive the error of this conclusion; for tho' the former part of his argument, that this solid must be less than a semi-spheroid is true, he has by no means proved or attempted it, that this solid must be greater than the cone, whereas we have clearly proved that the solid and the cone must be equal, upon the supposition that z is incomparably less than b .

C O R.

151. Since we have $v = \frac{mmx}{2mm - nn}$, by *Art. 148.*

when the orifice LM becomes so great as to be sensibly equal to BC , which will happen when the circle ab is placed in an unconfined fluid: then n may be taken equal to m ; and hence the equality above becomes $v = x$; that is the velocity of the fluid is equal to that which a body acquires by falling thro' the height cd of the fluid.

C O R.

152. Hence, if the water was at rest, and the little circle ab , did ascend, with the same velocity with which the water descended thro' the orifice, or that acquired by a body falling thro' the height cd ; the resistance of the little circle ab will be equal to the weight

weight it supports when at rest, and the water descends. Therefore the resistance of the little circle in the ascent will be equal to one third of the weight of the cylinder of the same base and altitude to cd .

C O R.

153. Since the solid acb is one third of the cylinder of the same base and altitude, by *Art.* 150. the circle ab must move with triple the velocity of the water in its descent, at the same time, in order to meet with a resistance equal to the weight of the cylinder of the same base and altitude cd : and since the heights of falling bodies are as the squares of the velocities acquired in the fall, by *Art.* 84. No. 2. the height thro' which a falling body acquires a velocity triple that of another, must be nine times greater. Consequently a body must fall thro' a space nine times the height cd , to acquire the velocity with which the circle ab must ascend, in order to meet with a resistance equal to the weight of a cylinder of water of the same base ab , and altitude cd . Because the resistance of a cylinder, moving in the direction of its axis, is the same as that of its base: a cylinder whose base is equal to the little circle ab , moving in the same direction and with the same velocity, will meet with the same resistance as that circle.

If the length of the cylinder, moving in a fluid, be increased or diminished, its motion, as well as the time of describing nine times its length, will be increased or diminished in the same ratio; and therefore the force by which the motion so increased or diminished, can be destroyed or generated, will continue the same; because the time is increased or diminished in the same proportion; and consequently the force remains still equal to the resistance of the cylinder, since that resistance will also remain the same.

If

If the density of the cylinder be increased or diminished, its motion, as well as the force by which its motion may be generated or destroyed in the same time, will be increased or diminished, in the same ratio. The resistance of any cylinder whatever, will consequently be to the force by which its whole motion may be generated or destroyed, in the time it moves over nine times its length as the density of the medium is to the density of the cylinder nearly.

Sir *Isaac Newton* observes, that a fluid must be compressed to become continued; it must be continued and non-elastic, that all the pressure arising from its compression may be propagated in an instant, and so acting equally upon all the parts of the body moved, may produce no change in the resistance. The pressure arising from the motion of the body, is spent in generating a motion in the parts of the fluid, and this creates the resistance. But the lateral pressure arising from the compression of the fluid, be it ever so forcible, if it is propagated in an instant, generates no motion in the parts of a continued fluid, nor produces no increase or decrease in the resistance.

L E M M A.

154. *The resistance of a body moving in a compressed non-elastic fluid, will always be the same, let its figure be what it will, when the greatest transverse section is the same.*

For whether *ab* be the diameter of a circle, or the section of any surface, the solid *acb* of the fluid supported by that surface, will be always one third part of the cylinder of the same base or altitude as has been proved in *Art.* 150. Consequently the resistance of any solid is always equal to that of its transverse section.

N. B.

N. B. We suppose the friction the solid meets with in its motion to be inconsiderable, and the particles of the fluid to have no cohesion.

T H E O R E M.

155. *If a globe moves uniformly forward, in an infinite compressed non-elastic medium, its resistance will be to the force that may generate or destroy its whole motion in the time it moves over the length of six of its diameters, as the density of the fluid is to the density of the globe.*

For since the resistance of the globe is equal to that of its circumscribed cylinder, and the resistance of the cylinder is to the force that may generate or destroy its whole motion, in the time it describes nine diameters, as the density of the medium is to the density of the cylinder, by *Art. 153*. Therefore as the globe is the two thirds of the cylinder, the resistance of the globe is to the force that may generate or destroy its whole motion, in the time it describes six of its diameters, as the density of the medium is to the density of the globe.

C O R.

156. Hence the resistances of globes, in an infinite compressed non-elastic mediums, are in the compound ratio of the squares of their velocities, of the squares of their diameters and the densities of the mediums directly, and the densities of the globes inversely. For the resistance increases in proportion as the squares of the velocities increase, it increases also as the surface of the globe, or the square of its diameter increases, as likewise as the density of the medium increases, and decreases as the density of the globe increases.

C O R.

C O R.

157. If different globes move in the same medium, and with the same velocity, their resistances will be as the squares of their diameters directly, and their weights inversely; and if the globes are of the same density, their resistances are as the squares of their diameters directly, and the cubes of the same diameters inversely; that is as their diameters inversely.

Thus the resistance of a nine pound iron ball, whose diameter is four inches, is to the resistance of an iron ball, whose diameter is seven inches, as seven to four.

C O R.

158. The greatest velocity a globe can acquire in a resisting medium by its comparative weight, is equal to that which it may acquire by falling with the same weight, from a state of rest in a non-resisting medium, thro' a space that is to three of its diameters as the density of the globe is to the density of the medium. For the globe, in the time of its fall, will acquire a velocity, which, being uniformly continued, will describe a space double that fallen thro' in that time. And therefore the resistance will then be equal to its gravity, by *Art.* 153: and after this the motion can no more accelerate and become uniform.

C O R.

159. Hence if d denotes the diameter of the globe, m its density, n that of the medium, and r twice the height from which a falling body acquires the greatest velocity; then will $r : 6d :: m : n$, or $r = \frac{6md}{n}$, by

Art. 158.

N

R E.

R E M A R K I.

Sir *Isaac Newton* found $r = \frac{8md}{3n}$, in the 2d *Cor.* after *Prop.* 38, Book II. in consequence of his finding the weight supported by a small circle placed in the middle of an orifice, to be one half; whereas we proved it to be but one third of the cylinder, which has the same base, and whose altitude is equal to the height thro' which a falling body may acquire the velocity of the water. Now as the laws of motion in a resisting medium absolutely depend on the determination of the greatest velocity of bodies, it will be shewn hereafter, in *Art.* 179, that our determination of this velocity is confirmed by some undoubted experiments, and that Sir *Isaac's* cannot be true.

R E M A R K II.

Sir *Isaac Newton* considers the air to be an infinite compressed fluid of the same class as water, which does not seem to agree with its compressibility and elasticity, but rather as a rare elastic medium. In this case, a globe moving forward uniformly meets with a resistance that is to the force by which its whole motion may be generated or destroyed, in the time it describes two thirds of its diameter, as the density of the medium is to the density of the globe. This Sir *Isaac* has demonstrated in *Prop.* 35, Book II. and I in my *Artillery, Theor.* 5, of the Introduction. But it is supposed, in this last cited proposition, that the particles of the medium vanish immediately after the stroke, so as not to prevent the next from striking the body in the same manner: this certainly is not the case in the air, which compresses in a heap before the body, and forms a kind of pointed

ed solid, and if it is nearly the same as we have shewn it does in water, the globe must then move over nine times the two-thirds of its diameter, that is six of its diameters, by *Art.* 150, to meet with a resistance that is to the force which may generate or destroy its whole motion, in the same time as the density of the medium is to the density of the globe; and consequently the same as above. Hence, whether the air be considered as an infinite compressed fluid, or as a rare elastic medium, the conclusion will be the same.

SECTION VII.

Of MOTION *in a* RESISTING MEDIUM.

THE laws of motion considered in a non-resisting medium, and which we have treated of in the last section, have formerly been looked upon as universal, and were applied to all cases in general, till *Sir Isaac Newton* found that they could only hold good in those bodies which moved at a great distance from the earth, such as the planets and their satellites, which are supposed to describe their orbits in a free space, void of all resistance: it is true, that no absolute vacuum can be conceived, yet as the medium in which the heavenly bodies move, must be infinitely rare, so as to cause no retardation perceivable in many ages, it may be neglected. But in the motion of bodies near the surface of the earth, the case is quite different; since the particles of air resist the bodies considerably, and therefore all that had been said before of projectile motions is erroneous and useless.

To discover the laws by which bodies are resisted, *Sir Isaac* made a great many experiments with pendu-

ums, moving in the air and in water; he did likewise let fall several globes of different densities from the cupola of St. *Paul's*, and compared the times of their fall with the times they would fall in a non-resisting medium, by which he was able to compute their resistance. It is true that these balls were too small, and the distances fallen thro' too little, to determine the resistance in a very exact manner, yet they were sufficient to confirm the theory he had established before hand, in relation to the motion when the bodies were resisted by a medium.

In the theory of powder, we are obliged to suppose that when it is fired, it acts in the same manner as an elastic fluid, because we have no other principles to go by; but as powder does not light instantaneously, and we do not know by what law the inflammation proceeds, it is evident it cannot act with the same force as an elastic fluid of the same density, and consequently, all the conclusions drawn from the supposition, are uncertain and erroneous.

It has hitherto been supposed that the more powder a piece was loaded with, the longer it should be, in order to give time for the powder to light; and it seems to be reasonable that it should be so, but notwithstanding experiments have proved the contrary; for if a piece be loaded with various charges, it will be found that there is one which will produce a greater effect than any other that is more or less: again, if several pieces of the same caliber, and of different lengths, are tried, it will be found that there is one certain length which produces a greater effect than any other.

It has likewise been found by experiment, that a mortar loaded with a quantity of powder, which weighs but one sixth part, the weight of its shell when filled, carried its shell farther than a gun loaded
ever

ever so much, when both are elevated at an angle of the same degrees: this alone is sufficient to shew our ignorance in regard to the force of powder; for as the bore of a mortar is considerably wider than the chamber, it seems that as soon as the shell has moved from its place, the powder can act but very little upon it, and therefore the length of a piece seems to contribute but little to increase the force of powder; yet in experiments made with guns, it is found that there is a certain length, which is necessary to produce the greatest effect. This will be proved hereafter, when we treat of projectile curves described in the air, and confirmed by unexceptionable experiments.

Mr. *Robins*, in his treatise of gunnery, tried by a great many experiments to determine the force of powder, but finds the velocities in general much greater than they possibly can be; which seems to be owing to the very small quantities of powder that he was able to use, for the least error or oversight in so small a quantity, must necessarily become very considerable in such large charges as are used in pieces of artillery. These experiments were intended to confirm his theory, which, in itself, was built upon erroneous principles, as we have shewn in our treatise of artillery; and it required no small degree of dexterity to confirm a theory built upon wrong principles, by experiments no less faulty in their contrivance.

The learned professor *Euler*, in his Comment upon *Robins's* work, page 566, gives a table of velocities, produced by pieces of different lengths, and loaded with various quantities of powder; wherein he committed the same error as his author, by making them much greater than they possibly can be; besides he supposes that all cannon of the same number of calibers long, and loaded proportionally, will carry their
shots

shots equally far, which we have proved not to be true; neither does he seem to take notice of Sir *Isaac's* theorem, where he shews how to find the greatest velocities; nor has he given any experiments to confirm his assertions. Thus, for instance, when the piece is twenty-two calibers long, and loaded with half the weight of the shot he finds 1477 Rhinland feet, or 1521 *English*, moved over in a second; whereas we prove that this velocity can never amount to 914.7 feet in a forty-two pounder, and consequently much less in a smaller caliber.

He gives likewise another table in page 600, which contains the greatest and best charges of pieces, whose lengths are expressed in calibers; thus when the piece is twenty-two calibers long, he finds the greatest charge to be 5.96 calibers in length; now as in our guns, four calibers in length contain an equal weight to that of the shot, therefore the best charge of an eighteen pounder is 26.82 lb. According to this table, that of a twenty-four pounder 33.96 lb. and that of a forty-two, is 62.58 lb. Now as it has been found by experiment, that nine pound of powder in an eighteen pounder of about this length was the best charge; and therefore 12 lb. for a twenty-four pounder, and 21 lb. for a forty-two pounder: the charges given in this table are above double those they ought to have. Besides if these charges were true, they could never be put in practice, since these pieces would become twice heavier than they are now to resist so great a force as would be produced by these charges. As to the charges given to battering pieces, it has been found by experiments, that one third of the weight of the shot is sufficient; in all other occasions, a less quantity will do.

In

In this section, after having given some general theorems, and a general problem of projectile curves, we give several examples of bodies moving in right lines, from whence we deduce some useful remarks relating to artillery, in order to determine the best length and charge of fire arms; this being done we treat of the curve described in the air, which is the most difficult part of this work: As no author has yet determined this curve in all cases, we have been so much the more attentive, and given various numeral examples, by three different methods, in such a manner as to leave no doubt or obscurity, to prevent any misunderstanding of the subject.

T H E O R E M. *Fig. 53.*

100: *The fluxion of the velocity generated by any power in a resisting medium, is as the rectangle made by the difference between the measure of the power and the resistance, and the fluxion of the time elapsed from the beginning of the motion.*

Let the ordinate PM, of the curve AM constantly express the measure of the power, and the ordinate PN of another curve AN the corresponding resistance of the body; the abscissa AP expressing the time from the beginning of the motion: it is evident, that the difference MN, expresses the force which generates the velocity: for the area AMP, expresses the sum of all the powers from the beginning of the motion, and the area ANP the sum of all the resistances, and therefore the difference NAM between these areas the velocity at the end of the time AP: but the fluxion of this area is equal to the rectangle made by NM and the fluxion of AP, by *Art. 16*; consequently the fluxion of the velocity is equal to the rectangle made by the difference NM, between the measure of the power and the resistance and the fluxion of the time elapsed.

THEO.

THEOREM.

161. *The fluxion of the space described in a resisting medium, is as the rectangle made by the velocity and the fluxion of the time elapsed from the beginning of the motion.*

For since the space MAN is as the velocity, and the fluxion of the space described as the rectangle made by the velocity and the fluxion of the time, by *Art. 81*; the resistance having no other effect than to diminish the velocity, the demonstration is here the same as in that article.

THEOREM.

162. *The resistance acts only in the direction of the moving body.*

For in all other directions the body is pressed or acted upon, equally and in opposite directions; the resistance can therefore affect the motion no where but in the direction of the body in motion.

GENERAL PROBLEM. *Fig. 54.*

163. *If a body, projected from a given point A in a direction AE, with a given velocity, be acted upon by any centripetal force, in a resisting medium; to find the nature of the curve AM described by the body.*

Let $CM=y$, $MT=z$, the perpendicular TR to CM, $\dot{x}$, $RM=\dot{y}$, the perpendicular CS to MT, s ; p , the centripetal force; R the resistance; v , the velocity, and t the time: then will $\dot{z}=vt$, *Art. 161*; and as $\frac{\dot{y}}{z}p$ expresses the part of the centripetal force p which acts in the direction of the tangent, we have $\dot{v}=\pm\frac{\dot{y}}{z}p-R\times byi$, by *Art. 160*;

or

or writing the value of $\dot{t}$, it becomes $v\dot{v} = \pm p\dot{y} - R\dot{z}$; the sign $+p\dot{y}$ prevails in the descent, and $-p\dot{y}$ in the ascent. Consequently the laws of the centripetal force p and of the resistance R being given, the nature of the curve AM may be determined.

C O R.

164. When the directions of the centripetal force p are parallel, the last equation remains the same, only $\dot{y}$ expresses the fluxion of the ordinate and $\dot{x}$ that of the abscissa; we have besides $p\dot{z}^2 + vv\ddot{y} = 0$, by *Art.* 92. when $\dot{x}$ is constant, and $p\dot{x}\dot{y} = vv\ddot{x}$ when $\dot{z}$ is constant; since the resistance of the medium does not affect the equation, which does arise from the direction perpendicular to the tangent, by *Art.* 162.

C O R.

165. If r denotes twice the height from which a falling body acquires a velocity equal to the greatest which that body can acquire in the resisting medium, we have $R = \frac{vv}{r}$; for when v becomes the greatest velocity or $r = vv$, the resistance becomes equal to the force of gravity, as has been shewn before in *Art.* 158^a; and after that the motion becomes uniform in the descent, or diminishes in any other direction.

E X A M P L E.

166. Let a body descend from a vertical height by its own gravity, in a medium that resists according to the squares of the velocities; then as $z = y$ in this case, and $R = \frac{vv}{r}$, the equation $v\dot{v} = \pm p\dot{y} - R\dot{z}$ becomes $v\dot{v} = \dot{y} - \frac{vv}{r}\dot{y}$, or $\dot{y} = \frac{rv\dot{v}}{r - vv}$, supposing gravity to be unity: whose fluent is $y = A - \frac{1}{2}r \log.$

$\frac{r}{r-vv}$; but when $y=0$, then will also be $v=0$; and therefore $A=\frac{1}{2}r \log. r$; consequently the compleat fluent is $y=\frac{1}{2}r \log. \frac{r+v}{r-vv}$. Or if q expresses the number whose hyperbolic logarithm is $\frac{2y}{r}$; then will $q=$

$$\frac{r+v}{r-vv}, \text{ from whence we get } vv = q - 1 \times \frac{r}{q}.$$

Because $\dot{y}=\dot{z}$, $\dot{z}=vi$, becomes $\dot{y}=vi$, and since $\dot{y}=\frac{rv\dot{v}}{r-vv}$, we get $i=\frac{r\dot{v}}{r-vv}$. If we suppose a the greatest velocity, that is $r=aa$, the fluent of this fluxion will be $t=\frac{1}{2}a \log. \frac{a+v}{a-v}$; consequently the height y fallen through being given, the velocity v will be given, and from the known velocity the time t may be found.

APPLICATION.

Let $y=1848$, and the diameter of an iron ball four inches; then if $n=1$ we have $m=7425$, by our tables, and $r=\frac{6md}{p}$ (*Art.* 159.) becomes $r=14850$

feet. Hence $\frac{2y}{r}=.249$ nearly, this being the hyperbolic logarithm of the number q ; and to find its number, we must find first its tabular logarithm, which is done by multiplying the hyperbolic logarithm .249 by the constant number .43429, by *Art.* 59; therefore that logarithm is 108138, and the number answering in the common tables is $q=$

1.282746; hence $vv=q-1 \times \frac{r}{q}=3273.27$; this number multiplied by 32.2, twice the height fallen thro'

thro' in a second, by *Art.* 86. gives 105399.294, whose square root 324.65 will be the number of feet moved over uniformly in a second with the velocity acquired in the descent thro' 1848 feet. The descent thro' the same height in a non-resisting medium would be 345 feet.

To find the time, by means of the equation, $t = \frac{1}{2}a \log. \frac{a+v}{a-v}$, we have $a=121.86$, $v=57.21$, hence $a+v=179.07$, $a-v=64.65$, and $a+v$, divided by $a-v$, gives 2.7698, whose tabular logarithm is .442448; which being multiplied by the constant number 2.302585, gives 1.01877 for the hyperbolic logarithm of that number, by *Art.* 58; and this number multiplied by $\frac{1}{2}a=60.93$, gives $t=62.0736$ feet, which divided by 5.67, the square root of 32.2, (*Art.* 85.) gives 10.94 seconds.

In a non-resisting medium the body would have fallen thro' the same height in the time of 10.71 seconds, and therefore the difference is 0.23 seconds only.

EXAMPLE.

167. The body being projected upwards in a vertical direction, with a given velocity c ; to find the height to which the body will rise, as well as the velocity and time. We have $v\dot{v} = -\dot{y} - \frac{vv}{r}\dot{y}$ in this case, when gravity

is unity, or $\dot{y} = \frac{-r\dot{v}}{r+vv}$, whose fluent is $y = A - \frac{1}{2}r \log. \frac{r+vv}{r+cc}$; but when $v=c$, then $y=0$; and $A = \frac{1}{2}r \log. \frac{r+cc}{r}$; therefore $\frac{2y}{r} = \log. \frac{r+vv}{r+cc}$. If q expresses the number whose hyperbolic logarithm is $\frac{2y}{r}$, then $q = \frac{r+cc}{r+vv}$; from whence we get $vv = \frac{r+cc}{q} - r$.

The equation of the time will be $t = \frac{-rv}{r+vv}$; hence if A denotes an arc of a circle, whose radius is $r^{\frac{1}{2}}$, and tangent c and n another arc of the same radius, but v the tangent, we shall have $t = A - n$.

APPLICATION.

Let $cc = 5280$ feet and the rest as before; then when the body arrives to its highest point, we have $y = \frac{1}{2}r \log. \frac{r+c}{r} = \frac{1}{2}r \log. 1.3555$; the tabular logarithm of this number is .1311, which multiplied by the constant number 2.30258, by *Art.* 59, gives .30187 for the hyperbolic logarithm of that number, and this logarithm multiplied by $\frac{1}{2}r$, 7425, gives $y = 2241$ feet. As the body would have risen to the height of $\frac{1}{2}cc$ or 2640 feet in a non-resisting medium, the difference 379 feet is the part lost by the resistance. In the equation of the time $t = A - n$, the arc n vanishes when $u = 0$; and the radius $r^{\frac{1}{2}} = 121.860$ is to the tangent $c = 72.663$ as unity is to .596282, which answers to an angle of 30 degrees and 48 minutes, that is to 30.8 degrees. Now 180 degrees is to 30.8 as half the circumference $3.14159\sqrt{r}$ is to 65.507 feet for the arc A , and being divided by the square root 5.67 of 32.2, by *Art.* 85, gives 11.55 seconds for the time t , which time in falling thro' the same height in a non-resisting medium, would have been 10.7 seconds, and the difference is .85 seconds.

REMARK.

It must be observed, that a falling body in a resisting medium can never acquire a velocity equal to that we have called the greatest, *Art.* 158, but may perpetually approach it. For when $r = vv$, in *Art.* 166, both the space descended thro' and the time of the

the fall become infinite; so that the greatest velocity is only a limit to which the other velocities may approach without ever attaining it. This appears likewise, from observing, that both the acceleration and retardation or resistance, are in proportion to the squares of the velocities; this being constantly the case, the diminution can never arrive to that degree as to become entirely equal to the augmentation: and consequently the resistance of falling bodies can never become geometrically equal to the force of gravity, nor the velocity acquired in the fall to the greatest of all.

E X A M P L E.

168. If the body moves in a canal or horizontal right line with a given velocity; then because $y=z$ in this case, and supposing gravity to be unity, the equation in *Art.* 163. becomes here $v\dot{v} = y - \frac{vv}{r} \dot{y}$ or

$\dot{y} = \frac{rv\dot{v}}{r-vv}$, whose fluent is $y = A - \frac{1}{2}r \log. \frac{r-vv}{r}$; but

when $y=0$, the velocity v becomes equal to the given one c ; hence $A = \frac{1}{2}r \log. \frac{r-cc}{r}$, and therefore the compleat fluent is $\frac{2y}{r} = \log. \frac{r-vv}{r-cc}$. Because the equation

of the time is $\dot{y} = v\dot{t}$, and as $\dot{y} = \frac{rv\dot{v}}{r-vv}$, by what has

been shewn above, we get $\dot{t} = \frac{r\dot{v}}{r-vv}$; and if $r=aa$,

the fluent of this equation will be $t = A + \frac{1}{2}a \log. \frac{a+v}{a-v}$; and when $t=0$, then $v=c$; hence $-A = \frac{1}{2}a$

$\log. \frac{a+c}{a-c}$, and therefore $\frac{2t}{a} = \log. \frac{a+v}{a-v} - \log. \frac{a+c}{a-c}$ or

$$\frac{2t}{a} = \log. \frac{a-c \times a-v}{a+c \times a-v}$$

C O R:

C O R.

169. Hence it is manifest, that the given velocity, with which the body begins to move, must always be less than the greatest velocity that the body can acquire in the medium. For when $r=ce$, or $a=c$, then both the space moved over and the time elapsed vanish. Although this is very remarkable, yet no author has taken notice of it; on the contrary they often suppose this velocity much greater.

It must be observed that the same thing is true, whether the body moves in an horizontal line or in one that makes a given angle with it, since every thing will be the same as above. This we thought necessary to take notice of, to prevent an uncautious reader from taking what has been here said as only applicable to horizontal ranges.

P R O B L E M. *Fig. 55.*

170. *Supposing the powder all fired before the shot is sensibly moved from its place, to find the velocity with which it issues out of the piece.*

Let n express the ratio of gravity to the force of powder, the height of the charge $AB=a$, any distance $AC=x$, $BD=b$; then because the force of the powder confined in the space AB is to its force, when extended to the space AC , reciprocally as the square of the distance AC to the square of AB , by Cor. I. after Theo. IV. of our Artillery; that is, the force of the powder at C will be expressed by $\frac{aan}{xx}$. Now this value being wrote for p , and $\dot{x}$ for $\dot{y}$ and $\dot{z}$ into the general equation, *Art.* 163. as well as $\frac{vv}{r}$ for R ,

gives $v\dot{v} = \frac{aan\dot{x}}{xx} \times 1 - \frac{vv}{r}$ or $\frac{rv\dot{v}}{r-vv} = \frac{aan\dot{x}}{xx}$, whose flu-

ent

ent is $A - \frac{1}{2}r \log. \frac{r-vv}{r-vv} = -\frac{aan}{x}$; but when x becomes $=a$, then v vanishes, so that $A = \frac{1}{2}r \log. r - an$. Hence $\frac{1}{2}r \log. \frac{r-o}{r-vv} = an - \frac{aan}{x}$, or because $b = x - a$, we get $\frac{2abn}{rx} = \log. \frac{r-o}{r-vv}$. But if q denotes the number whose hyperbolic logarithm is $\frac{2abn}{rx}$,

then $q = \frac{r}{r-vv}$, or $vv = q - 1 \times \frac{r}{q}$.

C O R.

171. From whence it follows, in the same manner as in the last problem, that the velocity v must always be less and never can be equal to the greatest velocity the ball may acquire in that medium: and consequently the resistance $\frac{vv}{r}$ of the body arising from its motion in the medium can never be equal, much less exceed the weight of the shot, as has erroneously been supposed by some authors.

OBSERVATIONS.

We have neglected the resistance of the shot rubbing against the inside of the gun, in the last problem, the force lost thro' the vent and windage, and supposed the powder to fire all at the same instant, which is not true by experiments; besides, if it was so, the shortest guns would be the best: but as we do not know by what law the flame proceeds, we are not able by theory alone to determine the exact velocity, the best length of a piece, or the proper charge. All therefore we are certain of is, that the velocity found by theory

is

is greater than it should be, and yet less than that it may acquire in that medium : this appears from the equation $vv = \overline{q-1} \times \frac{q}{r}$; for unless q is so great that unity may be neglected, we cannot have $vv = r$.

Again, if we consider that the hyperbolic logarithm of $\frac{2abn}{rx}$, or $\frac{ab}{x}$, because $2n$ and r are constant numbers, should have certain limits to determine the length and charge : which it has not, for the rectangle made by two lines a , b when the sum x of their length is given, is the greatest of all, when they are equal, that is when the powder fills up half the vacant cylinder, which we know not to be true from undeniable experiments.

Tho' the theory is not sufficiently exact, yet it serves as a guide to experiments; for which reason it ought to be well understood by those who undertake to improve artillery.

Before this theory can be applied to particular examples, the ratio of the elastic force of powder to that of gravity must be determined. Let the pressure of the atmosphere be equal to a column of quicksilver of 29.5 inches high : then as the specific gravity of quicksilver is to that of cast iron as 14000 to 7425, the pressure of the atmosphere will be equal to a column of iron 55.6 inches high, or 4.635 feet. Whence if we suppose with Mr. Belidor and most authors, Mr. Robins excepted, the elastic force of powder to be 4000 times greater than that of air; if we multiply 4.635 by 4000, we get 18540 feet for a column of iron, expressing the elastic force of powder : and as an iron ball, whose diameter is denoted by d , is equal to a cylinder of the same base and $\frac{2}{3}d$ high, the force of gravity is therefore to the force of powder as $\frac{2}{3}d$ is to

to 18540, or as unity to $\frac{27810}{d}$, when the diameter d is expressed in parts of a foot.

EXAMPLE.

172. *Let an eighteen pounder of nine feet long be loaded with nine pound of powder, to find the velocity at the mouth of the cannon.*

The diameter d of an eighteen pound iron ball is five inches; and the height a of the charge ten inches or twice the diameter; then if we multiply 27810 by $\frac{5}{12}$, we get $n=66744$, $x=9$, $b=8\frac{1}{6}$; and as the specific gravity of cast iron is to that of air as 7425 is to $\frac{7}{8}$, we get $r=21214$, *Art.* 159; hence $\frac{2abn}{rx}=4.758$, for the hyperbolic logarithm of the number q ; and this number multiplied by the constant number .43429, *Art.* 59. gives 2.06635 for the tabular logarithm of that the number, which gives $q=116.5$.

Whence $vv=q-1 \times \frac{r}{q}=21032$, or 823 feet per second for the velocity required; which differs from the greatest 826 by 3 feet only. Tho' this agrees very exactly with the experiments made at Minorca, as will be seen hereafter, when we treat of projectiles in the air, yet it does not answer in other examples; for the reasons given in the observations above.

REMARK I.

What has hitherto been said in this work, was only to prepare the reader to understand what relates to Artillery and Gunnery, as being the subject for which it was chiefly intended, and could not be explained without the preceding principles.

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Not-

Notwithstanding the great importance of artillery to a state, neither the best length nor charge of guns, so as to carry their shot to the greatest distance, have as yet been determined to any tolerable degree of perfection. This is not to be wondered at, because all the experiments made here and abroad were so injudiciously contrived, that no true conclusions could be made from them; to this may be added, that the practical artillerists are entirely ignorant of the theory, and the theorists of the practice; and without being acquainted with both, it is obvious that no improvements can be made in this science.

Authors, for want of having a proper guide to go by, have greatly over rated the velocities of cannon shot, as likewise the resistancethey meet with moving in the air; and as long as this error subsisted, no improvement could be expected. For we have proved in *Art.* 171, that a body impelled by any force whatever can have no greater velocity than that it may acquire by falling in that medium, which we have shewn, in *Art.* 159, cannot exceed a certain limit, which we have determined by means of the ratio between the specific gravities of the medium to that of the body. Thus if the body is of cast iron, its specific gravity is to that of air as 7425 to $\frac{7}{8}$, by the Remark after *Art.*

147; whence if we make $m=7425$, and $n=\frac{7}{8}$ in the

equation, $r=\frac{6md}{n}$, *Art.* 159. we get $r=50914d$;

which shews that twice the height thro' which an iron ball descends so as to acquire the greatest velocity, is less than 50914 times its diameter expressed in feet; it shews likewise that the greatest velocities which iron balls can acquire are in proportion to their diameter.

From

From whence we can correct an error commonly committed in the construction of guns; which is, that those of smaller calibers, under a twenty-four pounder, are made longer in proportion than larger ones. The practitioners imagine, that by this means the velocities of their shot may be increased, and go as far as larger ones, by loading them with greater charges. A few examples will prove what we here advance.

Let the cannon be a six pounder, whose diameter of the shot is 3.498 inches, then 50914, multiplied by 3.498, and divided by 12, gives 14841.4 feet; this multiplied by 32.2, and the square root of the product extracted, gives 691 feet for the space moved over uniformly by the greatest velocity a six pounder shot can possibly have in a second.

If the cannon be a forty two pounder, whose diameter of the shot is 6.684 inches, then we get $r = 28359$ feet; this multiplied by 32.2 and the square root extracted from the product, gives 955 feet in a second for the greatest velocity a forty-two pound ball can possibly have.

When the ball is of lead, then as the specific gravity of lead is to that of air as 11325 to $\frac{7}{8}$; the equa-

tion $r = \frac{6md}{n}$ gives $r = 77657d$ feet, for the shot whose diameter d is expressed in feet. Thus for example, the diameter of a bullet, twelve of which weigh a pound, is $\frac{3}{4}$ of an inch nearly; so that writing $\frac{1}{10}$ for d , we get $r = 4854$ feet nearly; this multiplied by 32.2, and the root extracted from the product, gives 395 feet in a second, for the greatest velocity which that shot can have.

Again, let the diameter of the shot be 1.69 inches, whose weight is a pound; then $r = 77657d$ becomes

P 2

r =

$r=10937$ feet in this case, and proceeding in this manner as before we get 593 feet in a second, for the greatest velocity this shot can have.

As a cannon shot weighing 55 pounds, whose diameter is 7.32 inches nearly, will by proceeding as before, give 1000 feet in a second for the greatest velocity this shot can have, and as no shot of this size has hitherto been used, nor will hardly hereafter. we may draw this conclusion, that no cannon or musket shot can acquire a velocity, by any force impressed upon them, of a 1000 feet in a second; consequently those authors who have supposed, that they may have a velocity of 2000 feet and upwards, in a second, were greatly mistaken: and whatever has been deduced from this supposition must be erroneous.

It will not be improper to take notice here of a gun which the French made use of in the last war in Germany called *Amufette*, which fired a leaden ball of a pound weight; and it was said that it carried its shot farther than any of our light six pounders. Now, we have shewn, that the greatest velocity of a six pound shot approaches to 691 feet in a second, and that the greatest velocity of a leaden ball weighing one pound, approaches to 593 feet in a second; consequently, if the fact be true, our six pounders have not their proper length, nor charged with a proper quantity of powder. This appears likewise from other circumstances, as will be shewn hereafter.

We have hitherto spoke of solid shot, we shall now consider the velocities of shells; whose specific gravity is to that of air as 4892 to $\frac{7}{8}$, supposing them all similar to the thirteen inch ones, as we have done in our general construction: hence $r = \frac{6md}{n}$,
gives

gives $r=33545d$, in respect to all shells made of cast iron. Now, as the diameter of a thirteen inch mortar shell is $12\frac{3}{4}$ inches, we get $r=35641$ feet, and proceeding in the same manner as before, we shall find 1072 feet in a second, for the greatest velocity which this shell may have.

The diameter of a ten inch mortar shell being $9\frac{3}{4}$ inches, we get $r=27255$ feet, and 936.8 feet in a second, for the greatest velocity. Hence the same thing is true in regard to shells as in solid shots; that is, the greatest shells may have the greatest velocities.

R E M A R K II.

For the reason that authors have supposed the velocities of bodies moving in a resisting medium, may increase without any limits; the resistance may likewise increase without end, which is a great mistake, as we are going to prove. Suppose the density of the air is the same every where as near the surface of the earth, and a body was to fall in it, by the force of gravity; it is proved, that if there was no resistance, it would descend through spaces which are as the squares of the velocities acquired in the fall; but, by reason of the resistance of the air, the spaces described will be diminished in some proportion, which can be no other than that of the squares of the velocities: since the number of particles of the air which the body strikes in a given time, are as the velocities, and as the body strikes each particle with a force that is likewise as the velocity, the resistance arising from these causes must therefore be as the squares of the velocities. In oily or unctuous mediums, the resistance will be increased in proportion to the difficulties of separating the particles; which is not the case in the air. Now the descent of bodies
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in the air, being accelerated in the duplicate ratio of the velocities, and resisted in the same duplicate ratio; the body can never arise to such a degree of swiftness, as that the force of resistance becomes equal to the force of gravity, altho' it may continually approach it. If these two forces should become equal at last, the motion would then become uniform; this has been shewn in *Art.* 153. and 158. consequently, the force of resistance can never become equal, much less greater, than the force of gravity in this case.

What has been said here in respect to falling bodies, is equally true in bodies impelled by any force in any direction, by *Art.* 169. and 171; besides, the force by which the body is impelled, let it be what it will, cannot be instantaneous, and must act gradually by degrees; and whenever its action becomes equal to that of the resistance, it cannot increase beyond it. Consequently, the notion, that the resistance may increase without any limit, is without foundation; and those experiments that have been made on that account must be fallacious, and can avail nothing,

An objection is made against the resistance being in all cases as the squares of the velocities; when the motion of the body is so great, as that the air cannot fill up immediately the space left by the shot, the resistance must be greater than it would be in slower motions; from thence is inferred, that the resistance is, in slow motion, as the square of the velocity, but in very swift motion, in the triplicate ratio of the velocities.

But as in the estimation of the resistance, the force by which the motion of bodies are retarded, was determined from the number of the particles the body strikes, and from the velocity with which they are struck, and not from any consideration of their moving
round

round and striking the body behind; so that where this last action takes place, and is sufficient to produce an effect that can be estimated, the contrary conclusion to that objection may be made, viz. that in slow motion the resistance is diminished, by reason of the particles striking the body behind in the direction of its motion. This seems to be confirmed by experience, for in slow motion the curve described by the body is sensibly a common parabola. This will be found by measuring the ranges, as I have done several times at La Fere; and always found in distances not exceeding 300 yards, that the ranges never differed above a few feet from that computed by means of the parabola.

The next query is, whether a shot can move so fast as to leave any vacuity behind it, or not? If not, the objection is without foundation. Now, to find this recourse must be had to experiments. We have shewn in *Art. 147*, that the sound moves over a space of 1142 feet in a second; and as the sound is propagated by the elastic force of the air, therefore the elastic force of the air is such as to produce a velocity which moves over the same space in the same time: we have likewise proved in the last remark, that a thirteen inch shell may have the greatest velocity of any shot or shell that are now in use; and as this velocity is such as to move over a space of 1072 feet in a second, when uniformly continued, which being less than the velocity produced by the elastic force of air, it is manifest, that no common shot can move so quick as to leave a vacant space behind it.

R E M A R K III.

Having determined the greatest velocities that cannon shot can possibly have, let the charge be what it will, in the first remark; and the greatest resistance they

they can meet with in their motion, in the second; it remains now to apply these principles to the improvement of artillery, which were necessary to be known; and for want of them authors have been greatly misled in their computations and conjectures.

Of all questions that ever have been made in respect to artillery, the most important ones are those concerning the proper length of guns and their best charges, so as to carry their shot as far as possible; but authors, as well as practitioners, differ widely in their opinions: which could be no otherwise; for the practitioners, for want of theory, have been misguided in their inquiries, and most experiments that were made are such as no conclusions that are really useful can be drawn from them; for those made at Dunkirk by Dumetz, with various calibers of the same length, and with charges weighing two-thirds of their shot, can answer no other purpose than that a twenty-four pounder carries its shot farther than any other caliber of the same length; and those made at Woolwich by the late General Armstrong, with six twenty-four pounders of various lengths, can serve only to shew that the piece which was 9.5 feet carried farther than those of any other length with the same charge. These are the most remarkable experiments that I know of, excepting those made at Minorca by General Williamson, and are the only ones to the purpose, as will be shewn hereafter.

In the first period of time, guns were made very heavy and long, so as to bear charges equal in weight to that of their shot, imagining that the greater the charge was, the longer the gun should be, in order to give time for the powder to take fire before the shot leaves the piece, and to produce the greater effect. In Queen Elizabeth's time, we and the French made them of excessive length, as appears by the cul-

verin in Dover Castle, and that at *Nancy*; whilst the Spaniards made them shorter and lighter than they have been ever since. For at the siege of Breda by the Spaniards, under the command of Spinola, in 1625, Count Mansfield, a German in their service, cast several six pounders which weighed 180 lb. only, and twenty-five pounders of 700 lb. weight, besides another Spanish gun, which we had at Woolwich, that weighed 2100, and carried a ball of 41 lb. of our weight. In King Charles the Second's time, we made thirty-two pounders of iron that weighed 4200, and forty-two pounders of 5300 weight; whereas our present twenty-four pounders brass weigh 5200, and the forty-two of the same metal, 6600. To make brass guns so much heavier than the iron ones of the same caliber, is inexcusable in the land of liberty, where reason and good sense prevails in almost every thing else.

It is true we make at present our six pounders for the field-service so light and short, that four of them weigh no more than one of the long and usual guns; but we are to thank the late Duke of Cumberland for it, who, by his authority and judgment in artillery, insisted upon the diminishing the weights. This great variation in the length of guns sufficiently shews how little improvement has been made in artillery since its first invention.

The theorist having no proper experiments, and being unacquainted with the practice, it is not strange that he should be so much mistaken in his calculations; but that he should fall into the vulgar error of the practitioners, by admitting the velocities to increase in proportion as they are charged with a greater quantity of powder, without any limitation, is inexcusable. As they all admit the resistance which moving bodies meet with, one should think it required no great degree of penetration to

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perceive

perceive that this resistance must diminish the velocities in some cases, so as not to be augmented.

What can be the reason that no nation has endeavoured to improve their artillery, which is essentially useful and so very expensive, unless it is the prepossession in favour of an old established custom, in such a manner, as neither reason or good sense can overcome? It will therefore be in vain for an author to endeavour all he can for improving artillery, unless some great man, who has it in his power to undertake it, and to give proper encouragement to those that are qualified for carrying on a series of proper experiments.

To facilitate such an undertaking, we shall here insert the experiments made at Minorca, in 1745, with two iron eighteen pounders, one of eleven feet long, weighing 51 : 0 : 5, and the other nine feet long, and weighing 39 : 1 : 3, under the direction of General Williamson, assisted by Major Hislop and several other artillery officers.

These two pieces were firmly fixed in the ground, and so contrived as to have any degree of elevation; they were loaded with different charges and elevations, from six to forty-five degrees. After having made a great many trials, it was found, that, when they were loaded with nine pounds of powder, they both carried their shot farthest, in all the different elevations; and that the gun of nine feet long carried always farther than that of eleven feet long. These experiments being authentick, furnish a sure direction to determine the proper length of guns, and their best charges. For as the diameter of an eighteen pound shot is 5 inches nearly, this gun is $21\frac{3}{8}$ diameters of the shot long; and therefore the best length of any gun is about 21 diameters of its shot, and its best charge equal in weight to half that of its shot. But before a state is resolved to reform
their

their artillery, it would be necessary to repeat those experiments, and try various calibers in the same manner, in order to be certain that this alteration, in so essential a thing, may be done upon solid and firm considerations; and when the dimensions of guns are fixed, a law should be made to prevent any one to alter them for the future.

Having thus ascertained the limits of the greatest velocities, as well as the resistance of bodies, the ranges of shells and shots may be determined in a greater degree of exactness than by any other theory yet published; and as their computation is very intricate and laborious, tables may be constructed so as to apply the theory to practice with ease.

P R O B L E M. Fig. 56.

173. If a body projected with a given velocity c , in a given direction AE , to find the nature of the curve ADK described by the body.

If $AP=x$, $PM=y$, the arc $AM=z$, v the velocity at M , t the time and $p=1$; then $v\dot{v}=-\dot{y}-R\dot{z}$, by *Art.* 163. for the ascent, and $\dot{x}\dot{y}=vv\dot{x}$, by *Art.* 164.

when $\dot{z}$ is constant, and as $R=\frac{vv}{r}$, by *Art.* 165. If the values of y and R are wrote into the equation above, we get $v\dot{v}=-\frac{vv\dot{x}}{x}-\frac{v\dot{v}\dot{z}}{r}$, or dividing by vv , and transposing the first term of the second side to the first, we shall have $\frac{\dot{v}}{v}+\frac{\dot{x}}{x}=-\frac{\dot{z}}{r}$, whose fluent is

$\log. \frac{v\dot{x}}{\dot{z}}=-\frac{Az}{r}$; because $\dot{z}$ has been supposed constant, it was necessary to divide by it, to make the terms homogeneous: and if q expresses the number whose hyperbolic logarithm is unity, that is, $lq=1$;

then $\log. \frac{v\dot{x}}{\dot{z}} = -\frac{A\dot{z}lq}{r}$, and hence by logarithms $\frac{v\dot{x}}{\dot{z}} =$

$Aq^{-\frac{z}{r}}$. Now at the beginning A of the motion $v=c$, $z=0$; and if the radius is to the cosine of the angle EAP as unity to k ; then $\dot{z}:\dot{x}::1:k$, we get $A=ck$, consequently the last equation becomes $\frac{v\dot{x}}{\dot{z}} = ckq^{-\frac{z}{r}}$.

And since $\dot{z}^2 + vv\ddot{y} = 0$, by *Art.* 164, when $\dot{x}$ is constant, and if for convenience we make $\dot{x}=1$; by comparing the values of vv , in the two last equations, we get $q\frac{2\dot{z}}{r} + cck\ddot{y} = 0$, for the equation of the curve.

174. Since $\dot{z}=vt$, by *Art.* 161, and $v\dot{x}=ck\dot{z}q^{-\frac{z}{r}}$, by above; the two values of v being compared, give $\dot{x}q^{\frac{z}{r}}=ckt$. Hence if we suppose the resistance to vanish, r becomes infinite, and therefore $q^{\frac{z}{r}}=1$; so that the equation $\frac{v\dot{x}}{\dot{z}}=ckq^{-\frac{z}{r}}$ becomes $v\dot{x}=ck\dot{z}$, and $\dot{x}q^{-\frac{z}{r}}=ckt$, gives $\dot{x}=ckt$, or $x=ckt$, which values of v and t are exactly the same as we have found in *Articles* 94, 95, where we have treated of bodies moving in a non-resisting medium.

Several authors have found the equation of the curve, nearly expressed like this here; but in finding the relation between the abscissa and ordinate, no one that I have as yet seen has done it in a general manner. Thomas Simpson has, it is true, an equation that will answer when the projectile velocity is but small; but when it is great, he has recourse to the differential method. But as some quantities are to be found by trials, which in many cases is not to be done without an infinite deal of labour, it is obvious that his method is often impracticable.

Professor

Professor Euler solved this equation by an infinite series in his Mechanics, as likewise in his Comment upon Robins's Gunnery, but neither of them converge so as to answer all cases; we shall here give three different solutions of this problem, and it will appear that they all give very nearly the same ranges.

FIRST METHOD.

PROBLEM.

175. To solve the equation $q \frac{2z}{r} + cckk\ddot{y} = 0$, by an infinite series.

The fluxion of $-q \frac{2z}{r} = cckk\ddot{y}$ is $-\frac{2\dot{z}}{r} q \frac{2z}{r} = cckk\ddot{y}$, and these two equations multiplied crossways, so as to expel the exponential quantity, gives $r\ddot{y} = 2\dot{z}\ddot{y}$.

If now we suppose y equal to an infinite series composed of the powers of x and constant but indetermined coefficients a, b, c, d, f , and make $\dot{x} = 1$, for convenience's sake; then the first, second, and third fluxions, give *

$$y = ax + bx^2 + cx^3 + dx^4 + fx^5$$

$$\dot{y} = a + 2bx + 3cx^2 + 4dx^3 + 5fx^4$$

$$\ddot{y} = 2b + 6cx + 12dx^2 + 20fx^3$$

$$\ddot{\dot{y}} = 6c + 24dx + 60fx^2.$$

As the square of $\dot{y}$ added to the square of $\dot{x}$, is equal to the square of $\dot{z}$; by making $1 + aa = ss$, we get $\dot{z}^2 = ss + 4abx + \frac{6ac + 4bb}{s} x^2$, and if $m = \frac{3ac}{s} + \frac{2bb}{s^3}$; the square root of the last equation will

be $\dot{z} = s + \frac{2abx}{s} + mx^2$. Hence, by writing the values

of

of $\dot{z}$, $\ddot{y}$ and $\ddot{y}$, into the equation $r\ddot{y} = 2\dot{z}\ddot{y}$, we get,

$$\left. \begin{aligned} 3cr + 12rdx + 3orfxx = 2bs + 6cs \\ + \frac{4abb}{s} \end{aligned} \right\} x \quad \left. \begin{aligned} + 12ds \\ + 2bm \\ + \frac{12abc}{s} \end{aligned} \right\} xx.$$

And by comparing the coefficients of the homologous terms, then will

$$3rc = 2bs, \quad \text{or } 3c = \frac{2bs}{r}$$

$$6rd = 3cs + \frac{2abb}{s} \quad 3d = \frac{bss}{rr} + \frac{abb}{rs}$$

$$15rf = 6ds + bm + \frac{6abc}{s} \quad 15f = \frac{2bs^3}{r^3} + \frac{8abb}{rr} + \frac{2b^3}{rs^5}.$$

Consequently the value of y will be expressed by the powers of x and known quantities, when the coefficients a and b are determined: in order to which we are to consider, that at the beginning A of the motion, where $x=0$, the second equation gives $\dot{y}=a$, or $\dot{y}=ax$, because $\dot{x}$ has been supposed unity. Hence, $\dot{x}:\dot{y}::1:a$; but $\dot{x}$ is to $\dot{y}$ as the radius to the tangent of the angle EAP of elevation; therefore a denotes the tangent of that angle, when the radius is unity, and because $1+aa=ss$, s expresses its secant.

The third equation gives $\ddot{y} = 2b$, when $x=0$, and writing $2b$ for its equal $\ddot{y}$ into the equation $\dot{z}^2 + vv\ddot{y} = 0$, *Art.* 164, we get $\dot{z}^2 + 2bvv = 0$, but at A , v becomes c and $\dot{z}$ the secant s ; so that $ss + 2bcc = 0$, or if b expresses double the height thro' which a falling body acquires the projectile velocity c , then $b=cc$, and $ss + 2bb = 0$.

As the values of y converges but slowly in many cases, without some alteration being made in the equation, the most simple and shortest way to do this, is to write bx and by , for x and y ; or which is the same, by multiplying the coefficients of each term,

term, by the powers of b one less than that of x ; that is, a, b, c, d, f , by $1, bb, b^3, b^4$, and making $\frac{b}{r} = n$, which gives $y = ax - bx^2 - cx^3 - dx^4 - fx^5$, and $2b = ss, 3c = ns^3, 12d = 2nns^4 - ans^3$, and $6of = 4n^3s^5 - 8an^2s^4 + ns^3$.

It must be remembered, that when the values of x and y are found by this equation, to multiply them by the value of b , according to the supposition we have made, to get their true values. Hence, when the ratio n of the greatest and projectile velocities is given or remains the same, the values found for x and y will serve to determine the various parts of the projectiles in general, and if after the diameters of the shots are given, the particular cases of each will be obtained; which renders the equation very general; and it will be shewn hereafter that the five first terms are sufficient to find the value of x very nearly.

C O R.

176. When the resistance becomes insensible, all the terms divided by r , or, which is the same, multiplied by n , vanish, and therefore the equation becomes in this case $y = ax - bx^2$, which is precisely the same as that found in *Art.* 98, because $ak = l$, $2bb = ss$, and b is here double to what it is there.

C O R.

177. When the body is projected in an horizontal direction, the tangent a vanishes, the secant s becomes equal to the radius or unity, and y negative. Therefore $y = bx^2 + cx^3 + dx^4 + fx^5 + gx^6 + kx^7$; and $2b = 1, 3c = n, 6d = nn, 6of = 4n^3 + n, 18og = 4n^4 + 7nn, 315k = 2n^5 + 16n^3 - \frac{3}{8}n$.

The equations $v\overline{q\overline{r}} = c\overline{k\dot{z}}$, $\overline{q\overline{r}} + c\overline{k\ddot{y}} = 0$, found in *Art.* 173 and 164, become here $v\overline{q\overline{r}} = c\overline{k\dot{z}}$ and $\overline{q\overline{r}} = \overline{b\ddot{y}}$; because the sign k is here unity, $b = cc$, and $\ddot{y}$ negative;

negative; but it must be observed that c denotes the projectile velocity at the vertex D where the motion is supposed to begin.

C O R.

178. Since $\dot{x} = vt$ and $\dot{x}^2 + vv\ddot{y} = 0$, by *Art.* 164, we get $-\ddot{y} = it$. Now as the third equation $\ddot{y} = 2b + 6cx + 12dx^2 + 2ofx^3$, gives $-\ddot{y} = ss + 2ns^3x + 12dx^2 + 2ofx^3$, because $-2b = ss$, and $-3c = ns^3$. And if $2ms = 12d - nns^4$; $25p = 2of - 2mnss$; the square root of this equation will be $i = s + nssx + mx^2 + px^3$; whose fluent is $t = sx + \frac{1}{2}nssx^2 + \frac{1}{3}mx^3 + \frac{1}{4}px^4$.

The velocity v is found by writing the ratio of $\dot{x}^2$ to $-\ddot{y}$ into the equation $\dot{x}^2 + vv\ddot{y} = 0$.

P R O B L E M. *Fig.* 56.

179. *The angle of elevation being 45 degrees, and the ratio n unity, to find the horizontal range AK.*

Here we have $a = 1$, $ss = 2$, and $n = 1$, by supposition; and hence $b = 1$, $c = .9428$, $d = .4309\frac{2}{3}$, $-f = .10907\frac{1}{4}$. Therefore the equation in *Art.* 175, becomes here $y = x - x^2 - .9428x^3 - .431x^4 + .109x^5$. Now when $y = 0$, then $1 = x + .9428x^2 + .431x^3 - .109x^4$. And by the general rule for finding the roots, *Art.* 78, we get $1 : 1 : 1.9428 : 3.3166 : 5.4703 : 9.3255 : 15.7 : 26.4883 : 44.7133 : 75.4367 : 127.2975 : 214.8034$, and the penultiem term divided by the last gives $x = .59262$, true in all its places. For $x^2 = .351198$, $x^3 = .2081269$, $x^4 = .12334$, and these values wrote into the equation gives $1 = 1.0000$.

P R O B L E M.

180. *To find the abscissa AB, corresponding to the greatest ordinate BD.*

By taking the fluxion of the last equation, and making $\dot{y} = 0$, we get $1 = 2x + 2.8284x^2 + 1.7238x^3 - .5454x^4$. By finding the root as before, then $1 : 2 : 6.8284 : 21.0374 : 64.2904 : 198.763 : 611.905 : 1885.3413$; and $x = .32456$, nearly. For $x^2 = .10534$,

10534, $x^3 = .034189$, $x^4 = .011096$; and these values wrote into the equation gives $1 = .999946$.

P R O B L E M.

181. To find the greatest ordinate B D.

If we write the values of x and its powers here found into the equation $y = x - x^2 - .9428 x^3 - .431 x^4 + .109 x^5$, we get $y = .1826$.

E X A M P L E.

182. We have found $r = 50914 d$, in the first remark after *Art.* 172; when the shot is of cast iron: now, if it weighs 18 pounds, the diameter d will be five inches, and hence $r = 21214$ feet, and as b has been supposed equal to r ; if we multiply 21214 by .59262, the value of x found in *art.* 179, we get 12572 feet or 4190 yards for the greatest horizontal range A K, which an 18 pounder can have when elevated at an angle of 45 degrees.

If the value of x (.32456) in *Art.* 180, be multiplied by 21214, we get 6885 feet or 2295 yards for the abscissa A B corresponding to the greatest ordinate, and .1826, multiplied by 21214, gives 3873.6 feet or 1291.2 yards for the greatest ordinate B D.

R E M A R K.

In order to shew the near agreement of our Theory with the best experiments made by General Williamson, &c. at Minorca in 1745, as mentioned before; who, after many trials with various charges, found that nine pounds of powder carried an 18 pound ball farther than any other charge: having also tried two 18 iron pounders, one of nine and the other eleven feet long, found that of nine feet long carried its shot farther than that of eleven, under the same circumstances: after this, he set down six

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experi-

experiments, made with that of 9 4160 yards,
 feet long, loaded with 9 pounds of 4135
 powder, and elevated at an angle of 4100
 45 degrees: their ranges were found 4015
 as in the margin. 3980

This shews how nearly the range 4190, we found in *Art.* 182, by our theory agrees with these experiments; for we have supposed a greater projectile force than the shot could have, as shewn before: and hence it is plain, that the specific gravity of air is $\frac{7}{8}$ nearly, as we determined it after, *Art.* 147. For if we suppose it to be $\frac{5}{4}$, as Cotes has it, then

$r = \frac{6md}{n} = 14850$ feet, which multiplied by .59262, the value of x , gives 2933 yards; this being less than any of those above, shews that this supposition is contrary to these experiments, since this range should be rather more than less, because we have supposed a greater projectile velocity than the shot could have.

If we make use of Sir Isaac Newton's rule $r = \frac{8md}{3n}$ for determining the greatest velocity, and suppose $n = \frac{7}{8}$ according to our rule, we shall find $r = 9429$ feet, which when multiplied by .59262, the value of x , we get 1862.6 yards for the horizontal range, which plainly proves this rule to be deficient likewise: and if we were to make $n = \frac{5}{4}$, as Cotes has it, the range comes out to be 1303.4 yards only; and consequently our determination of the specific gravity of the air, and that of the resistance, is the nearest that possibly can be found.

In all arts and sciences of a mixt nature, a criterion is absolutely necessary to try their certainty; for the least error, either in the facts or suppositions, misleads an author in such a manner as that the conclusion

clusion comes out quite contrary to what it ought to be : this has been the case of all the authors that have wrote upon motion in a resisting medium, for want of having good experiments to compare with their calculations : as to those experiments which some of them made, they were too diminutive to draw proper consequences from them, whereby they could not establish their theory upon a proper foundation.

P R O B L E M.

183. *To find the velocity at the highest point D, and the time of describing the arc A D.*

We have $\ddot{y} = 2b + 6cx + 12dx^2 + 20fx^3$, by *Art.* 175 ; and as $-b = 1$, $-c = .9428$, $-d = .4309\frac{2}{3}$, $+f = .10907\frac{1}{3}$; the equation becomes $-\ddot{y} = 2 + 5.6568x + 5.1716x^2 - 2.1816x^3$: and if we write the values of x and its powers found in *Art.* 180, we get $-\ddot{y} = 4.30616$; and as $\dot{z}$ is here unity, the equation $\dot{z}^2 + vv\ddot{y} = 0$, gives $vv = .2322$, which multiplied by 21214, the value of b , gives 4925.89 feet for twice the height through which a falling body acquires the velocity sought : and if we multiply this height by 32.2, twice that of a body fallen through^r in a second, and extract the square root of the product, we shall have 398 feet moved over uniformly in a second by this velocity.

By a method which will be given hereafter, for finding the velocities independant of the foregoing infinite series, we find $vv = 2328$; which differs therefore by .0006 parts from that above : this shews how nearly the values of x have been found by the former operations.

We have $t = sx + \frac{1}{2}nssx^2 + \frac{1}{3}mxx^3 + \frac{1}{4}pxx^4$ by *Art.* 178, $2ms = 12d - nns^4$, $2ps = 20f - 2mnss$; and because $ss = 2$, $n = 1$, $d = .4309\frac{2}{3}$, $-f = .10907\frac{1}{3}$, we get $m = .4142$, $-p = 1.357$: therefore $t = 1.4142x + x^2 + .1381x^3 - .339x^4$. Now if we

write the values of x and its power found in *Art.* 180, we get $t = .565$. But if this value be multiplied by the square root of b (21214) we shall have the height fallen through in that time, and if that height be divided by the square root of 32.2, twice the height fallen through in a second, we shall have the number of seconds employed in the fall, by $2s = ptt$ in *Art.* 84.

If therefore we divide 21214 by 32.2, and extract the square root of the quotient, we get 25.667, which, multiplied by .565, gives 14.5 seconds for the time required.

P R O B L E M.

184. *To find the velocity with which the shot strikes the ground at K, and the time of describing the whole curve.*

We have — $\dot{y} = 2 + 5.6568x + 5.1716x^2 - 2.1816x^3$, by *Art.* 183; and by writing the values of x and its powers found in *Art.* 179, we get — $\dot{y} = 6.7145$. But the value of z^2 must be known before we can find that of vv : as we have $y = 1 - 2x - 2.8284x^2 - 1.7238x^3 + .5454x^4$, by *Art.* 180; by writing the values of x and its powers found in *Art.* 179, we get — $y = 1.47026$, whose square added to unity gives $z^2 = 3.1609$; and by $z^2 + vv\dot{y} = 0$, we get $vv = .4558$, and this multiplied by 21214, the value of b , gives 9669 feet for twice the height through which a falling body acquires this velocity; or if this number be multiplied by 32.2, the square root of the product, gives 558 feet moved over uniformly in a second by this velocity.

We have $t = 1.4142x + x^2 + .1381x^3 - .339x^4$; and by writing the values of x and its powers in *Art.* 179, gives $t = 1.176$; and this multiplied by 25.667, *Art.* 183, gives 30 seconds for the time required.

It

It is evident that the value 1.47026 of y is the tangent of the angle in which the shot strikes the ground, which therefore is 55 degrees and 47 minutes.

P R O B L E M.

185. *Let the projectile velocity be half the greatest possible, and the rest the same as before, to find the horizontal range A K.*

As $n = \frac{1}{2}$ by supposition, we have $c = .4714$, $d = .0488$, $-f = .0626$; and therefore the general equation becomes $y = x - xx - .4714x^3 - .0488x^4 + .0626x^5$, and when $y = 0$, then $1 = x + .4714x^2 + .0488x^3 - .0626x^4$. Hence by the general rule for finding the roots 1 : 1 : 1.4714 : 1.9916 : 2.6714 : 3.6194 : 4.8845 : 6.5966 : 8.9087 : 12.03, and $x = .7405$ nearly. For if the value of x and its powers are wrote into the equation, we get $1 = 1.00000$.

P R O B L E M.

186. *To find the abscissa A B, corresponding to the greatest ordinate B D.*

The fluxion of the last equation, when $y = 0$, gives $1 = 2x + 1.4142x^2 + .1952x^3 - .3131x^4$, and finding the root of this equation, as usual, we get $x = .4$ nearly.

In searching the velocity at the highest point D, I found $vv = .3179$, and by the method given hereafter $vv = .3177$; so that the difference is .0002 only; which shews how nearly the value of x has been found by the infinite series.

P R O B L E M.

187. *Let the projectile velocity be but one tenth of the greatest, and the rest as before, to find the horizontal range A K.*

Because

Because $n = .1$, by supposition, the general equation becomes here $y = x - x^2 - .09428x^3 + .0169x^4$; and when $y = 0$; then $1 = x + .09428x^2 - .0169x^3$; hence we get $1 : 1 : 1.09428 : 1.17166 : 1.25793 : 1.3499 : 1.4487 : 1.5562 : 1.6699$; and $x = .9318$ nearly.

P R O B L E M.

188. To find the abscissa A B, corresponding to the greatest ordinate B D.

By making $y = 0$, in the former equation, we get $1 = 2x + .2828x^2 - .0676x^3$, and $1 : 2 : 4.2828 : 9.0636 : 19.2032 : 40.68$; hence $x = .472$ nearly. For $x^2 = .22278$, $x^3 = .10515$; and these values wrote into the equation give $1 = 1.000$.

C O R.

189. Hence, if these values of x and its powers are wrote into the equation found in *Art.* 187, we get $y = .2405$ nearly for the greatest ordinate.

And these values wrote into $-y = 2 + .5656x - .2028x^2$, we get $-y = 2.2178$. Whereas, by the method hereafter, it should be $-y = 2.22956$; which therefore differs by .0117.

P R O B L E M.

190. Let the angle of elevation be 30 degrees, and the projectile velocity the greatest possible, to find the range A K.

Here we have $a = .57735$, $s = 1.1547$, by the tables of sines, $ss = \frac{4}{3}$, and $n = 1$, by supposition; hence $y = .57735x - \frac{2}{3}x^2 - .5132x^3 - .168x^4 + .0745x^5$. Now if we make $y = 0$, and divide by .57735, we get $1 = 1.1547x + .8889x^2 + .29098x^3 - .129x^4$; whence we get $1 : 1.1547 : 2.2222 : 3.8833 : 6.6664 : 11.6472 : 20.2181 : 35.1379 : 61.0748 : 106.1377$. and $x = .57543$ nearly

nearly. For the values of x and its powers wrote into the equation give $1 = 1.00008$.

If this number be multiplied by 21214, the value of b , we get 4069 yards for the horizontal range A K.

PROBLEM.

191. *Let the angle of elevation be 60 degrees, and the rest as before, to find the horizontal range A K.*

By the tables of sines we have $a = 1.732$, $s = 2$, and as $n = 1$ by supposition, we get $b = 2$, $c = \frac{8}{3}$, $d = 1.512$, $f = 1.428$; and therefore $y = 1.732x - 2x^2 - \frac{8}{3}x^3 - 1.512x^4 + 1.428x^5$. Now if $y = 0$, and we divide by 1.732, we get $1 = 1.1547x + 1.5396x^2 + .873x^3 - .8245x^4$. The root being found in the usual manner gives $x = .49323$, very nearly: for the values of x and its powers wrote into the equation give $1 = 1.000003$. This number, multiplied by 21214 feet, gives 3487.7 yards for the horizontal range, short by 581 yards to that when the angle of elevation is 30 degrees.

PROBLEM.

192. *Let the inclination of the plane A N, either above or below the horizon, be given, to find the point N, where the shot will hit that plane.*

Suppose p the tangent of the angle N A Q, then will $Q_N = y = px$, and this value of y wrote into the equation of Art. 175, with a positive sign when the point N is above the horizon, or negative when below, gives $\pm p = a + bx + cx^2 + dx^3 + fx^4$; and if $p = .1$, $n = 1$, $ss = 2$, we get $.9 = x + .9428x^2 + .43095x^3 - .10907x^4$. If both sides are divided by .9, we shall have $1 = 1.1111x + 1.0475x^2 + .4788x^3 - .1212x^4$. Hence $1 : 1.1111 : 2.282 : 4.1782 : 7.4436 : 13.6262 : 24.6707 : 44.7425 : 81.178 : 147.2255$, and $x = .55138$

.55138 nearly. For the values of x and its powers wrote into the equation give $1 = 1.00016$.

If the object is below the horizon and the rest as before, then $-p = .1$, and the equation will then be $1 = .90909x + .85709x^2 + .39177x^3 - .0991x^4$, and $1 : .90909 : 1.68353 : 2.70142 : 4.15596 : 6.6631 : 10.51112 : 16.62728 : 26.32379 : 41.64033 : 65.89056 : 104.1574 : 164.871 : 260.8211 : 412.70435 : 653.0166 : 1033.24188$, and therefore $x = .632$; for if the values of x and its powers are wrote into the equation, give $1 = 1.000000$.

P R O B L E M.

193. *The angle of elevation and the horizontal range being given, to find the ratio n of the greatest velocity to the projectile one.*

If we suppose $y = 0$, in *Art. 175*, we get $a = bx + cx^2 + dx^3 + fx^4$; and $2b = ss$, $3c = ns^3$, $12d = 2nns^4 - ans^3$, $60f = 4n^3s^5 - 8an^2s^4 + ns^3$; these values, being wrote into the equation, give

$$a = \frac{1}{2}ssx + \frac{1}{3}ns^3x^2 + \frac{1}{6}nns^4x^3 + \frac{1}{15}n^3s^5x^4 \\ - \frac{1}{12}ans^3x^3 - \frac{2}{15}annx^4 \\ + \frac{1}{60}ns^3x^4$$

Now when the angle of elevation and the range x are given, the value of n may be found in the usual manner.

E X A M P L E.

194. *Let the angle of elevation be 45 degrees.*

Then as $a = 1$, $ss = 2$: we get

$$1 = x + .9428nx^2 + \frac{2}{3}nnx^3 + .37712n^3x^4 \\ - .2357nx^3 - \frac{8}{15}nnx^4 \\ + .04714n^3x^4.$$

Now if the range $x = .59262$, then $x^2 = .351198$, $x^3 = .208127$, $x^4 = .12334$; and these wrote into the

the equation give $1 = .59262 + .28787n + .07297n^2 + .04651n^3$; which gives $n = 1$ nearly: for we get $1 = .99996$. wanting 00004 only.

EXAMPLE.

195. Let the angle of elevation be 45 degrees, and the horizontal range $x = .7405$; then as $x^2 = .54834$, $x^3 = .40604$, $x^4 = .30067$. These values wrote into the equation of the last article give $1 = .7405 + .43544n + .11034n^2 + .113388n^3$, whose root is $n = .5$; for $n^2 = .25$, $n^3 = .125$, $n^4 = .0625$; and these values wrote into the equation give $1 = 1.0000$.

These two examples agree exactly with the *Art.* 179, 185, as it must happen; but when the value of n cannot be guess'd at, the value of x must be subtracted from unity, and both sides divided by the difference, then the values of n may be found by the general rule for finding the roots.

PROBLEM.

146. *The ratio n of the velocities and the horizontal range being given, to find the sine of the angle of elevation.*

If in *Art.* 175, we make $y = 0$, the equation there found becomes $a = bx + cx^2 + dx^3 + fx^4$, and $b = \frac{1}{2}ss$, $c = \frac{1}{3}ns^3$, $d = \frac{2nns^4 - anz}{12}$, $f = \frac{n^3s^5}{15}$, &c.

Now if z denotes the sine and u the cosine of the angle of elevation, then will $a = \frac{z}{u}$, $s = \frac{1}{u}$, by the

property of the circle; and hence $b = \frac{1}{2uu}$, $c = \frac{n}{3u^3}$,

$d = \frac{2nn - nz}{12u^4}$, and $f = \frac{n^3}{15u^5}$. These values being

wrote into the equation, and the whole multiplied by u , give $z = \frac{x}{2u} + \frac{nx}{3uu} + \frac{2nn - nz}{12u^3}x^3 + \frac{n^3x^4}{15u^4}$.

S

Whence,

Whence, if we suppose $x = az + bz^2 + cz^3 + dz^4$, and we write the value of x and those of its powers, we get $x = 2uz - \frac{8nu^2}{3} + \frac{40nnuz^3}{9} -$

$$\frac{1088n^3uz^4}{135} + \frac{4nnuz^4}{3}. \text{ But if } a \text{ denotes the sine of}$$

the angle double that of elevation, then will $a = 2uz$, and $a = 2z - z^3 - \&c.$ by the property of the circle; by writing these values into the last equation

$$\text{it becomes } x = a - \frac{2naa}{3} + \frac{5nna^3}{9} - \frac{68n^3a^4}{135}. \text{ From}$$

$$\text{hence we get } a = x + \frac{2nx^2}{3} + \frac{nnx^3}{3} + \frac{2n^3x^4}{15}.$$

N. B. This series is not exact, and gives the sines only nearly, as the following examples will shew; thus, if $n = 1$, and $x = .59262$, then will $a = .91257$, which should be unity by *Art.* 179; the difference being .08743.

Again, let $n = 1$, and $x = .57543$, then will $a = .8660$; but should be 8743 by *Art.* 190; and therefore the difference is .0083 only.

With all the skill I may have in infinite series, I could not make this to answer exactly; it would require two terms more at least, which makes the calculation so very operose and intricate, that it is hardly worth while to attempt it: besides in practice too much exactness is not required.

Second METHOD of solving the same Problem by the differences.

PROBLEM.

197. *The ratio of the greatest and projectile velocities being given, together with the angle of elevation, to find the velocity at the highest point D.*

Because

Because $q \frac{2z}{r} = b\ddot{y}$ by Art. 177, and $\dot{z} = \sqrt{1 + \dot{y}^2}$; now if the first side be multiplied by $\dot{z}$, and the second by its equal $\sqrt{1 + \dot{y}^2}$, we get $\dot{z} q \frac{2z}{r} = b\ddot{y} \sqrt{1 + \dot{y}^2}$; and if $\frac{1}{2}m$ denotes the fluent of $\ddot{y} \sqrt{1 + \dot{y}^2}$; we shall have $r q \frac{2z}{r} = 2A + bm$ for the fluent by the remark we made after Art. 80; but at the vertex D, where $z = y = 0$, and $q \frac{2z}{r} = 1$, we get $\frac{r}{2} = A$, and therefore $r q \frac{2z}{r} = r + bm$, or $q \frac{2z}{r} = 1 + \frac{bm}{r}$; this value being wrote into the square $vvq \frac{2z}{r} = b\dot{z}^2$, gives $vv \times 1 + \frac{bm}{r} = b\dot{z}^2$. From whence it appears that the ratio of the velocity at the vertex, and any other point where the angle made by the tangent and abscissa is known, is given.

C O R.

198. Hence, we may deduce the values of $\dot{z}$ and $\dot{y}$; for if we write the value $1 + \frac{bm}{r}$ of $q \frac{2z}{r}$ into $q \frac{2z}{r} = b\ddot{y}$, we get $1 + \frac{bm}{r} = b\ddot{y}$, or $1 = \frac{b\ddot{y}}{r + bm}$, or $\dot{z} = \frac{rb\ddot{y}}{r + bm}$, or else multiplying both sides by $\dot{y}$, $\dot{y} = \frac{rb\dot{y}\ddot{y}}{r + bm}$.

It is to be observed that the equations just found are for the descent from the highest point D; but when the body ascends, then m becomes negative;

consequently $vv \times 1 - \frac{bm}{r} = b\dot{z}^2$, $\dot{x} = \frac{rb\ddot{y}}{r-bm}$ and $\ddot{y} = \frac{rb\ddot{y}}{r-bm}$ in this case.

C O R.

199. Since we have $\dot{z} = vi$ and $v \dot{x} q \frac{z}{r} = c\dot{z}$ by Art. 177, we get $\dot{x} q \frac{z}{r} = ci$; or because $q \frac{z}{r} = \sqrt{1 - \frac{mb}{r}}$ and $\dot{x} = \frac{r\ddot{y}}{r-bm}$, we get $ci = \frac{r^{\frac{1}{2}}\ddot{y}}{\sqrt{r-bm}}$.

E X A M P L E.

200. Let the angle of elevation be 45 degrees, and the projectile velocity at the point A equal to the greatest possible, to find the velocity c at the highest point.

Then $vv \times r - \frac{bm}{r} = b\dot{z}^2$ gives $b = \frac{vv}{\dot{z}^2 + nm}$; because m is here negative and $\frac{vv}{r} = n$. Now as $\dot{z}$ expresses the secant at A, we have $\dot{z} = 2$, and $n = 1$, whence $b = \frac{vv}{2+m}$; but the fluent of $\ddot{y}\sqrt{1+\dot{y}^2}$ gives $m = \dot{z}\dot{y} + \log. \dot{y} + \dot{z}$, because $\dot{z} = \sqrt{1+\dot{y}^2}$; or because $\dot{y} = 1$, $m = \dot{z} + \log. 1 + \dot{z} = 2.2956$; we have $b = \frac{vv}{4.2956}$. Hence, in an 18 pounder, $vv = 21214$ feet, by Art. 182; consequently, $b = 4938.6$ feet.

Before we apply this, to find the values of $\dot{x}$ and $\ddot{y}$, it will save much trouble to insert the following table, where the values of m are under that letter, according to those supposed of $\dot{y}$ under it against each other.

other. The use of this table will appear in the following

EXAMPLE.

201. To find the abscissa A B corresponding to the greatest ordinate B D.

Because $b = \frac{vv}{4.29559}$, or as $vv = r$ by supposition $b = \frac{r}{4.29559}$; and if this value be wrote into $\dot{x} = \frac{rb\ddot{y}}{r-bm}$, we get $\dot{x} = \frac{r\ddot{y}}{4.29559-m}$.

y	m	y	m	y	m	y	m
1	2.29559	$\frac{1}{2}$	1.04027	$\frac{3}{5}$	1.26642	$\frac{1}{8}$	0.25077
$\frac{1}{3}$	0.67872	$\frac{3}{4}$	1.63064	$\frac{4}{5}$	1.75720	$\frac{3}{8}$	0.76722
$\frac{2}{3}$	1.42638	$\frac{1}{5}$	0.40265	$\frac{1}{6}$	0.33500	$\frac{5}{8}$	1.32727
$\frac{1}{4}$	0.50688	$\frac{2}{5}$	0.82091	$\frac{5}{6}$	1.84322	$\frac{7}{8}$	1.95340

Now if the space of which $\dot{x}$ is the fluxion be divided into six equal parts by 7 ordinates, then the parts of the base y will be $0 : \frac{1}{6} : \frac{2}{6} : \frac{3}{6} : \frac{4}{6} : \frac{5}{6} : \frac{6}{6}$, or $0 : \frac{1}{6} : \frac{1}{3} : \frac{1}{2} : \frac{2}{3} : \frac{5}{6} : 1$; by subtracting the corresponding values of m from 4.29559, and dividing unity by the difference, we get .23279 : .25248 : .27648 : .30719 : .34853 : .40776 : .5. Now making the sum of the first and last equal to A, that of the second and last but one equal to B, that of the third and last but two C, and the fourth D, then $A = .73279$, $B = .66024$, $C = .62501$, $D = .30719$; and writing these values into the form $41A + 216B + 27C + 272D$, found in Art. 76, for

840

seven ordinates, and dividing the Sum 273.0872, by the denominator 840, give $\dot{x} = .32511$ for the number

number required ; which exceeds .32456, found in *Art.* 180, by .00055 parts.

EXAMPLE.

202. *To find the greatest ordinate BD.*

Since $y = \frac{r\ddot{y}}{4.29559 - m}$, from whence it appears, that if we multiply the ordinates found above each by the respective values of $\ddot{y}$, we get $0 : .04288 : .09216 : .15359 : .23233 : .33980 : .5$. Hence $A = .5$, $B = .38188$, $C = .32451$, $D = .15359$. These values wrote into the form above, and the sum 153.5243, divided by 840, give $y = .1827$; that found in *Art.* 181, is .1826,

EXAMPLE.

To find the part BK of the range.

In order find this part, the tangent of the angle at K must be known first ; to effect which, let ab be drawn parallel to the base BK, so as the tangent y at b be unity ; then by means of the equations

$$\dot{x} = \frac{r\ddot{y}}{4.29559 + m}, \quad y = \frac{r\ddot{y}}{4.29559 + m}, \quad \text{we are to find}$$

Da and ab : now if we suppose the base y divided into five ordinates, then will $0 : \frac{1}{4} : \frac{1}{2} : \frac{3}{4} : 1$: be the values of y ; and dividing unity by the sum of 4.29559, and the values of m , corresponding to those of y , we get .23279 : .20822 : .18741 : .16874 : .15171; and hence $A = .3845$, $B = .37696$, $C = .18741$. These values wrote in the general form $\frac{7A + 32B + 12C}{90}$, give $ab = .188923$.

Now if each of these ordinates be multiplied by the respective values $0, \frac{1}{4}, \frac{1}{2}, \frac{3}{4}, 1$, we get $0 : .05205 : .0937 : .12655 : .15171$, hence $A = .15171$, $B = .1786$, $C = .0937$; these values, when

when wrote into the general form above of five ordinates, give $y = aD = .08779$.

Now the question is to find such a value for y , that the space found by means of the equation $y = \frac{rj\ddot{y}}{4.29559 + m}$, may be equal to Ba ; for this value being once obtained, it need only be added to unity, which is that at b , to have that at K , by which the difference between BK and ab may be found.

We have found $BD = y = .1827$, in *Art.* 202, and $aD = .08779$ just now, whence $Ba = .0949$; which being divided by the last ordinate .15171, found above, gives .6256, which must be greater than the value of y required, which we shall suppose .47; whence if the space .47 be divided into two equal parts .235, then $1 : 1.235 : 1.47$ will be the parts corresponding of the base: then the values of m , when $y = 1.235$, and 1.47 , will be 2.0005, 3.79, which being added to 4.29559, and unity divided by the sum, give .15171 : .15898 : .12376 : and these values wrote into the general form $\frac{A + 3B}{6}z$, for three

ordinates, give .1519; which, multiplied by the base .47, gives $Ba = x = .07139$, which being less than .0949, shews the value 1.47 for y is too small, and hence a greater must be assumed; and if we add this to 816, we get $BK = .2603$ instead of .2680, so that the difference is .0077, so small it would be needless to extend the explanation of this method any farther. The reader must be sensible, that the latter part is attended with great difficulties, the trials to be made for finding the tangent of the angle, in which the body strikes the ground, are so very tedious and operose, that this method will hardly ever be used, at least the latter part; as to the first,

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it may be applied to good purposes, and the rest supplied by what we shall say in the

THIRD METHOD.

PROBLEM.

204. *The projectile velocity at the vertex D, being given, to find the abscissa AB corresponding to the greatest ordinate BD.*

We have $y = xx + cx^3 + dx^4 + fx^5 + gx^6 + kx^7$, by *Art.* 177; and $b = \frac{1}{2}$, $3c = n$, $6d = nn$, $6of = 4n^3 + n$, $18og = 4n^4 + 7nn$, $315k = 2n^5 + 16n^3 - \frac{3}{8}n$. It is to be observed, that this series was continued further than the former, because all the terms, which contain the odd powers of x , become negative in the ascent of the body, and n expresses the ratio of the greatest and projectile velocity at the vertex D.

Now since $b = \frac{vv}{4.29557}$, by *Art.* 200; and $\frac{b}{r} =$

$\frac{1}{4.29559} = .23273$, because $r = vv$ by supposition; and $n^2 = .05419$, $n^3 = .01261$, $n^4 = .00293$, $n^5 = .00068$. Hence $b = \frac{1}{2}$, $c = .07759$, $d = .00903$, $f = .00472$, $g = .00217$, $k = .00037$; these values wrote into the equation before give $y = \frac{1}{2}x^2 + .07759x^3 + .00903x^4 + .00472x^5 + .00217x^6 + .00037x^7$; and the fluxion of this equation will be $\dot{y} = x + .23279x^2 + .03613x^3 + .0236x^4 + .01304x^5 + .00194x^6$. By means of these two equations all the several parts of the projectiles may be determined; but as the value of y is greater than unity, the series will not converge without that of x being diminished, which is done in the shortest manner by writing $2x$ instead of x ; we shall then have $\dot{y} = 2x + .9312x^2 + .289x^3 + .3776x^4 + .4173x^5 + .1242x^6$: by this means the value of x is diminished

diminished to one half. It must be observed that every other term in these series become negative in the ascent.

EXAMPLE.

205. *The angle of elevation at A being in the case above 45 degrees.*

The fluxion y being the tangent when x is the radius, and therefore $y = x = 1$, by supposition, we get $1 = 2x - .9312x^2 + .289x^3 - .3776x^4 + .4173x^5 - .1242x^6$, in the ascent; and by the method of finding the roots we have $x = .697$; for this value wrote into the equation gives $1 = 1.00513$.

Now $2x = 1.394 = AB$, when multiplied by b , and as $b = 4938.6$ feet by *Art. 200*, that abscissa will therefore be 6884.4 feet; and as we have found 6885 feet in *Art. 182* for that distance, the difference is only .6 of a foot. Now, to find the greatest ordinate BD, we must write the value 1.394 of x , and that of its powers, into $y = \frac{1}{2}x^2 - .07759x^3 + .00903x^4 - .0047x^5 + .00217x^6 - .00037x^7$, which give $y = .78283$; and this value, multiplied by $b = 4938.6$, gives $BD = 3866$ feet; which is less by 7.6 feet than that found in *Art. 182*.

EXAMPLE.

206. *Let it be required to find the remainder BK of the horizontal range.*

Because $y = .5x^2 + .07759x^3 + .00903x^4 + .00472x^5 + .00217x^6$, and $y = .78282$; if we divide by this value of y we get $1 = .638708x^2 + .0991x^3 + .0115x^4 + .006x^5 + .002776x^6$: the square root of this equation is $1 = .7992x + .062x^2 + .0024x^3 + .0007x^4$; whose root is $x = 1.1514$. For by writing this value of x , and those of its powers into the equation, give $1 = 1.00728$ nearly; which shews that this value is some very small matter more than the real one.

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If this value 1.1514 of x be multiplied by 4938.6 feet, that of b , we get 5686.3 feet for BK; and as we have found 5687 feet for that distance, in *Art.* 182, the difference is not above one foot.

It is plain that $y = x + .2327x^2 + .0961x^3 + .0276x^4$ expresses the tangent of the angle in which the shot strikes the ground, and as $y = 1648$, that angle will be 58 degrees and 45 minutes, which exceeds that found in *Art.* 184 by about two degrees.

We have thus shewn how the several problems of this kind may be solved by three different methods, all sufficiently accurate for practice, wherein no mathematical exactness is required more than to come as near as it can be done upon the ground; for as the batteries and the object are seldom upon an exact level, the gunner goes by guess only, and when the object is either much higher or lower, the difference is only judged at by the eye: and when the shot or shell falls short or exceeds the mark, the powder is increased in one case and diminished in the other.

There are many other problems given by authors; such as to find the ratio of the resistance to the force of gravity, so that the body may describe a given curve, or to find the curve described by the body, when the resistance is to its gravity in any other ratio than that of the squares of the velocities; but as there are no other laws in nature, as far as I can find, and serve only to puzzle the reader, and to lose his time in useless inquiries, we have not mentioned them, and refer such readers as are fond of needless speculations to authors who wrote professedly on them.

SECTION

SECTION VIII.

The true Figure of the Earth determined from actual Measures.

SINCE Sir Isaac Newton has first determined the figure of the earth, in his Principles of Natural Philosophy, most mathematicians of the first class have made it their principal study. Many works have since appeared upon this subject without agreeing upon the ratio of the axis to the diameter of the equator, notwithstanding the measures of a meridian degree made since in different latitudes, and the great number of experiments made upon the length of pendulums vibrating seconds in different parts of the world.

Sir Isaac considering the earth as a fluid body, and upon this supposition, that any meridian section must be an ellipsis, otherwise the particles of matter could not remain at rest amongst each other in the revolution of the earth about its axis, and besides that there can be no other figure in nature which has that property; from thence, and the measure of a meridian degree in France, found that the ratio of the axis to the diameter of the equator was as 229 to 230.

But finding afterwards that, by some experiments on pendulums made in the West Indies, this ratio must be something greater; and yet most authors make it less, as Maclaurin, Thomas Simpson, and Clairaut; as to Bouguer, he makes it much greater than it should be.

It is for this reason, I have attempted to find this ratio more accurately, by means of some properties of the ellipsis, and the two best measures of a meridian degree, near the pole and the equator, and

found it to be as 215 to 216. That this ratio is the nearest possible, expressed by three figures, appears by its agreeing as nearly as can be with the actual measures and the experiments made on the length of pendulums vibrating seconds in different latitudes.

Had authors compared their theories with the actual measures and experiments, they would have found, that their ratios were not sufficiently exact; which is the best and most certain method in philosophical inquiries. The French have spared no pains or expence to determine this question so useful in geography and navigation, and though they have succeeded beyond expectation in their measures near the pole and at the equator, the case is quite different with those in France; for there are no less than five different measures of the same degree, &c. 57060, 57292, 57183, 57074, and 56926 toises: sometimes it is said to be the measure of a meridian degree at Paris, and at others the mean meridian of France. In short, these measures are so uncertain, that it is not safe to apply them to any particular meridian, for want of knowing where they begun or where they ended; for Clairaut, Maupertius, and Bouguer do not mention any thing about it in their works.

In what follows we shall suppose, with other authors, that all celestial bodies attract each other in the ratio of their quantities of matter directly, and the squares of their distances inversely: this Sir Isaac has proved to be so in all the planets, and from thence that this law was universal.

T H E O R E M. Fig. 57.

207. *A particle of matter placed on the surface of a sphere A M Q D is attracted towards the Center C, by a force proportionally to the radius C M.*

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For the quantity of matter in a sphere of an uniform density, is as the cube of its radius ; and since the attraction is as the quantity of matter directly, and the square of the distance inversely, the attraction of a particle of matter placed on the surface of a sphere is therefore, as the cube of the radius directly and the square inversely ; that is, as the radius directly.

This has likewise been proved by Sir Isaac, in *Prop. 72. Book I.*

C O R.

208. Hence, if the attraction should cease at a certain distance CI from the center, or be destroyed by another force in an opposite direction, it is evident, that there would be a hollow within of a concentric sphere.

That there is a repelling force in nature, is confirmed by many experiments ; but whether it reaches so far as the center of bodies, is not certain : it is true the attractive force seems to reach so far, and by a parity of reason the repelling force may do the same.

T H E Q U E R E M. Fig. 58.

209. *If a sphere, whose particles are freely disposed, turns round one of the diameters A a as an axis, with such a force as may be compared to that of attraction, the sphere will be changed into a spheroid.*

Let PM be perpendicular to the axis A a, and reduce the force of attraction, in the direction CM, into the perpendicular PM and the parallel CP, to the axis A C ; then the force of attraction at M, in the direction MP, will be to the force of attraction at D, in the direction DC, as PM is to CD ; and since the centrifugal forces in different circles, in equal times, are as their radii, *Art. 114*, the centrifugal force at M is to the centrifugal force at D,

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are likewise as PM to CD. Now as the centrifugal force acts in a contrary direction to that of gravity, it will diminish it in proportion as it approaches to the equator CD : whereby the sphere will be protracted at the poles, and rise towards the equator, otherwise the particles of matter could not remain in a state of rest.

Let MN be to DB as the centrifugal force at M to the centrifugal force at D, then will $PM : CD :: MN : DB$, and by composition $PM : CD :: PN : CB$. Which is the property of the ellipsis.

C O R.

210. It is manifest that the spheroid becomes prominent at the equator in proportion as the centrifugal force increases, or, which is the same, as the periodic time decreases, till the centrifugal force becomes equal to that of gravity, and then the particles of matter would remain suspended between the two equal forces ; but if the centrifugal force should prevail the least, they would fly off without end.

C O R.

211. It is no less evident, that if the particles of matter were of different densities, it would not change the nature of the figure. For in the formation of the body, the particles would naturally dispose themselves so as to be at rest with each other, and the heaviest would by their greater force approach the nearer to the center ; and no other change could happen, than that the spheroid would be denser near the center than at the surface: the centrifugal force must produce the same effect in both cases.

C O R.

212. If there was an empty space about the center of the sphere at rest, when that sphere is changed into a spheroid by its revolution about its axis, the hollow sphere will at the same time be changed into a spheroid.

a spheroid similar to the outward one. For the centrifugal force, as well as that of attraction, will act in the same manner upon every particle of matter, whether the spheroid was hollow within or not.

R E M A R K.

It is highly probable that all the planets are hollow within, from this general maxim in philosophy, that nothing has been created in vain, and more is in vain when less will do; and since after a certain thickness of matter near the surface, there seems to be no more necessary, at least as far as we can judge; yet as men can reason but from appearances, and they are in favour of an emptiness within the planets, there is no absurdity to suppose they are so in effect. It appears likewise that all bodies formed by the mutual attraction of the particles will assume the form a sphere, if no impediment intervenes and they do not turn about an axis, this appears from drops of water and quicksilver: but if these bodies revolve about an axis, they will become spheroids, more elevated at the equator than at the poles: this is evident in fluids, and even when some parts are only so, as has been found by observation in all the planets, although it is not certain whether they contain any fluid matter or not.

T H E O R E M.

213. *The direction of gravity on the surface of planets is every where perpendicular to the surface.*

Experience shews it to be so in regard to the earth, and it is manifest from the laws of hydrostatics, that the direction of gravity in all fluids must be perpendicular to the surface, otherwise the particles could never be at rest amongst each other: it is no less evident, that the direction of gravity is the same every where else, since bodies could not remain at rest after their fall. For suppose *A M*, *Fig. 54*,

to be a part of the surface, CM the direction of gravity; then it is manifest, that if the angle CMT, made by the direction CM and the tangent MT to the surface at M, be less than a right angle, the body, instead of remaining at rest after the fall, would move in the direction MT, which is contrary to experience. As all planets revolve about their axes, and the attraction acts upon them in the same manner, and by the same law, as upon the surface of the earth, it must evidently follow that the same causes produce the same effect.

N. B. Several authors endeavour to find the direction of gravity by a long and tedious investigation, without having regard to the laws of hydrostatics, and some pretend, that it deviates from the perpendicular; this arises from their supposing the spheroid to be at rest: and as this is contrary to fact, their conclusion is evidently erroneous.

T H E O R E M. Fig. 59.

214. *The attraction towards the center of a particle, placed on the surface of a planet, is as its distance from the center.*

Because three powers or forces remain in a state of rest, when they are to each other as the respective sides of a triangle formed by their directions: if Aa be the axis, CB the radius of the equator, MK perpendicular to the tangent at M, and CM a semidiameter; then because MK is the direction of gravity at M by the last, CK that of the centrifugal force at that point, and CM that in which the attraction tends towards the center: therefore CK is to KM as the centrifugal force is to the force of gravity, and KM is to CM as the force of gravity is to its part by which it tends towards the center. This has been proved by Sir Isaac in the 3d Cor. after Prop. 91, Book I, from different principles.

C Q R.

C O R.

215. Hence, if MK be produced till it meets the conjugate diameter CN to CM in L; then if the centrifugal force, expressed by CK, be reduced into two others CL, LK, the first perpendicular and the second parallel to the direction MK of gravity, ML will express the total force of gravity, and MK the part by which it acts on a particle placed at M, after being diminished by LK or the centrifugal force. This has likewise been proved by Clairaut, p. 188. If the gravity in any given latitude be known, the total gravity ML and the centrifugal force CK may be found.

C O R.

216. When the point M coincides with the point B in the equator, the perpendicular MK becomes equal to the parameter of the semi-axis CB, by the property of the ellipsis: hence, if a expresses the force of gravity at B, then will $CB : CA :: CA : a$; and when the point M coincides with the point A, then MK becomes equal to CA; that is, if p expresses the gravity at the pole, then will $CB : CA :: p : a$; which shews, *that the force of gravity at the pole is to the force of gravity at the equator reciprocally as the radius of the equator is to the radius of the axis.* This is likewise proved from statics.

C O R.

217. If PM be drawn perpendicular to the equator CB; then because the centrifugal force CK at M is to the centrifugal force at B, as CP is to CB, by the property of the ellipsis; or as the circumference described by the point M is to the circumference described by the point B. This agrees with what has been proved in *Art. 114*; and therefore confirms what has been proved in the *Articles 213, 214, 215, and 216.*

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These are the only elementary propositions necessary to determine the true figure of the earth from observations, which being deduced from the nature of the subject, and some properties of the ellipsis, are so plain and easy as to leave not the least doubt of their certainty; which is not so with most authors, who make use of fluxions and infinite series: for if the least mistake happens, it is difficult to find it out in long and tedious algebraic computations, the deducing a wrong direction of gravity is a convincing instance, besides many others equally as absurd.

P R O B L E M.

218. *The ratio between the actual gravities in two different latitudes being given, to find the ratio between the axes of the ellipsis.*

Let $CB = 1$, $CA = b$, $CP = u$, $PM = y$, and $1 - bb = n$; then will $yy = bb \times 1 - uu$ be the equation of the ellipsis, and $PK = bbu$; and as the sum of the squares of KP and PM is equal to the square of MK , we have $MK^2 = bb \times 1 - uu + bbuu$. But to find such a value for MK as may be applied to practice, let r be the sine and s the cosine of the latitude MKB , the radius being unity; then will $PK (bbu) : PM (y) :: s : r$, or $sy = bbur$, and writing the value of y into the equation $yy = bb \times 1 - uu$, we get $bbrruu = ss - uuss$, and this value of uu into that of MK^2 , gives $MK^2 = \frac{b^4}{bbrr + ss}$, when reduced under the same denomination, and unity wrote for its equal $rr + ss$, or because $1 - bb = n$, $bb = MK \sqrt{1 - nrr}$.

C O R.

219. Since, when the point M coincides with point B in the equator, then will $r = 0$, and the last equation becomes $bb = MK$; hence, if p is to a as the

the gravity at any point M is to the gravity at the equator, we get $a : p :: bb : \sqrt{1-nrr}$, or $a = p \sqrt{1-nrr}$; which shews, that the force of gravity in any latitude M is to the force of gravity at the equator, as unity is to $\sqrt{1-nrr}$, when r denotes the sine of that latitude. From whence the ratio of the axis to the diameter of the equator may be determined.

C O R.

220. Since $\sqrt{1-nrr}^{1-\frac{1}{2}} = 1 + \frac{1}{2} nrr$, nearly, the last equation $a = p \sqrt{1-nrr}$ gives $\frac{1}{2} anrr = p - a$, and as $\frac{1}{2} an$ is constant in the same sphere, the difference $p - a$, between the gravities in any latitude M and at the equator, is as the square of the sine r of that latitude nearly. From whence follows the common rule given first by Sir Isaac, that the increase of gravity from the equator towards the pole, is as the square of the sine of the latitude nearly.

From this Corollary it appears, that the decrease of gravity near the equator is much greater than towards the poles, because the sines increase much faster at the beginning of the quadrant than near the end; for which reason, experiments made on gravity near the equator are much more uncertain than farther off; because the errors in point of latitude are more there than elsewhere: as pendulums which vibrate in equal times, and by which the ratio between gravities are discovered, their exact lengths are extremely difficult to be found, as will appear from those two in the latitudes of $51^{\circ} 32'$ and $53^{\circ} 11'$; for their difference is no more than .004 parts of an inch. It will therefore be more convenient to find that ratio by the measures of two meridian degrees in two different latitudes, as far distant from each other as can be; for which reason we shall use

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those measured at the equator and near the pole, which at the same time are the most exact of any we know of.

P R O B L E M.

221. *From the given ratio between the force of gravity at Paris and at the equator, to find the ratio between the axis and the diameter of the equator.*

Mr. Clairaut says, in page 193, that a body descends through a space of 2174.11 lines in a second at Paris, and in the next page, through a space of 2168.75 lines at the equator; these numbers divided by 5, in order to reduce them to a less denomination, gives $p = 43482$, and $a = 43375$. Now as Paris is in the latitude of $48^{\circ} 50'$, according to Clairaut, we have $r = .752798$, or $rr = .5667$ nearly; and these values wrote into the equation $\frac{1}{2}anrr = p - a$, Art. 220, give $n = \frac{1 \frac{070}{132803}}$: from whence we get $1 - n = bb = \frac{131733}{132803}$, whose square root is $b = \frac{5185}{5208}$, after having been divided by 7.

Hence, the nearest ratio between the axis and diameter of the equator, expressed by three figures, is that of 248 to 247, according to these gravities; and not that of 230 to 231, as Maclaurin, Simpson, and Clairaut would have it, by Clairaut's own principles.

P R O B L E M.

222. *From the known meridian degrees in two different latitudes, to find the ratio of the axis to the diameter of the equator.*

Because similar circular arcs are as their radii, and the single degrees of an ellipsis are sensibly equal to a degree of the circle described by the radius of curvature; the degree of a meridian at any point M is to the meridian degree at any other point as the radius of curvature at M is to the radius of curvature at the other point. But the radius of curvature at the point M is equal to the
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cube of the perpendicular MK, divided by CA^4 , *Art.* 42, and since $bb = MK \sqrt{1 - nrr}$, *Art.* 218, the radius of curvature at M will be expressed by $bb \times \sqrt{1 - nrr}^{-\frac{3}{2}}$; for the same reason, if s denotes the sine of any other latitude, $bb \times \sqrt{1 - nss}^{-\frac{3}{2}}$, will express its radius of curvature.

C O R.

223. Hence, as the sine s at the equator vanishes $bb \times \sqrt{1 - nss}^{-\frac{3}{2}}$, becomes bb : therefore the meridian degree at the equator is to the meridian degree at M as bb is to $bb \times \sqrt{1 - nrr}^{-\frac{3}{2}}$, or as $\sqrt{1 - nrr}^{-\frac{3}{2}}$ to unity. Hence, the meridian degrees at the equator and in any latitude being given, the ratio of the axis to the diameter of the equator will be found, and on the contrary, the ratio of the axis to the diameter of the equator being given, that of the meridian degrees may be determined.

E X A M P L E.

224. Suppose $b = \frac{230}{231}$, according to Clairaut and others, then will $1 - bb = n = \frac{461}{53361}$, and as the sine of the angle $66^\circ 20'$, which answers to the middle of the meridian degree measured by Maupertuis and others, is $r = .91589$, and its square $rr = .83885$; this value and that of n wrote into $\sqrt{1 - nrr}$ give .99637, whose cube is .98914 nearly; and as the meridian degree, in the latitude $66^\circ 20'$, has been found to be 57838 toises, see page 162, Maupertuis, we have by *Art.* 223 this proportion $1 : 98914 :: 57438 : 56814$ toises for the meridian degree at the equator, which exceeds the measured one 56753 by 61 toises.

If we suppose $b = \frac{229}{230}$ with Sir Isaac, we shall find 56811 toises, which exceeds the measured one by 58 toises: if we take 57422 toises for the meridian degree at the pole, as corrected by Bouguer, we shall find

find 56798 toises, which exceeds the measured one by 45 according to Clairaut, and 42 according to Newton.

P R O B L E M.

225. *From the known meridian degrees at the equator and near the pole, to find the ratio of the axis to the diameter of the equator.*

Because $\sqrt{1 - nrr}^{\frac{3}{2}} = 1 - \frac{3}{2}nrr$ nearly by the binomial rule, we have by the proportion in *Art.* 223, $57838 : 56753 :: 1 : 1 - \frac{3}{2}nrr$, and because $rr = .83885$, by *Art.* 224, we get $n = \frac{1}{144545} \frac{370}{45}$ and $1 - n = bb = \frac{1}{144545} \frac{43175}{45}$, or dividing by 5, $bb = \frac{28635}{38909}$, whose square root is $b = \frac{169218}{170026}$.

The nearest ratio expressed by three figures is greater than that of 211 to 210, and less than that of 632 to 629; but nearer to the last. This is true upon the supposition, that the measures of the meridian degrees are exact; but as this ratio does not agree with the observations on pendulums, there must be some small error or other: if the correction mentioned by Bouguer, page 290, be admitted, which makes the degree near the pole 57422 toises only, we shall find in the same manner as before, that $n = \frac{1}{144505} \frac{338}{45}$ and $bb = \frac{1}{144505} \frac{43167}{45}$, whose square root gives $b = \frac{378374}{380138}$.

Hence the nearest ratio between the axis is less than that of 214 to 215, and greater than that of 215 to 216; but nearest to the last.

That this ratio is the nearest to the true one, of all those found by other authors, appears from its exact agreement with all the measures of a meridian degree in different latitudes, and with all the best experiments made on the lengths of pendulums vibrating seconds, as the following examples will shew.

EXAMPLE.

EXAMPLE.

226. *The meridian degree at the equator being given, to find that near Pellow in the latitude of $66^{\circ} 20'$.*

Since $b = \frac{215}{46656}$ by the last, or $bb = \frac{46225}{46656}$, we get $n = \frac{431}{46656}$; and because $rr = .83885$, by *Art.* 224, we get $\sqrt{1 - nrr} = .99611$, whose cube is .98837, and by the proportion in *Art.* 223, we have .98837 : 1 :: 56753 : 57421 toises for the meridian degree at Pellow, which wants but one toise of the corrected one.

EXAMPLE.

227. *The meridian degree at the equator being given, to find that at Paris in the latitude of $48^{\circ} 50'$.*

We have $rr = .5667$, by *Art.* 221, and as $n = \frac{431}{46656}$, by the last, we get $\sqrt{1 - nrr} = .99738$, whose cube is .99216, and by the proportion in *Art.* 224, we have .99216 : 1 :: 56753 : 57201 toises for the meridian degree at Paris, which exceeds the corrected one 57183 by 18 toises.

That the length of the meridian degree at Paris, we have just now found, is pretty exact, appears from the measure of gravity at Pellow and at Paris, made by Maupertuis, which he says is as 1.0014 to unity, page 226, English edit. As Pellow is in the latitude of $66^{\circ} 48'$, whose sine r is $= .91913$, and $rr = .8448$ nearly, hence we get $\sqrt{1 - nrr} = .996$ nearly, the cube of which is .988, and by the proportion in *Art.* 223, .988 : 1 :: 56753 : 57442 toises for the degree in that latitude. And as the meridian degrees are as the cubes of gravity, *Art.* 222, the cube 1.0042 of 1.0014 is to unity, as 57442 is to 57201 toises, for the meridian degree at Paris, which is exactly the same as that computed before; consequently, the true meridian degree of
Paris

Paris cannot differ from that here computed more than in the last figure.

EXAMPLE.

228. *The meridian degree of the equator being given, to find that at London in the latitude of $51^{\circ} 32'$.*

The sine of this latitude is $r = .78297$, and $rr = .61304$ nearly, and hence $\sqrt{1 - nrr} = .997164$, whose cube is $.99152$; and therefore $.99152 : 1 :: 56753 : 57238$ toises for the meridian degree at London.

The meridian degree measured by Mr. Norwood from London to York is different from that here found; for the difference between these places is $2^{\circ} 28'$, half of which, added to $51^{\circ} 32'$ of London, gives $52^{\circ} 46'$ for the middle latitude of that distance, whose sine is $r = .796178$ and $rr = .6339$ nearly; hence $\sqrt{1 - nrr} = .99707$, whose cube is $.99123$; and therefore $.99123 : 1 :: 56753 : 57255$ toises for the mean meridian, which is less than 57300, that found by Norwood, by 45 toises; and as this difference is greater than any error that could proceed from any inaccuracy of computation, it must have been an error of measurement: that this is the case appears from the next example, where the length of a degree is less, though $4^{\circ} 15'$ more northerly.

EXAMPLE.

229. *The length of a meridian degree at the equator being given, to find that at Edinburgh, in the latitude of $55^{\circ} 47'$.*

The sine of that latitude is $r = .82692$, and $rr = .68379$, and hence $\sqrt{1 - nrr} = .99684$, whose cube is $.99055$, and therefore $.99055 : 1 :: 56753 : 57294$ toises, for the meridian degree required.

EXAMPLE.

EXAMPLE.

230. *The meridian degree at the equator being given, to find that at Dublin in the latitude of $53^{\circ} 11'$.*

The sine is here $r = .800557$, and $rr = .64089$, and so $\sqrt{1 - nrr} = .99703$, whose cube is $.99112$; and therefore $.997112 : 1 :: 56753 : 57261$ toises for the meridian required.

Having shewn how to find the meridian degree in any given latitude, in a more concise and exact manner than any author has done before, which is very useful in correcting geography, and in navigation, it remains now to shew how to find the length of pendulums vibrating seconds in any latitude, in as correct a manner as can be expected, and by which it will be confirmed how exact the ratio of the axis to the diameter of the equator has been determined.

PROBLEM.

231. *The length of a pendulum vibrating seconds in any latitude being given, to find the length of a like pendulum in any other latitude.*

We have shewn, in *Art. 107*, that the lengths of pendulums vibrating seconds in any latitude are as the force of gravity in that latitude, and in *Art. 219*, that these forces are inversely as $\sqrt{1 - nrr}$, when r denotes the sine of the latitude; and therefore at the equator, unity is to $\sqrt{1 - nrr}$ as the length of a pendulum vibrating seconds, in the latitude whose sine is r , to the length of a like pendulum at the equator.

EXAMPLE.

232. Mr. De Mairan found by a great number of very accurate experiments, that the length of a pendulum vibrating seconds at Paris is 440.567 lines French measure: a line is the twelfth part of an inch,

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and

and as $\sqrt{1 - nrr} = .99738$ in that latitude, by *Art.* 227, we have $1 : .99738 :: 440.567 : 439.4127$ lines for the length of a pendulum vibrating seconds at the equator.

Sir Isaac has found this length to be 439.468, as may be seen in his table, page 247, Eng. edit. but he observes afterwards, page 250, that the earth is more elevated at the equator than he had computed it; the length of his pendulum must of course be shorter: on the contrary, Bouguer makes it 439.21 lines only, page 342; but we shall shew hereafter for what reason he differs so much in his measure from other authors.

E X A M P L E.

233. *The length of a pendulum vibrating seconds at the equator being given, to find the length of a like pendulum at Pellow in the latitude of $66^{\circ} 48'$.*

We have $\sqrt{1 - nrr} = .996$ nearly, by *Art.* 227, and therefore $.996 : 1 :: 439.4127 : 441.177$ lines nearly, by *Art.* 231. This agrees with that found by Maupertuis, p. 226, Eng. edit.

E X A M P L E.

234. *The length of a pendulum vibrating seconds at the equator being given, to find the length of a like pendulum at London in the latitude of $51^{\circ} 32'$.*

We have $\sqrt{1 - nrr} = .99716$, by *Art.* 228, and hence $.99716 : 1 :: 439.4127 : 440.6642$ lines for the length required. Now as the French foot is to the English foot as 114 is to 107, this length will be 469.4927 lines, English measure; or dividing by 12, 39.1244 inches, which wants but .0006 parts of an inch of the length 39.125 found by Dr. Halley, from several accurate experiments.

The same length would have been found, if it had been deduced from that at Paris.

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The near agreement between the lengths of pendulums deduced from theory, with those found by experiments, shews how nicely the experiments must have been made, and how accurately we have determined the ratio between the axis and the diameter of the equator.

The length of a pendulum vibrating seconds at London, said to have been found by Sir Isaac, is 39.207 inches, which is absolutely erroneous. For when converted into French measure gives 441.594 lines; whereas that at the pole is only 441.387, as determined by Sir Isaac, as may be seen in his table, page 247.

EXAMPLE.

235. *The length of a pendulum vibrating seconds at the equator being given, to find that of a like pendulum at Edinburgh in the latitude of $55^{\circ} : 47'$.*

We have found $\sqrt{1 - nrr} = .9968$, in Art. 229; and hence $.9968 : 1 :: 439.4127 : 440.809$ lines for the length required; which, when reduced into English measure, gives 39.137 inches.

EXAMPLE.

236. *The length of a pendulum vibrating seconds at the equator being given, to find that of a like pendulum at Dublin in the latitude of $53^{\circ} : 11'$.*

Here we have $\sqrt{1 - nrr} = .99703$, by Art. 230; hence $.99703 : 1 :: 439.4127 : 440.7216$ lines for the length required; which, being reduced into English measure, gives 39.1294 inches.

EXAMPLE.

237. *The length of a pendulum vibrating seconds at the equator being given, to find that of a like pendulum at the pole.*

Here we have $\sqrt{1-nrr}=b=\frac{215}{216}$, and therefore $\frac{216}{215} \times 439.4127 = 441.4564$ lines for the length required; which exceeds that at the equator by 2.0437 lines.

P R O B L E M.

338. *To find the length of the radius of the equator, and that of half the axis.*

The radius of curvature at the equator B, added to the value of CK, when the point M falls on the point B, gives the radius CB of the equator. Now since the meridian degree at the equator is 56753 toises, this, multiplied by 180, gives 10215540 toises for half the circumference; which, being to the radius as 355 to 113, gives 3251707 toises for this radius.

But when the point M falls on the point B, the radius of curvature MK is to the value of CK as bb to n , by *Art.* 219; and since $bb=\frac{46225}{46656}$, and $n=\frac{431}{46656}$, by *Art.* 226, we have $46225 : 431 :: 3251707 : 30319$ for the value required; which, being added to the radius of curvature at the equator, gives 3282026 toises for the radius CB of the equator, and the semi-circumference 10310789 toises, and one degree 57282 toises; which therefore exceeds the meridian degree 56753 by 529 toises.

As CB is to CA in the ratio of 216 to 215, we get 3266830 toises for half the axis CA. An English mile is 1760 yards or 829 French toises, the radius of the equator will be 3959.01 English miles, half the axis 3940.68 miles, and their difference 18.33 miles.

Sir Isaac found this difference 17.1 miles, page 243; but observes, in page 250, that the earth is a little higher at the equator than what he had computed it. This shews that our difference agrees with Sir Isaac's; whereas Maclaurin, Simpson, and Clairaut, make it less, and Bouguer much greater.

P R O B L E M.

P R O B L E M.

239. *To find the force of gravity at the equator.*

We have found that a body falls through a space of 15.097 in a second at Paris, by *Art.* 107, and that the force of gravity at Paris is to the force of gravity at the equator as unity to $\sqrt{1 - nrr} = .99738$, *Art.* 227; hence $1 : .99738 :: 15.097 : 15.057446$ feet, for the height fallen through at the equator in the same time: but if we suppose the ratio of the axis to be as 230 to 231, then will this height be 15.059967 feet.

P R O B L E M.

240. *To find the centrifugal force at the equator.*

Let r denote the circumference of a circle whose diameter is unity, a the radius of the equator, and t the time of one revolution; then will $2ar$ express the circumference of the equator; which, divided by the time t , gives $\frac{2ar}{t}$ for the arc described in a second, when t is expressed in seconds. Now $r = 3.14159$; and as $t = 86164$ seconds, and $a = 3282025$ toises; we get $\frac{2ar}{t} = 239.3293$ toises for the arc described in a second, and the square 57278.513849 of that arc, divided by the diameter 6564050, and multiplied by 6 feet, gives .052356 feet for the versed sine of that arc, or the space described by the centrifugal force in a second. And since the force of gravity describes a space of 15.059967 feet, by the last *Art.* at the same time, the centrifugal force is to the force of gravity as .052356 to 15.059967, or as unity to 287.597.

If we suppose the ratio of 230 to 231 of the axis, then the radius of the equator will be 3280044 toises; and by the same process as before, the gravity
at

at the equator will be to the centrifugal force at that place as 287.81 to unity.

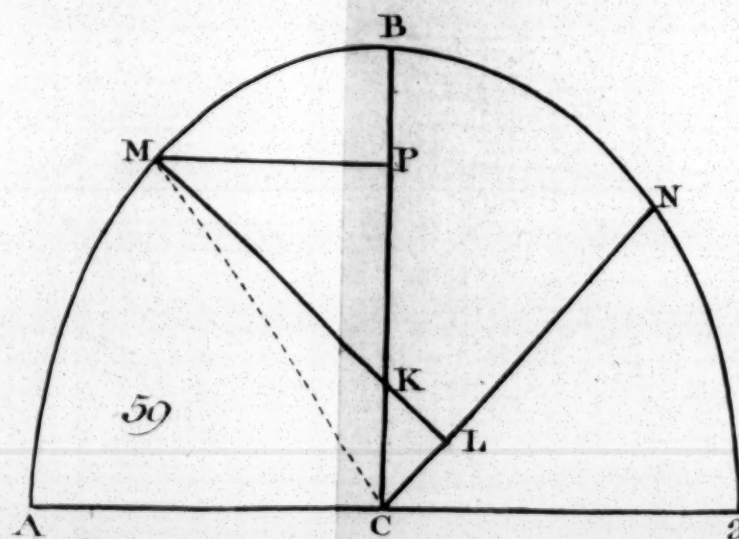
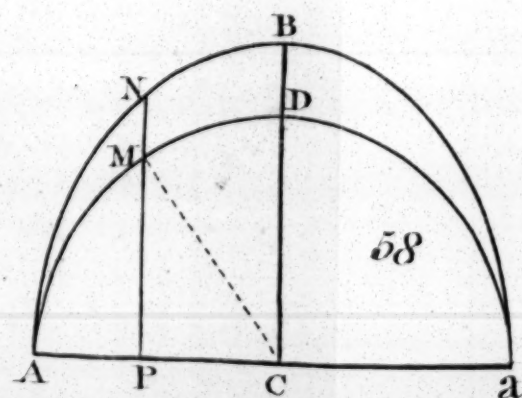
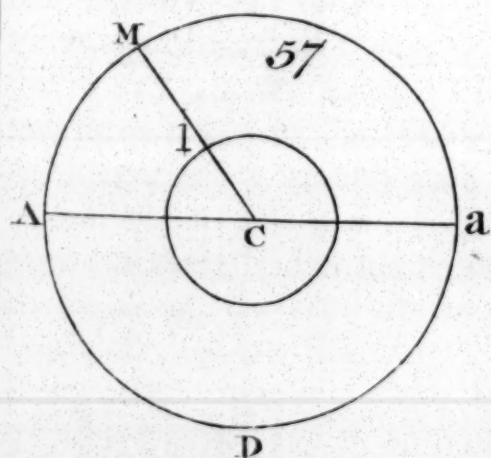
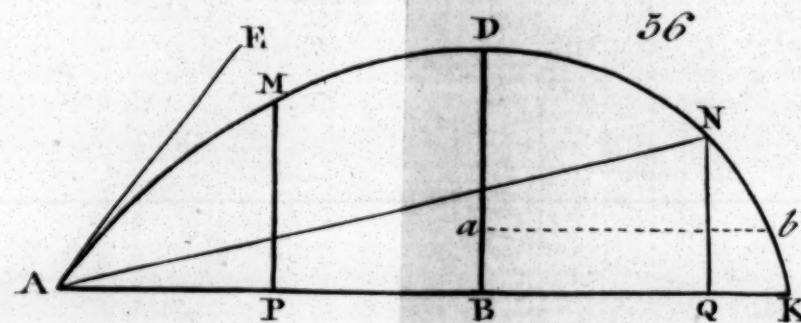
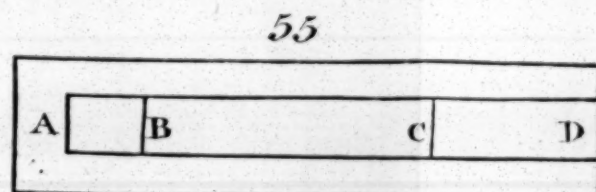
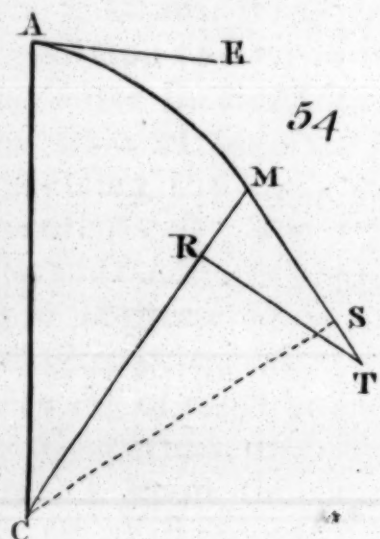
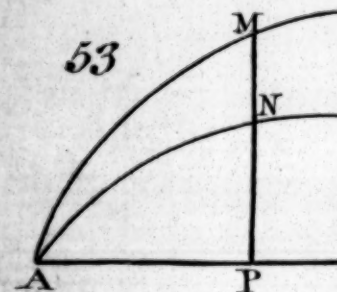
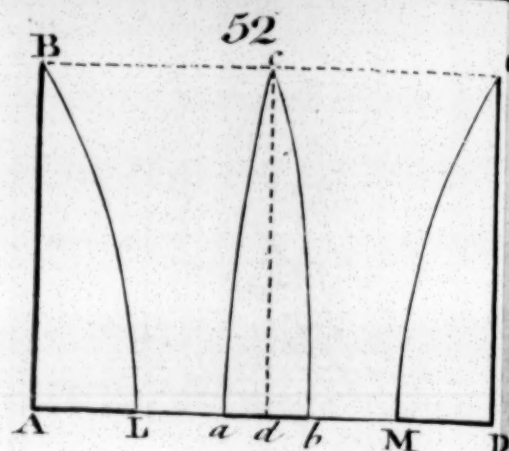
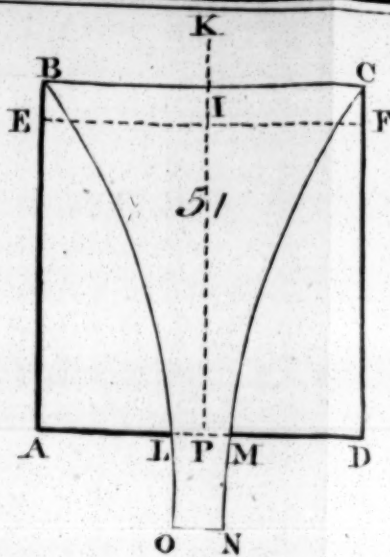
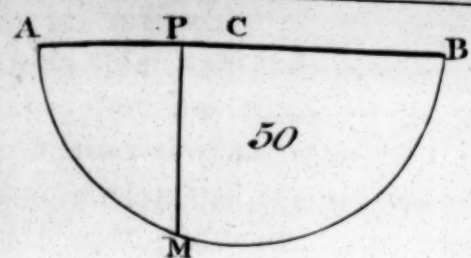
And if we admit Bouguer's ratio of 178 to 179, the gravity at the equator will be 15.04914 feet, the radius of the equator 3288345 toises, and the gravity at the equator to the centrifugal force as 286.885 to unity.

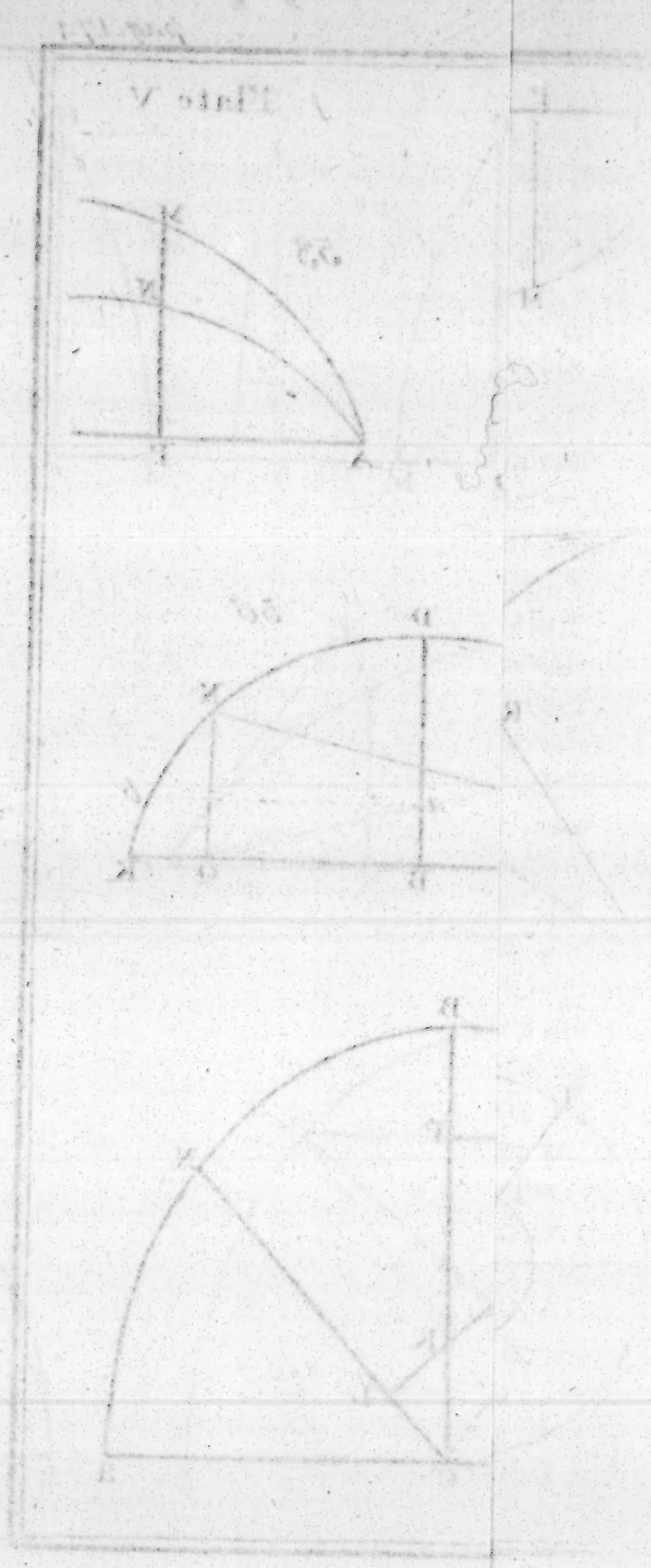
Observations on the authors who have wrote on the figure of the earth.

Maclaurin, Thomas Simpson, and Clairaut have found the ratio of the axis to the diameter of the equator to be as 230 to 231, which is less than that of Sir Isaac Newton, without considering whether theirs agreed better with the several measures of a meridian degree, or the most accurate observations on the length of pendulums, which yet is necessary to ascertain the truth of all such kind of speculations; it is true, they have made use of the most sublime part of the mathematics in order to shew their great skill and knowledge in computations, without considering that knowledge deduced from the most simple principles is least liable to mistakes, and more satisfactory; though they bestowed so much pains in their theory and algebraic calculations, they are in my opinion not sufficiently exact in their numeral approximations: for Simpson, in page 472, finds the ratio of the axis to the diameter to be as 1151 to 1156, which he says is as 230 to 231; whereas, by the common rule of approximation, it is less than that of 231 to 232, and greater than that of 232 to 233, but nearest to this last.

Clairaut, page 192, finds this ratio to be as 2304 to 2314, which he makes as 230 to 231; whereas it is greater than this, and less than that of 231 to 232; but nearest of all as 461 to 463: this author, finding the above ratio not to be exactly that of 230

Plate V





to 231, by supposing the ratio of the centrifugal force, and that of gravity at the equator, and to find this last he assumes the ratio of the axis, whereby he finds it to be as unity to 287.5; granting that the ratio of the axis is right, and his calculation just, then this last ratio is true; but it does not prove that the ratio of the axis is true by any means: had he found the ratio of the centrifugal force to that of gravity, without supposing that of the axis, the last must have been true also: a pleasant way of arguing backwards and forwards from one unknown thing to another, and at last draw a conclusion to have arrived at truth. This author has no better succeeded in page 188, where he says that the ratio of the axis of 230 to 231 is true, upon supposing the earth to be uniformly dense, and in no other case; and if it is not so, that ratio must be less and could not be greater: whereas we have found it to be greater, that is, as 215 to 216, and shewn it to agree better with the measures and observations than any of those hitherto given by others.

Maupertuis, in page 162, English edit. makes a remark upon the degree measured by Picard in France, and reduces it from 57060 toises to 56925.7, the difference is no less than 134.3 toises; which is still less by 366 toises than that of Picard's 57292 toises, published in 1701, and less by 257 toises than that 57183, corrected. This disagreement of the measures made in France by the most considerable mathematicians, creates such a confusion as renders it impossible to judge the right from the wrong; had they not been more successful in the north and near the equator, the true figure of the earth would have been less known than ever: what is surprising, Sir Isaac came nearer to the truth than all together, without having any good measure to depend upon.

We

We now come to Bouguer's work, the last published, which one should think would decide the question beyond all exception; but he makes the equator higher than it is, whereas the others make it too low. In page 296, he says, that by taking the measures at the equator and near the pole for granted, he found the ratio of the axis to be as 215 to 214; but was obliged to change his opinion, when he heard that the meridian degree near Paris had been corrected by the gentlemen who had been employed near the pole, and made it 57183 toises, whereby he found the ratio of 223 to 222, in which he was confirmed by the measure of a degree of longitude in the latitude of $43^{\circ} 32'$; which was found to be 51618 toises by measure, and 41607 by computation.

In page 299, he changes again when he heard that Picard's measure, instead of being 57183 toises, had been corrected again to 57074; from whence, in page 304, he concludes that the true ratio of the axis was that of 178 to 179. Now as his first ratio, deduced from the same datas as we did, came out nearly the same as ours, because we found that the true one was between that of 214 to 215, and of 215 to 216; but nearer the last: had he remained satisfied, he would have been right; but the various measures of a meridian degree in France, on which he depended too much, was what misled him.

To shew that his ratio is not true; the difference between the squares of 179 and 178 gives $n = \frac{357}{32041}$, by *Art.* 224, the square of the sine of $66^{\circ} 20'$, is $rr = .83885$, by the same *Art.* and hence $\sqrt{1 - nrr} = .99532$, whose cube is .98602; whence .98602 : 1 :: 56753 : 57557.6 toises for the meridian degree near the pole; which exceeds the measure 57438 by 119.6 toises, and if his own correction of that degree be allowed the difference will be 1361 toises.

Again, in the latitude of Paris, we have $rr = .56699$, by *Art.* 227, and hence $\sqrt{1 - nrr} = .99683$, whose cube is 99032; therefore $.99052 : 1 :: 56753 : 57296$ toises for the degree at Paris, which exceeds the measure corrected 57074 by 222 toises.

This evidently proves his ratio of the axis is the farthest of all from the true one: it does not appear from his work, that he so much as thought of comparing his ratio with any of the measures, no more than the observations with pendulums; for the length of a pendulum vibrating seconds at the equator is, according to his ratio, 439.17 lines; which not only disagrees with all others, but his own observations.

With regard to the experiments, he says, the weight of his pendulum was one ounce five drams and two scruples, suspended by a fibre of an aloe leaf, and fastened above by a pincher screwed into a wall; the form of the weight was two equal frustums of a cone joined by their greatest bases, whose center of oscillation, he says, was rather above the middle, because of the rod's weight, 10 toises of which weighed 4 grains; so that 3 feet weighed $\frac{1}{3}$ of a grain.

Now, whether so small a weight causes the center of oscillation of the bob to rise from below the center above it, is not very probable; the rod of the pendulum being fixed above cannot move so freely as about the center, let it be ever so supple; besides a small weight is more resisted in proportion in its motion than a heavy one. As his length is less by .2 of a line, or $\frac{1}{60}$ part of an inch, than it should be, which can be owing to nothing else than so imperfect an experiment; besides he judged by the eye of the length, and a ruler placed by the rod. But rather than allow of the imperfection of his experiment, he supposes that the universal law of gravity does not hold good under the equator.

T H E E N D.

A D V E R T I S E M E N T.

ALL spurious Lists must be very deficient and erroneous; the *Piracy* of the Army List, with the application for authority, in opposition to the only authentic authority, is absolutely the most audacious, invidious invasion on property that ever was attempted; and I hope the universe will never produce eighteen more booksellers who would be concerned in such a scandalous insult upon trade.

With the same conscience they would, if possible, strip mankind, even one another.

This is your Liberty and Property! — Booksellers, Plunderers, fomenters of defamation, sedition, treason, and blasphemy, and the very grave of liberty! O grief of griefs!

By means of such daring desperadoes the liberty and commerce of most nations have suffered.

Soon may be Published,

1. The *Filcheian Political Register*, shewing by what Artifice and a pretended Shew of Opposition, with the Assistance of Grubstreet Garreteers, and his own equal Genius, he amuses the Unwary with Scandal and Treason, void of Sense or Reason, as all his Publications hitherto have been.

2. The Last Dying Speech and Confession, Birth, Parentage, and Education, of the most notorious *Filch*, and his Chum *Sbylock*, with their Gang of piratical Booksellers, Stationers, &c. decorated with the Effigies of that motely Crew.

Tyburn will never have its due,

Whilst it groans for them and some of you.

Only think of an abandon'd Wretch behind his Counter, indiscriminately robbing and retailing, with Impunity, Defamation against the Innocent, and Treason against the Best of Kings. Who has most to answer, the needy or the greedy Thieves?

My Advertisements have been splot by means of those piratical Booksellers, to whom the least Offence has never been given by me.

Just published, for J. MILLAN.

1. A List of his Majesty's Land Forces and Marines, at Home and Abroad, &c. for 1768. By Permission, 5s.
2. General Wolfe's Instructions to Young Officers; also his Orders for a Battalion and an Army, &c. 2s.
3. Muller's Works, of Fortification, Engineering, Mining, Artillery, Mathematics, &c. &c. containing upwards of 200 cuts, 8 vol. 2l 10s. or any volume separate.
4. Manœuvres for a Battalion of Infantry upon fixed Principles, with 57 plates, 10s 6d.
5. New Exercise, by his Majesty's Order, 1s.
6. Recruiting Book proper for all Officers on that Service, 2s 6d.
7. Regimental Book, with proper Heads, beautifully engraved, 4l 4s.
8. General Returns for reviewing Horse, Dragoons, and Foot.

ARTILLERY DICTIONARY.

AFFUT, the French name of a gun carriage, for which we have no proper name; the only distinction from all other carriages, is that it belongs to a gun.

AMMUNITION, comprehends powder, shot, shells, carcasses, and every necessary in the field for the use of fire-arms.

AMMUNITION WAGGON, is a four-wheel carriage with shafts, and having the sides railed in with raves and staves, lined with wicker work, so as to carry bread and all sorts of tools. See plate xx.

As the fore horses draw two a-breast, that in the shaft cannot go in the same tract, it continually treads on one side or other the ridge left between the two; or if the road is quartered it shuns to tread in the rut, by which the wheels can never follow the same tract; for which reason it would be more convenient to have two pair of shafts, as waggons that carry great loads.

AMMUNITION CART, is a two-wheel carriage with shafts, and the sides as well as the fore and hind parts are inclosed with boards instead of wicker work.

APRON, a small thin piece of lead about six inches square, which is laid over the vent of a gun or mortar when loaded, till they are fired, to prevent the priming from dampness.

ARSENAL, a large building that contains store-rooms for every thing used in artillery.

ARMS, the two round ends of an axle-tree. See axle-tree.

ARTILLERY, comprehends cannon, mortars, howitzes, petards, shot, shells, carcasses, and all their appurtenances used at home and in the field.

ARTILLERY, REGIMENT, is composed of three battalions (and a company of cadets, of which the master-general is always captain) each of them commanded by a colonel, lieutenant-colonel, and a major; the master-general of the ordnance is commander in chief; the lieutenant-general, commander en second; and the three

A

colonels

colonels are called colonel-commandant each of his battalion; each company contains matrosses, gunners, and bombardiers.

ARTILLERY, EQUIPAGE, is a quantity of guns, mortars, shot, and shells, with all the necessary stores made out for a campaign, or an expedition by sea or land.

ASTRAGAL, a moulding, in architecture, like a round ring with a fillet on each side, used in guns and mortars. See guns.

AXES are of two sorts, one for felling wood, and a pick-axe for picking stones or digging trenches when the ground is too hard for the spade.

AXIS, in mechanics, is that line or beam about which any thing turns.

AXLE-TREE of a carriage about which the wheels turn, the middle part between the wheels is called the body; the two ends, which are round and enter into the naves, the arms.

They are made of dry young elm, beach, or oak.

AXLE-TREE BAR, in travelling carriages, is a flat iron bar that reaches from one end to the other, and let into the wood underneath quite flush.

AXLE-TREE BOLTS is an iron bolt passing through the middle of the axle-tree to bolt the axle-tree bar.

AXLE-TREE BANDS, used in field carriages, they go round the ends of the body, and are bolted to the under-side of the side-pieces of the carriage.

AXLE-TREE HOOPS, iron rings that go round the ends of the arms to fasten the iron bar, when there are any, and to secure the linch-pin holes.

AXLE-TREE STAYS, are used in ship or garrison carriages, made of iron in the form of an S to fasten the fore axle-tree firm to the under-side of the side-pieces, to prevent their giving way in the sudden jerk occasioned by the recoil when the gun is fired.

BAGGAGE WAGGONS are common waggons railed in with raves and staves, and have bales to cover the top, with strong canvasses to keep the baggage from wet.

BAGS, SAND, they are 16 inches diameter and 30 high, filled with earth or sand to repair breaches and the embrasures of batteries, when damaged by the enemies fire, or by the blast of the guns.

BAGS are also used on many occasions to hold small shot, flints, &c.

BALLS, CANNON, formerly made of stone but now of iron, they are fired out of cannons against an enemy, or against their works made for cover.

BALLS, FIRE, made of an iron shell, or paper, filled with a slow composition, to throw into the enemies works at night to see what they are doing, and to enable the gunners to point their guns on them.

BANNERS, or the ordnance flag fixed upon the fore part of the drum-major's kettle-drum carriage.

BARRACKS, a building to lodge soldiers in.

BARRELS, POWDER, are about 16 inches diameter, and 30 or 32 long, and hold 100 pounds.

BARREL, BUDGE, holds from 40 to 60 pounds; at one end is fixed a leather bag with brass nails: they are used upon batteries to keep the powder from firing by accident, for loading the guns and mortars.

BARROWS, to carry loads from one place to another.

BARROWS, HAND, serve for carrying loads by two men.

BATTERY, to place guns or mortars on in a siege; it consists of a breast-work or parapet of 20 or 22 feet thick, and 7 or 8 feet high, with embrasures to fire the guns through, behind which is a platform of wood 18 feet long, 8 feet broad before, and about 16 behind, sloping about six inches forwards in order to prevent the guns from recoiling too much, and bringing the guns easier forward when loaded.

They commonly allow 18 feet for the distance from gun to gun, and when the battery is seen on the sides by the enemies works they add flanks of 10 or 12 feet thick, from 15 to 18 long, and as high as the parapet. They are called by different names from the several manners of firing.

BATTERY, BARBET, when the breast-work is only three feet high, so that the guns may be traversed any way, as are commonly those near the sea or large rivers, to fire at the ships as they are coming, pass by, or passed; these batteries are also used upon the faces of the bastions in a siege while the besiegers are as yet at a distance, and because embrasures would help to make the breach in the faces sooner than without them.

BATTERY, RICOCHET, which are placed about 100 yards before the first parallel, perpendicular to the faces of the works, to enfilade them by giving the guns a small elevation to go over the opposite parapet, and loaded with a small quantity of powder.

This method of firing was invented by M. Vauban, and first put in practice at the siege of Aeth in 1697.

BATTERY MORTAR, they have a breast-work as before, without embrasures; the inside slope of it ought to be such that two pickets might be placed, one above and one in the middle, to mark the direction of the object. The inside of these batteries are sunk two or three feet into the ground, by which they are much sooner made than those of cannon.

BATTERING PIECES serve to make breaches; it is judged by most nations that no less than a 24 pounder is proper for that use; formerly much larger calibers were used, but as they were so long and heavy, became very troublesome to carry and manage, they have been rejected; but if larger calibers can be made so as to be worked and managed with ease, there is no doubt but that they will make a breach sooner than a 24 pounder; this we have proved to be very possible, and what is more of iron instead of brass, as they have hitherto been made.

BAYONET, a broad dagger without any guard, generally made with a round hollow handle, to fix to the muzzle of a musket, in which manner it serves instead of a pike to receive the charge of horse or foot.

BEDS,

BEDS, MORTAR, they are made of solid timber, consisting generally of four pieces fastened together with strong iron bolts and bars; those of the royal and coehorns are made of one solid block only.

BEDS, SEA MORTAR, these are likewise made of solid wood like the former, but differ in their form, having in the middle a hole to receive the pintle or strong iron bolt about which they turn.

BED, STOOL, is a piece of wood to lay the breech of a gun upon in a truck-carriage, with another piece fixed to it at the hind end, that rests upon the body of the hind axle-tree, and the fore part is supported by an iron bolt.

BOLSTER, is that piece of wood fastened upon the axle-tree of any wheel carriage.

BOLT, made of iron, to bolt wood or iron together; they receive additional names from the different parts of a carriage where they are used.

BOLT, AXLE-TREE, those which pin the hind axle-tree of a truck-carriage to the side-pieces.

BOLT, EYE, the fore bolt of a gun-carriage, or of a mortar-bed, to lock the cap-square by means of a key.

BOLT, JOINT, is that to which the cap-square is fixed by means of a hinge.

BOLT, TRANSOM, is that which goes through the transom of a gun-carriage to fasten the side pieces firmly to the transoms.

BOLT, BED, that which passes through the side of a truck-carriage, and supports the fore-end of the stool bed.

BOLT, BREECHING, with rings, used in ship-carriages to fasten the tackles which serve to stop the recoil of the gun, and to advance it when loaded.

BOLT, BOLSTER, that which fastens the bolster of a limber, or of any waggon to the bed under it, or to the axle-tree.

BOLT, BRACKET, those which fasten the upper and under parts together of the cheeks of a truck-carriage.

BOLT,

BOLT, STOOL BED, there are two in each to fasten the stool to the bed, riveted underneath.

BOLT, GARNISH, those which are used in field-carriages, whose heads terminate in a point, and have a neck under it, and fasten the fore part of the axle-tree band by means of a key.

BOMBARDIER, the third rank of private men in a company of artillery.

BOMB-KETCH, a vessel that carries sea mortars, each two, a 10 and a 13 inch one.

BOMB-TENDER, that vessel which carries the necessary stores for the use of the mortars.

To bombard a place, is to fling a great number of shells in it, either to destroy the buildings or to oblige it to surrender.

BORE of a gun or mortar, is the hollow part which receives the powder and shot in the guns, or the shell in the mortar.

Box, they are of several sorts, and for different uses.

BOXES, NAVE, made of iron, and fastened one at each end of the nave, to prevent the arms of the axle-tree about which the boxes move, from causing too much friction.

BOXES, WOOD, with lids, which serve to hold grape shot; each caliber has its own, distinguished by marks of the caliber on the lid.

BOXES, TIN, to fill with small shot for grape, the size of the gun.

BOXES, WOOD, made to slip on and off over the axle-tree of light field carriages, to hold 12 cartridges for firing upon a march in case of an attack.

BREAST, the fore ends of the side-pieces of a travelling-carriage; thus, breast transom, breast transom plates with hooks.

BREECH of a gun, is that solid piece of metal behind, between the end of the bore and the section, through the end of the base ring, or which terminates the hind end of the gun.

BURR,

BURR, a round iron ring which serves to rivet the end of the bolt so as to form a round head like that of the bolt.

BUTTON of the cascable, in a gun or howitz, is the hind part of the piece made round in the form of a ball.

BULLET, of lead for muskets, carabines, pistols, and grape shot.

CALIBER of a gun or mortar, is the same as their bore or opening; and the diameter of the bore is called the diameter of its caliber.

CAMP, is the spot of ground where a strong party of troops, or the army, settle and pitch their tents. The artillery camp is always upon the highest spot before or behind that of the army, to have room for placing their artillery and stores.

CANNON. See guns.

CANTEENS, a tin vessel used by the soldiers to carry their drink in.

CAPS, made of leather for covering the mouth of a mortar when loaded till it is fired, to prevent damp getting into it.

CARRIAGE, a general name for all machines that have wheels.

CARRIAGE, GUNS, the wooden frame upon which guns are fixed; there are three sorts commonly used, viz.

CARRIAGE, FIELD, or travelling, an assemblage of two long planks called side-pieces, joined together by three or four pieces called transoms, one end is supported by two wheels and the other touches the ground.

CARRIAGE, SEA, made of two side-pieces, one transom and bolts, and has four small wheels, each made of one piece of wood called trucks.

CARRIAGE, GARRISON, made the same as the former, with this difference, that their wheels or trucks are made of cast iron.

There has been some of these carriages made entirely of cast iron, and are esteemed very good, especially

cially in hot climates, where the wood is said to rot very soon.

CARTRIDGE, a bag to hold as much powder as is sufficient for loading a gun, they are made of a strong paper called by that name, or of parchment; but as the end generally remains in the bottom of the gun, the paper ones retain fire, and may fire the next, which makes it dangerous; and that made of parchment sticks close to the bottom and accumulates, if not removed, till it stops the vent, which it fills and sticks so fast in it as not to be removed without an iron drill; to prevent those inconveniencies, to the parchment case is fixed a flannel bottom, by which all accidents are avoided; but as the dust of powder may pass through the flannel and thereby cause accidents, a parchment cap is put over it to take off before the cartridge is put into the gun.

CARTRIDGE, for muskets, carabines, and pistols, are made of thin paper, so as to hold half an ounce or less of powder and a bullet.

CARTS, are two-wheel carriages to carry powder, tools, and all kind of ammunition.

CART, SLING, serves to sling a gun under it, and to carry it at a small distance.

CART, POWDER, inclosed round with boards, with lids on the top like a roof, and covered over with an oil-cloth to prevent any dampness coming to the powder.

CART, FORGE, serves to carry an anvil, tongs, hammers, and a poker upon a march, for mending any iron work that may break upon a march, or repair them in the field.

CART or TUMBREL, serves to carry all sorts of tools and other small necessaries to the artillery.

CENTER, the middle part of the side-pieces in a traveling carriage at the bending part.

CIMAISE, a moulding of architecture, applied in our construction of guns, near the mouth of the gun.

COLLAR,

COLLAR, the part of tackle of a horse that goes about the neck, or a belt that goes over the shoulder of a man, and the arm on the other side, when they draw a gun for a little way.

COINS, or **WEDGES**, to lay under the breech of a gun to raise or depress it.

COLOURS, used in the camp, with staves, to mark the front line of the parade, and in marking out the camp.

COLOURS, or **ENSIGNS**, of the various regiments to distinguish them.

The artillery colours are three guns upon their carriages, which are also called the ordnance arms.

CROSSS-BARS, in a limber or waggon, the two wooden bars mortised into the hind part of the shafts.

CROWS, are iron bars, a little bent on one end, and made use of as hand-spikes.

CROWS FEET, an iron of four points of about six inches long, which are used against the cavalry, for one point will always be uppermost, let it fall as it will.

DIMENSIONS, the length, breadth, and thickness of bodies.

DOG-NAILS, have a round flat head, and is itself round in part near the head.

DOLPHINS, are the two handles placed on the second reinforce ring of brass guns, resembling the fish of that name.

DOWLEDGES, iron plates let into the fellows of a wheel on the inside, and fastened with four iron pins, rivetted on the other side.

DOWL BARS, are two iron beams, placed in mortar beds in a vertical position to fasten the upper and under beds together.

DOWL PINS, are made of wood, about three inches long, one in diameter, placed in the middle of the joints of the fellows of a wheel to keep them together.

ther. These pins are also used in ship carriages, two on each side, to keep the side-pieces together.

DRAUGHT HOOKS, fixed on the sides of travelling carriages to the transom plates, three on each side: sometimes a draught hook is also joined to the washer.

DRAUGHT RINGS, fixed to the side-straps of a travelling carriage by means of staples rivetted on the inside.

DRIVER, a wooden cylinder for driving of fuses.

DRUM, a military instrument of music, to beat to arms, or to the several military motions.

DRUM, KETTLE, used by the cavalry as likewise in the artillery.

KETTLE DRUM CARR, of artillery, is a four wheel carriage drawn by four horses; on the fore-part stands the ordnance flag, and on the hind sits a drum-major with two kettle drums as in a chair of state.

EMBRASURE, that opening left in the epaulement for the guns to fire through.

ENGINE, to drive fuses, consists of a wheel with a handle to it, to raise a certain weight, and to let it fall upon the driver, by which the strokes become more equal.

ENGINE, to draw fuses out of shells, has a screw fixed upon a three legged stand; the bottom of this stand has a ring to place it on the shell, and at the end of the screw is fixed a hand screw by means of a collar, which being screwed on the fuze by turning the upper screw draws or raises the fuze.

ENGINE, for driving piles by means of horses, an ingenious invention of an acquaintance of mine for driving the piles at Westminster bridge.

ENGINE, for raising of water, either by fire, as at York Buildings or near Chelsea, others by means of water wheels.

ESSES, FIXED, to draught chains made in the form of an S, one end is fastened to the chain, and the other to hook to the horses harness or to a staple.

EYE BOLTS, see Bolt.

FATHOM, a measure called so, being equal to two yards or six feet, equivalent to the French word *toise*.

FELLOWS, the parts of a wheel which forms its circumference, made of dry elm; there are generally six in each wheel.

FLAGS, with staves, the union and the red, mostly used on board ships.

FLAMBEAUX, used in the artillery, at night, upon a march, or in the park.

FLINTS, for muskets, carabine, and pistols.

FORE-LOCK KEY, the flat iron, broader at one end than at the other, which passes through the eye-bolt to fasten the cap-square, and to secure the trunnions of a gun or mortar.

FORMERS, of different sizes, a cylinder of wood, used in laboratory work, to drive the compositions of fuzes and rockets.

FOUNDRY, for casting guns, mortars, shot, and shells.

FUNNELS, of different sizes, to pour the powder into shells, and the composition into fuze or rocket cases.

GALLOPER, a carriage of a light field-piece with shafts, to be drawn without a limber, but very little used now.

GARNISH BOLT, or nails. See bolts and nails.

GARRISON, the troops quartered in a fortified place to guard and defend it in case of an attack, are called collectively by that name.

GIN, an engine to raise guns composed of three long legs; two of them are kept at a proper distance by means of two iron bars fixed on one of the legs by a staple passing through the hole at one end; the other end has a hook which enters into a staple fixed into the other leg so as to put off or on at pleasure; at

three feet from the bottom is a roller upon which the cable is wound, and the three legs are joined together with an iron bolt, about which they move; to this bolt is also fixed an iron half ring to hook on a windless; when the gin stands upright so as the legs stand at a proper distance, one end of the cable is fastened to the gun, and the other passing through the pulleys and round the roller, the roller is turned round by means of hand-spikes passing through the holes in the ends of the roller; whilst a man holds the cable tight, the gun is raised to the height that the carriage may be put under it.

GRAPE-SHOT, is a certain number of small iron shot in a bag of coarse cloth, just to hold a bottom of iron or wood with a pin in the middle; about this pin are put as many shot as the grape is to consist of; then they are quilted with a strong packthread in order to keep the shot tight; their size is according to the bore of the gun, and their weight nearly equal to that of the shot of that gun.

Small iron or lead, the size of musket shot, are also put into tin boxes, to serve in the same manner as grape-shot.

GRAPLING IRONS, is composed of four, five, and six branches, bent round and pointed, with a ring at the root, to which is fastened a rope to hold by, when the grapple is thrown at any thing, in order to bring it near so as to lay hold of it.

GUIDEFORE, in a waggon, are two crooked pieces of wood fixed upon the axle-tree each with an iron bolt, joined before by a cross bar, to which the shafts are fastened with a bolt and key, and upon the hind part is fastened the sweep bar.

GUIDEHIND, is likewise composed of two crooked pieces of wood fastened to the hind axle-tree as the former, and fastened to the pole with a flat ring and an iron bolt.

GUN, the length is distinguished by three parts; the first reinforce, the second reinforce, and the chace; the first

reinforce is two-sevenths, and the second one-seventh and half a diameter of the shot.

The inside hollow wherein the powder and shot are lodged, the bore and the diameter of the bore is called the diameter of the caliber.

The part between the hind end and the bore, the breech; and the fore-part of the bore, the mouth.

The cascable is the part terminated by the hind part of the breech, and the extremity of the button.

The trunnions, the cylindric parts of metal which project on both sides of the gun, which rest in the grooves, made in the side-pieces of a carriage.

The mouldings are those behind the breech, and are looked upon to belong to the cascable, the first and second reinforce rings, ogees, astragals, and fillets.

Those of the first reinforce are a ring ogee joining to it, and an astragal with fillets; the part of the gun between the ogee and astragal is called the vent-field, because the vent is placed there.

The ogee of the second, a ring and ogee; and those of the chace, a ring-ogee; the astragal with fillets, the muzzle astragal, the swelling of the muzzle an ogee, or cimaise and two fillets.

The part between the ogee and chace astragal, the chace girdle; and the part from the muzzle astragal and the mouth, the muzzle.

Formerly guns were distinguished by the names of makers, culverins, cannon, demi-cannon, &c. but at present their names are taken from the weight of their shot; as for example, a 12 or 24 pounder carries a ball of 12 or 24 pounds weight.

Guns are made of brass or cast iron; the brass is a mixture of copper and tin: sometimes yellow brass is added, but is reckoned to make the metal brittle. The most common proportion is to an hundred pounds of copper twelve pounds of tin. But as copper requires a red heat to melt, and tin does melt in a common fire, when a gun is much heated by firing, the tin melts or softens

softens so much that the copper alone supports the force of explosion, whereby they generally bend at the muzzle, and the vent widens so much as to render the gun useless. If such a composition of metal could be found as required an equal degree of heat to melt, it would answer the intent: but as no such thing has been hitherto found, I look upon good iron to make better and more durable guns than any other composition whatever, as experiments and practice have shewn. For all our brass battering guns made use of this last war were soon rendered unserviceable, and iron ones substituted.

The necessary tools for loading and firing guns, are rammers, sponges, ladles, worms, hand-spikes, wedges, or screws.

The rammer is a cylinder of wood, whose diameter and axis is equal to that of the shot, and serves to ram home the wadds put upon the powder and shot; the sponge is the same, only covered with lamb-skin, and serves to clean the gun when fired: the rammer and sponge are fixed to the same handle.

The ladle serves to load the gun with loose powder.

The worm serves to draw out the wads when a gun is to be unloaded.

The hand-spikes serve to move and to lay the gun.

The coins or wedges, to lay under the breech of the gun, and to raise or depress it.

In field-pieces, a screw is used instead of coins, by which the gun is kept to the same elevation.

The tools necessary to prove guns, besides those mentioned for loading them, a priming iron, a searcher with a reliever, and a searcher with one point.

The first searcher is an iron, hollow at one end to receive a wooden handle, and on the other has from four to eight flat springs of about six inches long, pointed and turned outwards at the ends: the reliever is an iron flat ring, with a wooden handle, at right angles to it: when a gun is to be searched after it has been fired, this searcher is introduced and turned every way from one
end

end to the other; and if there is any hole, the point of one or the other spring gets into it, and remains till the reliever, passing round the handle of the searcher, presses the springs together and relieves it; and if any of the points catches in the vent, the priming iron is introduced to relieve it.

When there is any hole or roughness in the gun, the distance from the mouth is mark'd on the outside with chalk.

The other searcher has also a wooden handle and a point at the fore end of about an inch long, at right angles to the length about this point is put some wax mixed with tallow, and when introduced into the hole or cavity, is pressed in and drawn forwards and backwards, then the impression upon the wax gives the depth, and the length is known by the motion of the searcher: if the hole is a quarter of an inch deep, and downwards, the gun is rejected.

GUNNER, is the second rank of private men in the regiment of artillery.

HAIR-CLOTH, used on the floor of powder magazines, to prevent accidents of fire from mens shoes treading or rubbing upon or against sand or gravel, &c. may strike fire.

HAND-GRENADES, a small iron shell from 2.5 to three inches diameter, filled with powder, which being lighted by means of a fuze, are thrown by the grenadiers amongst the enemy when they come near.

HAND SCREWS, are composed of a toothed iron bar, that has a claw at the lower end and a fork at the upper; this bar is fixed in a stock of wood about 2.5 feet high, and moved by a rack-work; this claw or fork being placed under a weight raises it as far as the bar can go.

HAND-SPIKE, a wooden leaver five or six feet long, flattened at the lower end, and tapering towards the other, useful in moving guns to their place, after being fired and loaded again.

HAND-

HAND-MALLETS, a wooden hammer with a handle, to drive fuzes and picketts, in making fascine batteries.

HAMMER, IRON, made of various sizes and forms, the common have a head on one end and a claw on the other, to drive nails in or to draw them out.

HASP, a flat staple to catch the bolt of a lock.

HATCHET, for cutting wood to make fascines and pickets.

HEAD, the fore part of the side pieces of a gun carriage.

HELVE, or **HAFT**, the wooden hands of hatchets, hammers, or pick-axes, &c.

HINGES, are two iron bands, with a joint, nailed to the doors or lockers to fasten them, and move them back and forwards.

HOOKS, those which are fixed to the transom plates of a field-carriage; they serve to fix chains for drawing it occasionally back and forwards.

HOOPS, IRON, of several sorts, such as nave and axletree hoops.

HORN-BEAM, a wood much used for making the fuzes of shells.

HORNS, POWDER; the gunners sling them over the shoulders with a belt, to prime the guns or mortars.

HOWITZ, a sort of mortar, mounted upon a field-carriage like a gun; the difference between a mortar and a howitz is, that the trunnions of the first are at the end, and in the middle in the last.

HURTERS, a flatted iron fixed against the body of an axletree with straps, to take off the friction of the naves of wheels against the body.

HURTER, of wood, is a piece of five by six inches square, and eight feet long, laid upon a platform against the breast work, to stop the wheels from damaging the fascines.

JOINT-BOLTS, by which one end of a cap square is fixed to the carriage sea bolts.

IRON-GUNS, first made of hammered, but now of cast-iron.

IRON,

IRON, PRIMING, which the gunners carry with their horns to clear the vent, or to release the searcher in the proving of guns.

IRON, for marking with a broad R the horses that are taken for the government's service to draw the artillery ammunition and stores.

KETTLE OF COPPER, to boil compositions for fire-works.

KEYS, FORE-LOCK, serve to pass through the lower end of bolts, to fasten them.

KEYS, with chains and staples, fixed on the side-pieces of a carriage or mortar beds, they serve to fasten the capsquares, by passing through the eyes of the eye-bolts.

KEYS, SPRING, they serve the same purposes as the former, but instead of being one single piece they are of two, like two springs laid one over another; when they are put into the eye-bolts they are pinched together at the ends, and when they are in, open again, and cannot shake out by the motion of the carriages; they are used in travelling carriages.

LABORATORY is the place where the different materials for fire-works are lodged and prepared. In the field a tent is used for that purpose pitched on a place where the enemies shot or shells cannot approach.

LADLES, made of copper to hold the powder for loading of guns, with long handles of wood, when cartridges are not used.

LADLES, SMALL, of copper; with short handles, they fill the fuses of shells, or any other composition to fill the cases of sky-rockets, &c.

LANTHORNS, Muscovy, dark and common, used in the field when dark to light the gunners in the camp to prepare the stores.

LASHING RINGS, with hoops, fixed on the side-pieces of travelling-carriages to lash the tarpauling that covers the gun, as also to tie the sponge, rammer, and ladle.

LIMBER, a two-wheel carriage with shafts to fasten the trail of travelling carriages by means of the pintle or iron pin, when travelling, and taken off on the battery, or in the park of artillery, which is called unlimbering the guns.

LINCH-PIN, that which passes through the ends of the arms of an axle-tree, to keep the wheels or trucks from slipping out in travelling.

LINCH CLOUT, the flat iron under the ends of the arms of an axle-tree to strengthen them, and to diminish the friction of the wheels.

LINE of DIRECTION, formerly marked upon guns, by a short point upon the muzzle and cavity on the base ring, to direct the eye in pointing the gun, but are left off at present for no substantial reason.

LINT STOCK, the stick used by the gunners to fasten the match whereby they fire the guns or mortars.

LOCKS, of various sorts, common for lockers in travelling carriages, or for boxes containing shot and powder or cartridges.

LOCKER HINGES, serve to fasten the cover of the lockers in travelling carriages.

LOCKING-PLATES, thin flat pieces of iron nailed on the sides of a field carriage where the wheels touch it in turning, to prevent the wearing the wood in those places.

LOOP, in a ship carriage, made of iron, fastened one on the front of the fore axle-tree, and two on each side, through which the ropes or tackle pass, whereby the guns are moved backwards and forwards on board of ships.

MARLIN, tarred, skains, white; are long wreaths or lines of untwisted hemp, dipped in pitch or tar, with which cables and other ropes at sea are wrapped round to prevent their fretting or rubbing in the blocks or pullies through which they pass.

They serve the same purpose in artillery upon ropes used for draught for rigging gins, usually put up in small parcels called skains.

MATCH,

MATCH, used in artillery for firing guns and mortars; it serves likewise to trace batteries instead of ropes.

MATROSS, properly an apprentice to a gunner, and has the least pay of any soldier in the artillery.

MEASURES, POWDER, made of copper, holding from an ounce to eight or twelve pounds, and are very convenient in a siege when guns or mortars are loaded with loose powder, especially in ricochet firing.

MEDICINE CHEST, serves the surgeons to carry their medicines and instruments when in the field.

METAL, GUN, a composition of tin and copper; the most common proportion is, to a hundred pounds of copper twelve pounds of tin: the founders will sometimes vary from this proportion according to the hardness of the copper.

MORTARS, made of brass or iron, used in the land service and at sea, for throwing shells or carcasses; the land are shorter and lighter, and the chambers hold less powder.

They are distinguished by the diameter of their bores; thus, a thirteen, ten, or eight inch mortar are those whose diameter of their bores are thirteen, ten, or eight inches, the royal and coehorn excepted; the royal carries a shell whose diameter is 5.5 inches, and that of the coehorn whose diameter is 4.6 inches.

MORTAR, SEA, are longer to prevent the setting the rigings of the ships on fire, and heavier to lessen the shock upon the ship when the mortar is fired: none but thirteen and ten inch mortars are used at sea; and as these mortars are sometimes fired at a great distance, the chamber of the first holds 32 pounds of powder, and that of the second 12.5 pounds.

MOULDS, for casting shot for guns, pistols, muskets, and carabines; the first are of iron used by the founders, and the others by the artillery in garrison and the field.

MOULDS, of wood or brass, used in laboratory works for filling and driving all sorts of rockets, and cartridges; of many different sizes.

NAILS of various sorts are used in artillery.

1. **GARNISH**, in travelling carriages, have pointed heads finished like diamonds, with a small narrow neck; they serve to fasten the plates with roses, to cover the side-pieces from the ends of the trunnion-plates to five or six inches beyond the center of the carriage.
2. **DIAMOND-HEADED**, small nails whose heads are made like a flat diamond, and serve to fix the plates upon travelling carriages, besides the garnish nails, both upon the trail, and those at the head are near the trunnion-plates.
3. **ROSE-BUDS** are small round-headed nails drove in the center of the roses of the plates.
4. **COUNTER-SUNK**, those that have flat round heads sunk into the iron plates so as to be even with the outside of it.
5. **STREAK-NAILS**, are those which fasten the streaks to the fellows of the wheels.
6. **BOX-PINS**, small nails without heads to pin the nave-boxes to the naves.
7. **STUBS** are driven on the outside of the nave hoops to keep them in their places.
8. **FLAT-HEADED NAILS** to fasten the locker or any hinges.
9. **DOG NAILS** have flat round heads, and one part of the shank next to the head is also round.

NAVE, that part of a wheel in which the arms of the axle-tree moves, and in which the spokes are drove and supported.

NAVE HOOPS are flat iron rings to bind the nave; there are generally three on each nave.

NAVE BOXES, there are two, one at each end, to diminish the friction of the axle-tree against the nave: they were formerly made of brass, but experience has shewn that the cast iron cause less friction and are much cheaper.

OGEE,

OGEE, an ornamental moulding in the shape of an S, taken from architecture and used in guns, mortars, and howitzes.

ORDNANCE, a name given to all that concerns artillery; thus, the commander in chief is called Master General of the Ordnance instead of artillery; the Lieutenant General of the Ordnance.

ORDNANCE, Board of, consists of four officers, the surveyor general, clerk of the ordnance, store-keeper, and the clerk of deliveries, over which presides the master, or, in his absence, the lieutenant-general: this board deliberates, regulates, and orders every thing relating to the artillery.

Guns, mortars, and howitzes are often called ordnance.

PAILS, made of wood bound with iron hoops and handles, hold generally four gallons, and serve in the field to fetch water for the use of artillery works.

PANNELS are the carriages which carry mortars and their beds upon a march.

PETARD, an engine to burst open the gates of small fortresses, is made of gun metal, fixed upon a board two inches thick and about two feet square, to which it is screwed, and holds from nine to twenty pounds of powder, with a hole at the end opposite to the plank to fill it, into which the vent is screwed; the petard thus prepared is hung against the gate by means of a hook, or supported by three staves fastened to the plank, when fired bursts open the gate.

PICKETS, for several uses; those for pinning the fascines of a battery, are from three to five feet long, the heads two or three inches diameter.

PICKETS, used in the park of artillery, are about six or eight inches long, shod with iron, to pin the park line.

PICKET, an out-guard posted before an army, to give notice if an enemy approaches.

PINTLE, a long iron bolt fixed upon the middle of the limber bolster, to go through the hole made into the

trail

trail transom of a field carriage, when it is to be transported from one place to another.

PINTLE-PLATE, is a flat iron through which the pintle passes, and nailed to both sides of the bolster, nailed with eight diamond-headed nails.

PINTLE-WASHER, an iron ring through which the pintle passes, placed close to the bolster for the trail to move upon.

PINTLE-HOLE, is of an oval figure made in the trail transom of a field carriage, wider above than below, to leave room for the pintle to play in an uneven road.

PINTLE-PLATES, one above and one below the pintle hole, with a ring at the end, nailed to the trail transom with diamond-headed nails.

PIN, an iron nail or bolt, with a round head, and generally has a hole at the end to receive a key; there are many sorts, as axle-tree pins or bolts, bolster pins, pole pins, swing-tree pins, &c.

PLATFORMS, of a battery, is the bed of wood upon which the guns stand; each consisting of eighteen planks of oak or elm, a foot broad, two inches and a half thick, and from eight to fifteen feet long, nailed or pinn'd on five beams of about four inches square, called sleepers.

The platform must be higher behind than before by six or nine inches, to prevent too great a recoil, and to advance the gun easily when loaded.

POLE, in a four wheel carriage, is fastened to the middle of the hind axle-tree, and passes between the fore-axle-tree and its bolster, fastened with the pole pin so as to move about it, and fastens the fore and hind carriages together.

POINT-BLANK, is said of a gun when laid level; and point-blank range the distance the shot goes upon a level plain.

PONTOON, a flat-bottomed boat, whose carcase of wood is lined within and without with tin; they serve to lay bridges over rivers for the artillery and army.

The

The French pontoons are made of copper, and only on the outside; though they cost more at first, yet they last much longer than those of tin, and when wore out the copper sells nearly for as much as it cost at first.

PONTOON CARRIAGE, made of two wheels and two long side-pieces, whose fore ends are supported by a limber, and serves to carry the pontoon, boards, and the pieces of timber laid across for making the bridge.

PORT FIRE, a composition put in a paper case to fire guns and mortars upon particular occasions, or to light recreative fire-works.

POUNDER, said of guns to specify a certain caliber; thus, a twenty-four, twelve, or a six-pounder are those pieces whose balls weigh twenty-four, twelve, or six pounds.

POWDER, GUN, a composition of saltpetre, sulphur, and charcoal; the usual proportion is, to six of saltpetre, one of sulphur, and one of charcoal, well mixed and grained, and is put into barrels which contain 100 pounds each, for the use of artillery.

POWDER-CART, a two-wheel carriage, covered with an angular roof of boards, and to prevent the powder from dampness a tarred canvass is put over the roof; on each side are lockers to hold shot in proportion to the quantity of powder, which is generally four barrels.

The French use powder-waggons, which carry 12 barrels of 200 pounds each, but then the shots are carried by themselves on other carriages.

POWDER MAGAZINE, a building to hold the powder in fortified places, containing several rows of barrels laid one over another, and arched over so as to be bomb-proof.

POWDER MILL, where the materials are beat, mixed together, and grained; they are placed near rivers, and as far from any house as can be, for fear of accidents, which often happen.

PROOF

PROOF OF FIRE ARMS, the general rule for proving guns and mortars is to fire them three times with double the quantity of powder as they are loaded with in common service; but the rule of the ordinance is, that all guns under 24 pounders are loaded with powder as much as their shot weighs; a brass 24 pounder with 21 lb. a brass 32 pounder with 26 lb. 12 ounces, and a 42 pounder with 31 lb. 8 ounces; the iron 24 pounder with 18 lb. the 32 with 21 lb. 8 ounces, and the 42 with 25 lb.

The brass light field-pieces are proved with powder that weighs half that of their shot, except in the 24 pounder, which is loaded with 10 lb. only.

The government allows 11 bullets of lead in the pound for the proof of muskets, and 14.5, or 29 in two pounds, for service; 17 in the pound for the proof of carabines, and 20 for service; 28 in the pound for the proof of pistols, and 34 for service.

When guns of a new metal or of lighter construction are proved, then, besides the common proof, they are fired two or three hundred times as quick as they can be, loaded with the common charge given in service. Our light six pounders were fired 300 times in three hours and twenty-seven minutes, loaded with 1 lb. and 4 ounces of powder.

QUADRANT, an instrument made of brass or wood, divided into degrees, and each degree into ten parts, to lay guns or mortars to any angle of elevation.

The common sort is a quadrant, one of the radii projects the quadrant of about a foot, and a plummet suspended in its center, by means of a fine thread of silk, so that when the long end is introduced into the piece the plummet shows its elevation.

The best sort has a spiral level fixed to a brass radius, so when the long end is introduced into the piece, this radius is turned about its center till it is level, then its end shews the angle of elevation, or the inclination

clination from the horizon, whereas the first shews that angle from the vertical.

QUICK MATCH, made of three cotton strands drawn into lengths, and put into a kettle just covered with white wine vinegar, and then a quantity of saltpetre and mealed powder is put in it, and boiled till well mixed. Others put only saltpetre into water, and after that take it out hot and lay it into a trough with some mealed powder, moistened with some spirits of wine, thoroughly wrought into the cotton by rolling it backwards and forwards with the hands; and when this is done, they are taken out separately, drawn through mealed powder, and dried upon a line.

RANGE, the distance from the battery to the point where the shot or shell touches the ground.

RANGE, POINT BLANK, that when the piece lies in a horizontal direction and upon a level plane.

RANGE, RANDOM, when the piece is elevated at an angle of elevation of 45 degrees upon a level plane.

RAVES, are the upper wooden bars in a cart or waggon, supported by the round and flat staves which enter into them.

REGIMENT OF ARTILLERY, is composed of three battalions, and (a company of cadets) each battalion of ten companies, and in time of war two companies of miners are added to it.

REINFORCE, the part of a gun next to the breech, which is made stronger to resist the force of the powder: there are generally two, called the first and second; the second is something smaller than the first, upon the supposition that when the powder is inflamed, and occupies a greater space, its force is diminished: we have shewn, in our construction of guns, the absurdity of this practice.

REINFORCE-RING; there are three in each gun, called the first, second, and third; they are flat mouldings, like flat iron hoops, placed at the the breech end of

the first and second reinforce, projecting the rest of the metal by about a quarter of an inch.

RELIEVER, an iron ring fixed to a handle by means of a socket, so as to be at right angles to it: it serves to disengage the searcher of a gun, when one of its points are retained in a hole, and cannot be got out otherwise.

RIDER, a piece of wood something higher than broad, the length equal to that of the body of the axletree upon which the side-pieces rest in a four-wheel carriage, such as the ammunition waggon, block carriage, and sling waggons.

RINGS, of various uses, such as the lashing rings in travelling carriages to lash the sponge, rammer, and ladle, as well as the tarpauling that covers the guns; the rings fastened to the breeching bolts in ship-carriages; and the shaft-rings to fasten the harness of the shaft horse by means of a wooden pin.

RIVETTING-PLATES, small square thin pieces of iron, through which the ends of bolts pass, and rivetted upon them.

RODS, iron or wood, rammers, to drive the charges of muskets, carabines, and pistols home to the bottom of the pieces.

RODS, or **STICKS**, fastened to sky rockets, to make them rise in a strait line.

ROCKETS, made for recreative fire-works, both for the air and water; the sky have sticks fixed to them.

ROPES, of various lengths and thickness, according to the uses they are made for; such as drays, for the gin, for the sling cart and waggon.

SALTPETRE, the principal ingredient for making gunpowder, is found in great plenty in some of the Eastern provinces and in Europe, and is extracted from the earth dug out of cellars and stables, or else near old walls where urine has been made.

SALTING BOXES, for holding mealed powder to throw upon the shells, when in the mortar, in order that

the fuzes may take fire from the blast of the powder; but it has been found that the fuze takes fire without mealed powder by which these boxes are rejected.

SAWS, for cutting wood, very useful in artillery, to cut planks of travelling carriages or platforms, and in general for all carpenters work, &c.

SCALES, of several sorts, large and small, to weigh from a grain to four or five tons.

SCALING LADDERS, used in scaling when a place is to be taken by surprize: they are made several ways; here we make them with flat staves, so as they may move about their pins, and shut like a parallel ruler, for conveniently carrying them. The French make them of several pieces, so as to be joined together, and to be made of any necessary length; sometimes they are made of single ropes knotted at proper distances, with iron hooks at each end, one to fasten them upon the wall above, and the other in the ground; and sometimes they are made with two ropes, and staves between them to keep the ropes at a proper distance, and to tread upon.

SCREWS, fastened to the cascable of a light gun or howitz, by means of an iron bolt, and goes through a socket fixed upon the center transom, to lay the piece with instead of wedges, as is done in heavy pieces.

SEA MORTARS, are those put on board of bomb ketches, to be used in the bombarding of sea-ports; each ketch carries two, one of thirteen and one of ten, no less are used.

SETTERS, a round stick to drive fuzes, or any other compositions, into cases made of paper.

SHAFT-RINGS, see rings.

SHAFTS of a carriage, are two poles joined together with cross bars, whereby the hind horse guides the carriage, and supports the forepart of the shafts, the hind turning round an iron bolt.

SHAFT BARS, are two pieces of wood to fasten the hind ends of the shafts together, into which they are pinned with nails of wood.

SHELLS, a hollow iron ball to throw out of mortars or howitzes, with a hole of about an inch diameter, to load them with powder and to receive the fuze; the bottom or part opposite the fuze is made heavier than the rest, that the fuze may fall uppermost; but in small elevations is not always the case, nor is it necessary; for when it falls first, it sets fire to the powder in the shell. But whether it breaks or not, it would be better to make the shell every where of the same thickness, because it would then burst into a greater number of pieces than it does now.

SHEEP-SKINS, serve to cover the mortars or howitzes between firing, to prevent any wet or dampness getting in them.

SHOT-IRON, balls to load guns with; the shot of lead to load muskets, carabines, or pistols, which are called bullets.

SHOT-GRAPE, a certain number of small shots of iron or lead, quilted together with canvas and ropes about a pin of iron or wood, fixed upon a bottom of the same matter, so as the whole together weigh nearly as much as the shot of that caliber.

SHOVEL, a miner's or pioneer's tool to dig and throw up earth.

SIDE-PIECES, of gun carriages, see carriages.

SIDE-STRAPS, in a field carriage, are flat iron bands which go round the side-pieces, in those places where the wood is cut cross the grain, to strengthen them near the center and the trail.

SLEETS, are the parts of a mortar going from the chamber to the trunnions, to strengthen that part, see plate II. fig. 5. by the letter n.

SLING-CART, or waggon, serves to carry guns or mortars a small distance, by means of a rope fastened with one end round the middle of the pieces, and wound upon a roller, by means of handspikes, or a rackwork.

SOMMERS,

SOMMERS, in an ammunition waggon are the upper sides, supported by the staves entered into them with one of their ends and the other into the side-pieces.

SPIKES-HAND, is a leaver of four or five feet long, and about two and a half inches diameter, tapered on one end, and flat on the other, to be put under a gun or mortar to raise it from the ground, so as to put a skid or piece of wood under it.

SPONGE, of a gun or mortar, is a cylinder of wood fixed to a handle, and whose diameter and length are equal to the diameter of the shot in guns, covered with a lambskin to clean the piece with after it has been fired.

SPOKES of a wheel, are the roundish parts between the nave and fellows, or circumference of the wheel.

STAPLES, are drove in the side-pieces of gun and mortar carriages, to fasten the keys of eye-bolts by means of chains; there are also staples to the locker of gun carriages, and for several other uses.

STAVES, round and flat, used in ammunition and other waggons or carts, the round and flat sticks between the sommers and side-pieces, or else in common or scaling ladders to tread upon.

STAYS, axle-tree, in truck carriages, are the irons which are fixed one end under the fore axle-tree, and the other to the side-pieces, in the form of an S.

STEEL-YARD, a sort of scale whereby many things are weighed, by means of two weights only, one fixed near the end, and the other sliding along the beam upon which the weights are marked, according to the distances from the point of suspension.

STONE-MORTAR, made lighter than the common, with a cylindric space near the chamber to hold a basket filled with stones, which are flung amongst an enemy, instead of shells or grapes, when near.

STOOL-BED, in a truck carriage, serves as a stool for the guns to lay upon, and the wedges for raising them.

STREAKES, are the iron bands on the outside of a wheel to bind strongly the fellows together.

STRAKE-

STREAKE NAILS, are those drove through the streakes into the fellows.

SULPHUR, or brimstone, a mineral, very useful in making of powder, or in artificial fireworks.

SWEEP-BAR, of a waggon, is that which is fixed on the hind part of the fore-guide, and passes under the hind pole, which slides upon it.

SWING-TREE, the bar in a carriage fastened across the fore-guide, and to which the traces of the horses are fastened.

TABLES, a kind of register to set down the dimensions of carriages for guns, mortars, &c.

TARPAULINGS, are made of strong canvas thoroughly tarred, and cut into different sizes according to their several uses in the field; such as to cover the powder waggons, and tumbrils carrying ammunition from rain; each field-piece has likewise one to secure their ammunition boxes.

TENTS, made of canvas, for officers and soldiers to lay under when in the field.

TIN-TUBES, are used in guns to set fire to the powder instead of priming. Their advantage over priming is that they preserve the vent from wearing.

TOMPEONS, or tampion, a wooden cylinder to put into the mouth of guns in travelling, to prevent the dust or wet getting in.

Sometimes used to put into the chambers of mortars over the powder when the chamber is not full.

TONGS, of a waggon, a piece of wood fixed between the middle of the hind ends of the shafts, morticed into the fore-cross-bar and let into the hind-cross-bar.

TOOLS, of all sorts, for artificers, miners, laboratory, &c.

TRAIL, is the end of the travelling carriage opposite to the wheels, and upon which the carriage slides, when unlimbered, or upon the battery.

TRANSOM,

TRANSOM, the pieces of wood which join the side-pieces of gun carriages: there is but one in a truck carriage placed under the trunnion holes, and four in the wheel carriages; the trail, the center, the bed, and the breast transoms.

TRANSOM-PLATES, with hooks; there is one on each side of the side-pieces against each end of the transom, the bed-transom excepted, fastened by two transom-bolts.

TRANSOM-BOLT, with burrs; they serve to tie the side-pieces to the transoms.

TRAVELLING CARRIAGES, see gun carriage.

TRAVELLING FORGE, see forge.

TRAVERSING PLATES, two thin iron plates nailed on the hind part of a truck carriage of guns, where the handspike is used to traverse the gun.

TRUCKS of a ship carriage, are wheels made of one piece of wood, from twelve to nineteen inches diameter, and their thickness is always equal to the caliber of the gun.

The trucks of garrison carriages are made of cast iron.

TRUCK CARRIAGE, goes upon four trucks of twenty-five inches diameter, has two flat side-pieces of ten inches broad, and serves to carry guns, ammunition, boxes, or any other weights from the store-houses to the water side, or to any small distance.

TRUNNIONS, two cylindric pieces of metal in a gun or mortar which project the piece, and by which they are supported upon their carriages.

TRUNNION PLATES, are two plates in travelling carriages which cover the upper parts of the side-pieces, and goes under the trunnions, its fore-part ends with a flat rose.

TUGPINS, are the iron pins which pass through the fore ends of the shafts, to fasten the draught-chains for the fore horses.

VENT,

VENT, of all kinds of fire-arms, is a small hole pierced at the end or near it of the bore or chamber, to prime the pieces with powder, or to introduce a tube, in order when lighted to set fire to the charge.

VENT-FIELD, is the part of a gun or howitz between the breech moulding and the astragal.

VENT-ASTRAGAL, that which determines the vent-field.

WADD, made of hay or straw, and sometimes of tow rolled up tight in a ball, and serves to be put into a gun after the powder, and rammed home, to prevent the powder from being scattered, which would have no effect when unconfined.

WADD-MILL, a hollow form of wood to make the wadds in of a proper size.

WADDING, hay, straw, or any other forage generally carried along with the guns to be made into wadds.

WADD-HOOK, a strong iron screw, like those that serve for drawing corks, mounted upon a wooden handle, to draw out the wadd, or any parts of cartridges, which often remain in guns, and when accumulated stops up the vent.

WASHER, a flat circular ring put on the axle-tree, between the linch-pin and small end of the nave, to prevent the nave rubbing against the linch-pin and wearing it, as likewise to diminish the friction of the nave.

WEIGHTS, are of two sorts, troy weight is the least, and serves for all light goods, silver, gold, and medicine.

Averdupois serves for all heavy course commodities, and every thing is weighed by it in artillery.

WHEEL of a carriage, well known by every body; in artillery their strength is always, or should be, proportional to the weight they carry; the diameter of the wheels of heavy gun carriages is fifty-eight inches, and those for light field-pieces fifty-two only.

